

# Tunneling Leakage Current Dependent RDD Model Framework for Gate Oxide TDDB

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**Abstract:** The Reaction Diffusion Drift (RDD) model is used to simulate trap generation ( $\Delta N_{TG}$ ) kinetics during Time Dependent Dielectric Breakdown (TDDB) experiments. Several features of measured data, e.g., stress gate voltage ( $V_G$ ) dependence of mean time to breakdown ( $T_{BD}$ ) across temperature ( $T$ ) and gate oxide thickness ( $Tox$ ), Weibull slope ( $\beta$ ) across  $Tox$  are explained. The role of electron ( $J_E$ ) and/ or hole ( $J_H$ ) leakage current (with  $J_H$  from Anode Hole Injection or AHI) is explored as a trigger for RDD model. The polarity gap between positive and negative  $V_G$  stress is addressed.

**Keywords:** Tunneling leakage current, Anode Hole Injection, Reaction Diffusion Drift Model, Time Dependent Dielectric Breakdown.

## I. INTRODUCTION

TDDB is caused by gradual buildup of gate oxide defects, leading to formation of a percolation path between gate and channel and resulting increase in gate leakage current at time  $T_{BD}$  [1]-[13]. Percolation path formation is stochastic in nature, and  $T_{BD}$  shows Weibull distribution [2]. The  $V_G$  dependence of  $T_{BD}$  is modeled by different physical mechanisms, e.g., AHI [14], [15], Thermo-Chemical (TC) [16], or Anode Hydrogen Release (AHR) [17]. However, none addressed the trap generation time kinetics (Stress Induced Leakage Current or SILC) and resulting failure when a critical defect density ( $N_{BD}$ ) is reached. In this paper, a leakage current triggered RDD model is used to address the same and bridge the gap in co-simulation of SILC and TDDB.

## II. GATE LEAKAGE CURRENT

Fig.1 shows the band diagrams for (a) Fowler-Nordheim (FN) tunneling and (b) Direct Tunneling (DT) of electrons from cathode to anode via (a) thick and (b) thin gate oxide, energy gain from oxide field, impact ionization to generate holes in anode, and injection of holes into gate oxide for NMOS in (a) inversion ( $V_G > 0$ ) and (b) accumulation ( $V_G < 0$ ) stress. The extra energy gain across the bandgap of the p-type substrate (i.e., minority ionization, Fig.1(b)) results in a higher hole-to-electron current ratio, Fig.2, demonstrating the impact of stress voltage polarity on AHI process [18].

## III. RDD MODEL

H-passivated bonds are broken ( $K_{F1}$ ) via injection of electrons and/or holes, oxide field ( $E_{OX}$ ) and temperature ( $T$ ), released H atoms diffuse and subsequently break other bonds to generate  $H_2$  molecules ( $K_{F2}$ ) and  $H_2^+$ ,  $OH^-$  ions ( $K_{F3}$ ) that diffuse or drift away, Fig.3. Broken bonds results in bulk traps (density  $\Delta N_{TG}$ ), whose magnitude at a given time is governed by  $K_{F1}$ , and long-time power-law time kinetics slope ( $n$ ) is

governed by the relative ratio of molecules to ions ( $K_{F2}/K_{F3}$ ), Fig.4. The mean  $T_{BD}$  is calculated as the time taken to reach  $N_{BD}$  [2], [14]. Except  $K_{F1}$ , the parameters (reaction rates and diffusivities) are Arrhenius  $T$  activated, and except  $K_{F1}$  and  $K_{F30}$  pre-factor of  $K_{F3}$ , all other parameters are kept fixed across all cases studied in this paper.

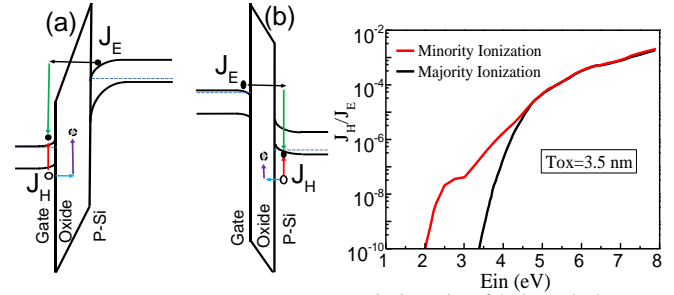


Fig.1 Energy band diagram showing electron tunneling, impact ionization, hole injection into bulk, and defect generation during (a) inversion and (b) accumulation for NMOS capacitor.

Fig.2 Ratio of hole and electron current density ( $J_H/J_E$ ) versus injection energy ( $E_{in}$ ) [18] showing the impact of bias polarity on majority and minority ionization.

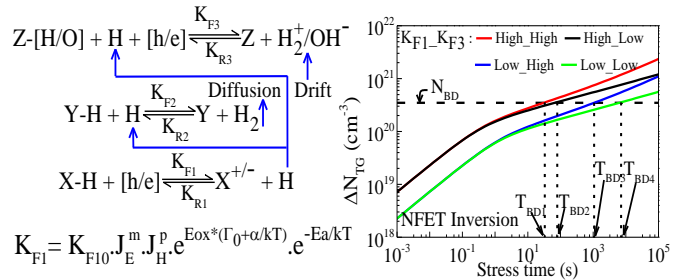


Fig.3 RDD model chemical equations showing diffusion of molecular, drift of ionic species, and first interface reaction rate  $K_{F1}$  dependence on tunneling leakage current, electric field, and temperature (bottom).

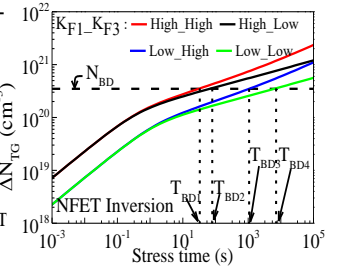


Fig.4 RDD model simulated bulk trap ( $\Delta N_{TG}$ ) kinetics for different  $K_{F1}$  (magnitude) and  $K_{F3}$  (slope,  $n$ ) demonstrating different mean time to breakdown ( $T_{BD}$ ) as it reaches critical defect density ( $N_{BD}$ ).

## IV. PERCOLATION MODEL

Fig.5 plots Weibull  $\beta$  versus  $Tox$  from various reports. The percolation model suggests  $\beta = n \cdot Tox/a_0$  ( $a_0$  being cell size) [2], [19], the  $\beta$  trend in Fig.5 can be modelled by  $Tox$  dependence of  $n$  and  $a_0$  as shown in Fig.6 and Fig.7 respectively. The  $n$  values are consistent with reported SILC slopes (after correction for sensing delay related discharging effect) at different  $Tox$ , and  $K_{F30}$  of RDD model can be varied to vary  $n$ . The reduction in  $a_0$  at lower  $Tox$  is reported elsewhere. Fig.8 plots the  $Tox$  dependence of  $N_{BD}$  from various reports (with fixed  $a_0$ ), and calculated by analytical percolation model [19] with fixed  $a_0$  and varying  $a_0$  as per Fig.7. A large  $N_{BD}$  reduction

at lower  $T_{ox}$  is observed for fixed  $a_0$ , the reduction of  $N_{BD}$  is much less when  $a_0$  reduces at lower  $T_{ox}$ .

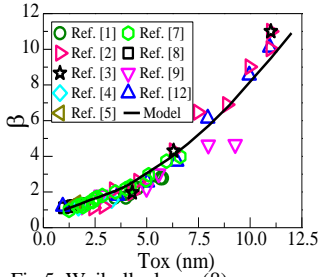


Fig.5 Weibull slope ( $\beta$ ) versus oxide thickness ( $T_{ox}$ ) from various reports and calculation ( $\beta=n*T_{ox}/a_0$ ), slope ( $n$ ) and cell size ( $a_0$ ) are taken from Fig.6 and Fig.7 trend, respectively.

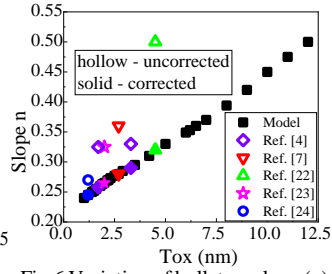


Fig.6 Variation of bulk trap slope ( $n$ ) and reported SILC slopes (after correction for sensing delay related discharging effect) with oxide thickness ( $T_{ox}$ ) from various reports and general trends.

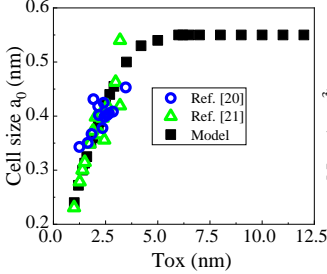


Fig.7 Variation of cell size ( $a_0$ ) as per percolation model with oxide thickness ( $T_{ox}$ ) from various reports and general trends.

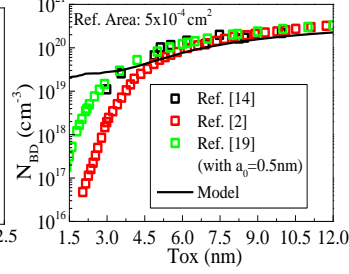


Fig.8 Mean critical defect density ( $N_{BD}$ ) versus  $T_{ox}$ . The Model line is calculated using nine neighbor model of [19] considering  $T_{ox}$  dependent  $a_0$  as shown in Fig.7.

## V. MODELING OF MEAN $T_{BD}$

RDD model is used for time kinetics of  $\Delta N_{TG}$  with  $K_{F1}$  related to  $J_E$ ,  $J_H$ ,  $E_{OX}$  and  $T$ , Fig.3, where  $K_{F10}$  is the reaction rate pre-factor,  $\Gamma_0$  is field acceleration factor,  $\alpha$  is bond polarization constant and  $E_A$  is the activation energy.  $J_E$  is calculated either from charge to breakdown ( $Q_{BD}$ ) data [6], [13] or FN and DT expressions [14], whereas  $J_H$  is calculated from the  $J_H/J_E$  ratio of AHI model, Fig.2 [18]. The  $K_{F1}$  term is generic and covers all physical cases, i.e., TC ( $m=p=0$ ), AHR ( $p=0$ ), and AHI ( $m=0$ ). The experimental and modeled mean  $T_{BD}$  versus  $V_G$  at a fixed  $T$  for various  $T_{ox}$  are shown in Fig.9 for NMOS inversion (NI), in Fig.10 for NMOS accumulation (NA), and in Fig.11 for PMOS inversion (PI). The experimental and modeled mean  $T_{BD}$  versus  $V_G$  at different  $T$  are shown in PMOS accumulation (PA) for relatively thin and thick  $T_{ox}$  respectively in Fig.12 and Fig.13. In all cases,  $\Delta N_{TG}$  time kinetics is calculated from RDD model, with a choice of  $K_{F30}$  to obtain  $n$  for a given  $T_{ox}$  as per Fig.7, and  $T_{BD}$  is noted when  $\Delta N_{TG}=N_{BD}$ . The pre-factor  $K_{F10}$  is shown versus  $T_{ox}$  for all cases (NI, NA, PI, and PA) for different choice of  $m$  and  $p$  in Fig.14. The variation of  $K_{F10}$  with  $T_{ox}$  for NA and PI is negligible compared to that for NI and PA which show opposite trends, Fig.14(a), for  $m=p=0$  (TC). For  $m=0.5$  and  $p=0$  (AHR), Fig.14(b),  $K_{F10}$  for NI has a decreasing trend with  $T_{ox}$  scaling, whereas, for  $m=0$  and  $p=0.5$  (AHI), Fig.14(c),  $K_{F10}$  increases with  $T_{ox}$  scaling for all cases.  $K_{F10}$  for all cases remains almost constant across  $T_{ox}$  for  $m=p=0.3$ , Fig.14(d), suggesting both  $J_E$  and  $J_H$  being responsible for initial dissociation of H bonds. A slight difference in  $K_{F10}$  is observed between NI versus NA and PI versus PA, which is possibly due to difference in gate oxide quality ( $\sim$ bond strength) near the poly-Si gate (NI and PA) and Si substrate (NA and PI), and doping of gate and substrate.

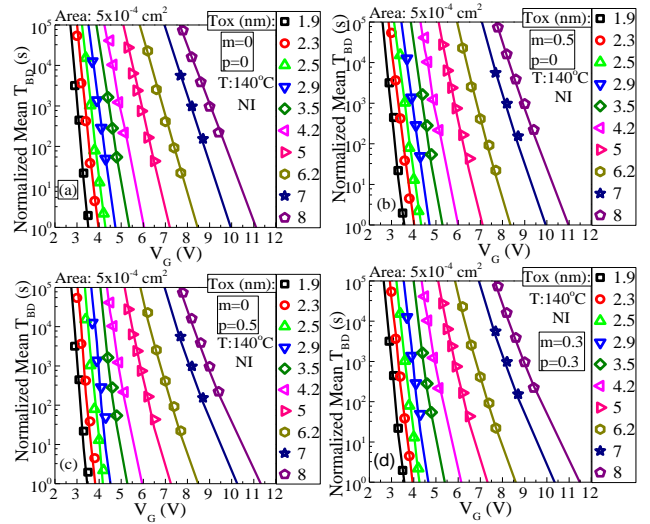


Fig.9 Mean  $T_{BD}$  versus  $V_G$  at a fixed  $T$  for NMOS capacitor with different  $T_{ox}$  stressed in inversion with  $K_{F1}$  (see the equation at the bottom in Fig.3) taking (a)  $m=p=0$ ; (b)  $m=0.5$ ,  $p=0$ ; (c)  $m=0$ ,  $p=0.5$ ; and (d)  $m=p=0.3$ . Data (symbols) from [6],[10],[11].

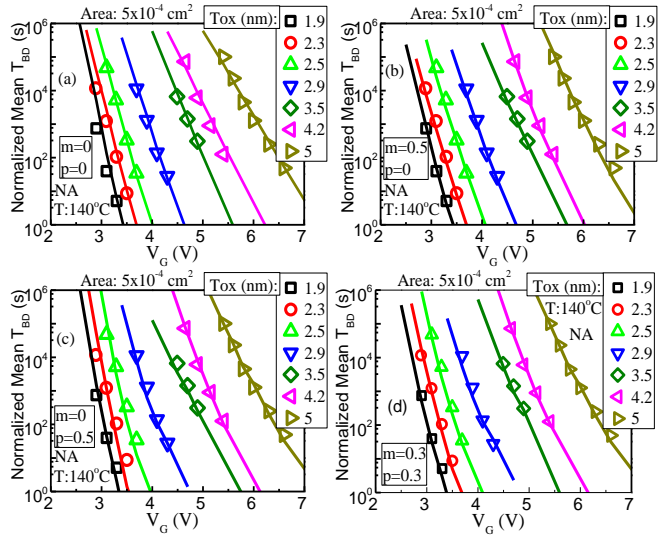


Fig.10 Mean  $T_{BD}$  versus  $V_G$  at a fixed  $T$  for NMOS capacitor in accumulation with different  $T_{ox}$  modeled with  $K_{F1}$  (see the equation at the bottom in Fig.3) taking (a)  $m=p=0$ ; (b)  $m=0.5$ ,  $p=0$ ; (c)  $m=0$ ,  $p=0.5$ ; and (d)  $m=p=0.3$ . Data (symbols) from [6].

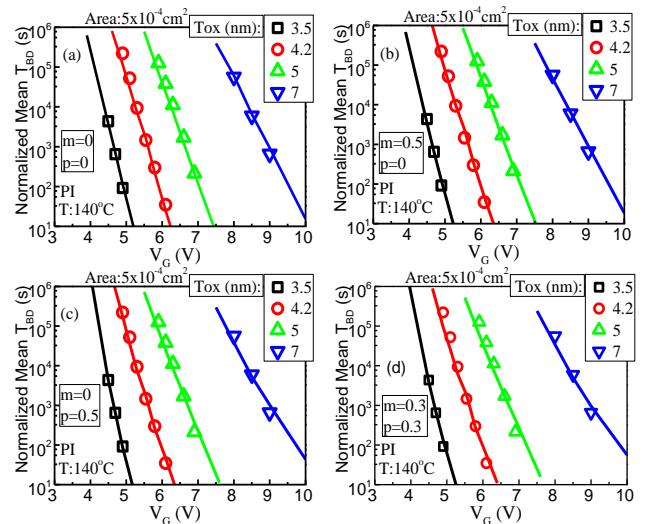


Fig.11 Mean  $T_{BD}$  versus  $V_G$  at a fixed temperature ( $T$ ) for PMOS capacitor with different  $T_{ox}$  stressed in inversion with  $K_{F1}$  (see the equation at the bottom in Fig.3) taking (a)  $m=p=0$ ; (b)  $m=0.5$ ,  $p=0$ ; (c)  $m=0$ ,  $p=0.5$ ; and (d)  $m=p=0.3$ . Data (symbols) from [10].

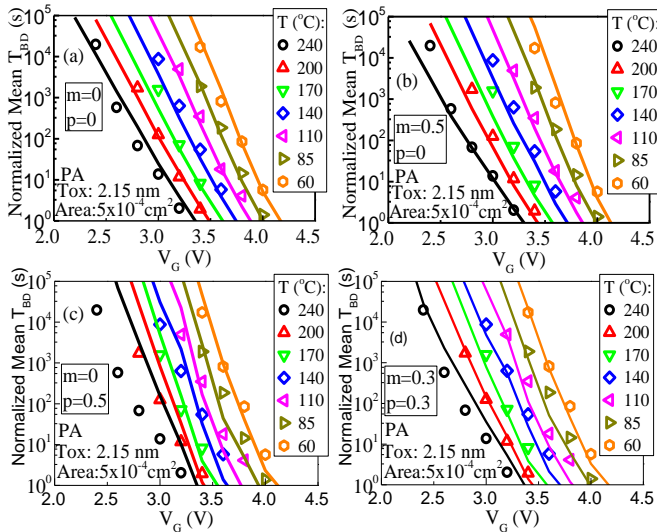


Fig.12 Mean  $T_{BD}$  versus stress  $V_G$  at different temperatures ( $T$ ) for thin PMOS capacitor stressed in accumulation (PA) with  $K_{F1}$  (see the equation at the bottom in Fig.3) taking (a)  $m=p=0$ ; (b)  $m=0.5$ ,  $p=0$ ; (c)  $m=0$ ,  $p=0.5$ ; and (d)  $m=p=0.3$ . Data (symbols) from [13].

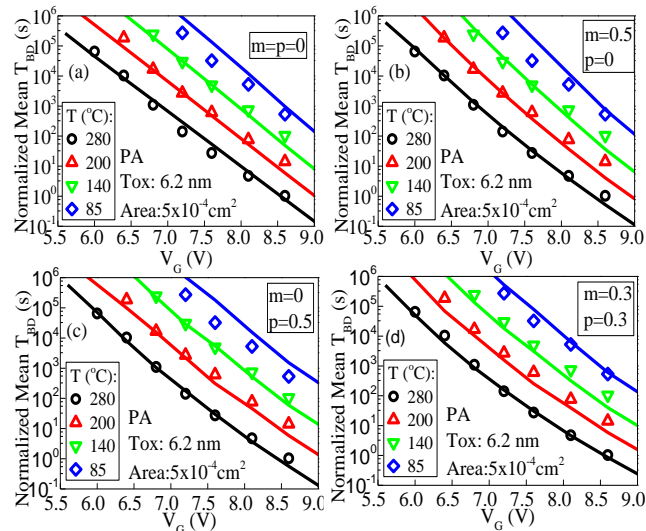


Fig.13 Mean  $T_{BD}$  versus  $V_G$  at various  $T$  for thick PMOS capacitor stressed in accumulation (PA) with  $K_{F1}$  (see the equation at the bottom in Fig.3) taking (a)  $m=p=0$ ; (b)  $m=0.5$ ,  $p=0$ ; (c)  $m=0$ ,  $p=0.5$ ; and (d)  $m=p=0.3$ . Data (symbols) from [10].

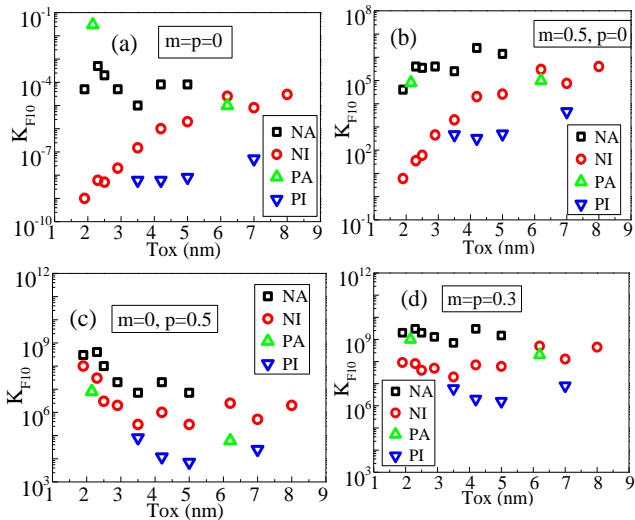


Fig.14  $K_{F10}$  variation with oxide thickness ( $Tox$ ) under various bias conditions with  $K_{F1}$  for (a)  $m=p=0$  (b)  $m=0.5$ ,  $p=0$  (c)  $m=0$ ,  $p=0.5$  (d)  $m=p=0.3$ .

## VI. COMPARISON OF END-OF-LIFE $V_{GMAX}$

The mean  $T_{BD}$  extrapolation up to 10 years ( $\sim 3.15 \times 10^8$ s) with gate voltage ( $V_G$ ) from exponential-1/E law, exponential-E law, and  $V_G$  power law for 1.9nm NMOS capacitor in inversion at a constant  $T$  is illustrated in Fig.15(a) while those calculated from RDD simulations for different  $m$  and  $p$  values are shown in Fig.15(b). The maximum gate voltage ( $V_{GMAX}$ ), corresponding to a mean  $T_{BD}$  of 10 years, versus  $Tox$  is shown in Fig.16(a) and Fig.16(b) respectively from extrapolation and RDD simulation. The exponential-1/E law projects the most optimistic  $V_{GMAX}$ , like AHI model ( $m=0$ ,  $p=0.5$ ), for all  $Tox$ 's whereas exponential-E law turns out to be the most conservative. The TC model ( $m=p=0$ ) becomes less optimistic at higher  $Tox$ 's whereas AHR model ( $m=0.5$ ,  $p=0$ ) does so at lower  $Tox$ 's. For all  $Tox$ 's studied, the  $m=p=0.3$  case shows a lesser optimistic but a consistent  $Tox$  dependence of  $V_{GMAX}$  values bounded by other cases. The variation of  $K_{F30}$  with  $Tox$  is shown in Fig.17 and the RDD parameters used are listed in Table-I. The tunneling electron leakage currents from  $Q_{BD}$  data [6], used for the RDD simulations, along with the leakage model [11] is shown in Fig.18.

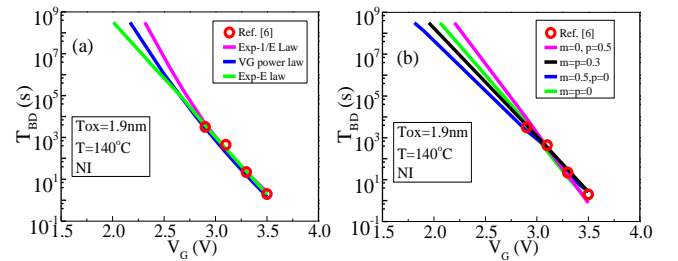


Fig.15 Illustration of gate bias extrapolation for 1.9nm up to 10 years ( $\sim 3 \times 10^8$ s) of mean  $T_{BD}$  from (a) exponential 1/E-law, exponential E-law, and  $V_G$  power law (b) RDD model simulation with  $m=p=0$ ;  $m=0.5$ ,  $p=0$ ;  $m=0$ ,  $p=0.5$ ; and  $m=p=0.3$ . Data (symbols) from [6].

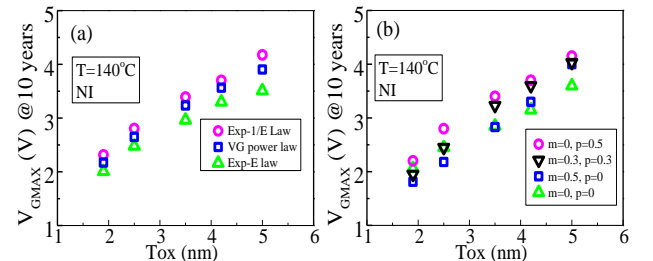


Fig.16 The maximum  $V_G$  variation with  $Tox$  for 10 years of mean  $T_{BD}$  from (a) exponential 1/E-law, exponential E-law, and  $V_G$  power law and (b) RDD simulation with  $m=p=0$ ;  $m=0.5$ ,  $p=0$ ;  $m=0$ ,  $p=0.5$ ; and  $m=p=0.3$ .

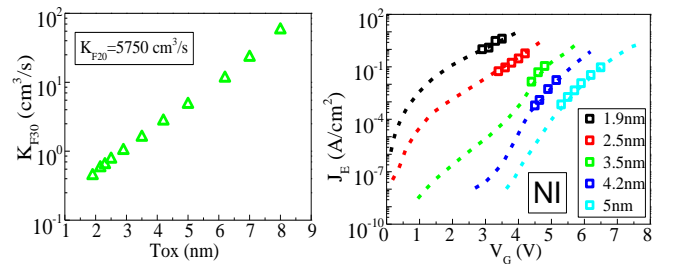


Fig.17 Variation of  $K_{F30}$  parameter of RDD model with oxide thickness  $Tox$ .

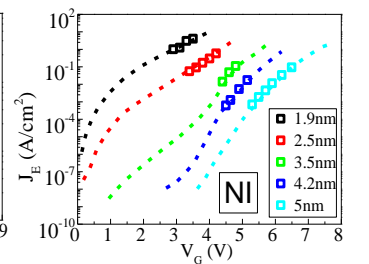


Fig.18 Tunneling leakage current variation with gate bias for various  $Tox$ . Symbols from  $Q_{BD}/T_{BD}$  data [6], dashed lines from tunneling model [11].

Table-I: RDD Model parameters

Tunneling leakage current exponent	$E_{AKF1}$ (eV)	$\Gamma_0$ (cm/MV)	$\alpha$ (qÅ)
m=p=0	0.6	2.4	1.5
m=0.5, p=0	1.0	0.6	4.8
m=0, p=0.5	0.6	1	1.0
m=p=0.3	0.8	0.1	3.6

## VII. CONCLUSION

RDD model is consistent with the requirement of reduction in  $\Delta N_{TG}$  time slope  $n$  at lower  $T_{ox}$ , in order to model Weibull  $\beta$  variation with  $T_{ox}$  and SILC measurements. The  $V_G$  dependence of mean  $T_{BD}$  is successfully modelled across  $T_{ox}$  and  $T$ , for different stress cases (NI, NA, PI, and PA). The co-injection of electrons and holes (from AHI), together with  $E_{OX}$  and  $T$  are found to be responsible for the initial trigger for the RDD model, leading to consistent set of model parameters across different  $T_{ox}$ .

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