

# A Physical Model of the Abnormal Behavior of Hydrogen-Terminated Diamond MESFET

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**Abstract** — Hydrogen-terminated Diamond MESFET exhibits abnormal gate capacitance-voltage (CV) relationship, with an insulator barrier-like plateau followed by an abrupt increase in capacitance at large gate bias. A new model is proposed by assuming inhomogeneous band gap in the Interfacial Layer (IL). With this model, various experimental gate CV, as well as gate leakage curves can be matched well with TCAD simulations. This model also offers possible explanation of the double-bumps in transconductance measurements in some experiments.

**Keywords**—Hydrogen-terminated Diamond, MESFET, CV, TCAD, Simulation, Interfacial Layer, RF, Power

## I. INTRODUCTION

Diamond is an ideal material for power and radio frequency applications due to its wide band gap (5.5eV), high hole mobility (1200cm<sup>2</sup>/Vs) and high thermal conductivity (22W/cm·K) [1]. However, activation of dopants in diamond is generally difficult due to their large ionization energies (e.g. B has activation energy of 0.36eV) [1]. Surface Hydrogen-terminated diamond (H-Diamond) provides an alternative, in which high areal density ( $p_s = 10^{13}$  to  $10^{14}$ cm<sup>-2</sup>) of Two-Dimensional Hole Gas (2-DHG) can be induced at the surface, despite the reduction of mobility to  $\sim < 100$ cm<sup>2</sup>/V·s [2][3]. High performance H-terminated MESFETs have been demonstrated [3].

However, H-terminated diamond MESFETs sometimes exhibit abnormal behaviors that are not well understood. Firstly, in capacitance-voltage (CV) measurements, the Schottky gate contact exhibits an insulator barrier-like plateau followed by an abrupt increase in capacitance at large gate bias [4][5]. Secondly, there are double-bumps in transconductance ( $g_m$ ) measurements [2][3], which are only possible in dual channel devices. Kasu et al. tried to explain it as the hole penetration through the Interfacial Layer (IL) barrier [6]. However, this is not compatible with the general expectation, that gate leakage will be much higher than the experimental values. Zhou et al. were able to fit experimental CV with their new model but it appears that

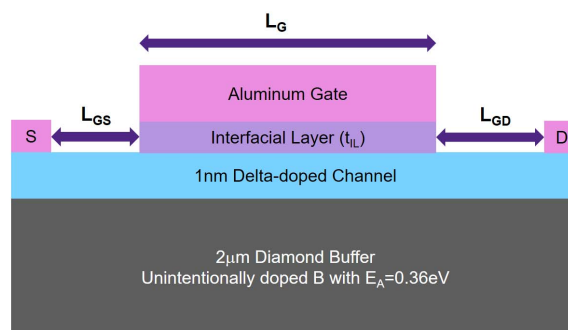


Fig.1. Structure used in the simulations. Experimental values of buffer doping, IL thickness  $t_{IL}$ ,  $L_G$ ,  $L_{GS}$  and  $L_{GD}$  are used in the corresponding simulations.

structure (MISFET) different from the experiment (MESFET) was used in the simulation [7].

Based on the ideas proposed in [6] and [7], in this paper, we present a new physical model to explain the phenomena and demonstrate good matching to experimental Schottky gate CV and leakage (AC and DC), from two different research groups, in the regimes of interest, through TCAD simulations [8].

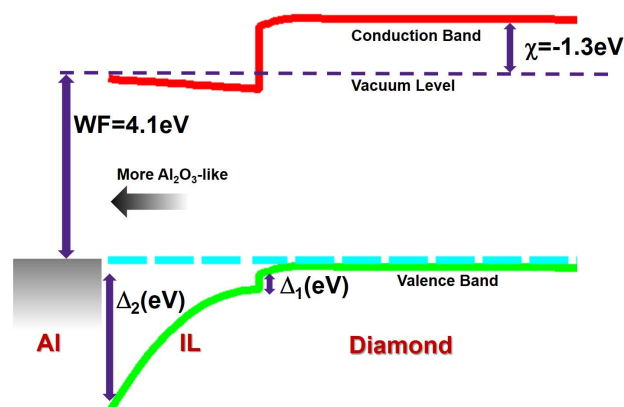


Fig. 2. Band diagram cutting vertically under the Al gate at zero bias, illustrating the setup of the new model. H-diamond with Negative Electron Affinity = -1.3eV is used [1].

## II. MODELING

Fig. 1 shows a typical H-terminated diamond MESFET structure. It is believed that an interfacial layer (IL) is generally formed between the Schottky gate metal (Aluminum) and H-Diamond in a MESFET [2][4][5]. Such layer is about several nm's ( $t_{IL}$ ) in thickness as found in TEM pictures [5]. Therefore, Schottky contacts in H-Diamond should be modeled with an IL as shown in Fig. 1.

However, with a homogenous IL, one cannot explain the two phenomena mentioned. We therefore postulate that the interfacial layer is composed of varying effective bandgap and thus varying valence band (VB) offset to the VB of diamond (Fig. 2). The VB offset increases from IL/diamond interface to Al/IL interface. This postulation is justified by the facts that 1) thin aqueous wetting layer is believed to exist at the H-Diamond surface to induce 2DHG [1][9] and 2) SIMS data show high Al and O concentration at the Al/IL interface [5]. So, we further postulate that near the IL/diamond interface, the IL "VB" is due to the Highest Occupied Molecular Orbital (HOMO) in the wetting layer [1] and near Al/IL interface, the IL "VB" has strong characteristics of the VB of  $Al_2O_3$ . Therefore, the bandgap of IL is chosen to increase quadratically from  $\sim 4.5eV$  at the IL/diamond interface to that of  $Al_2O_3$  ( $\sim 7.2eV$ ) at the Al/IL interface (Fig. 2).  $E_g \sim 4.5eV$  at the IL/diamond interface was chosen because it gives the best fitting.

The main fitting parameters of this model are the VB offset between diamond and IL ( $\Delta_1$ ) and the offset between the Fermi level of Al gate and the IL VB ( $\Delta_2$ ).

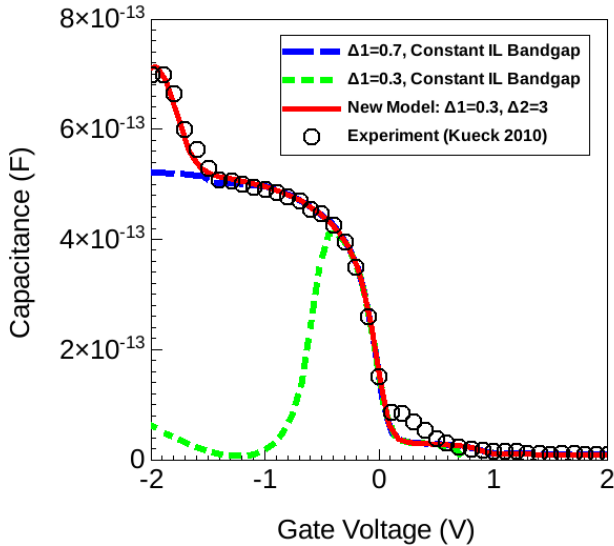


Fig. 3. Experimental [4] and simulated CV's with  $p_s=1.1 \times 10^{13} cm^{-2}$ ,  $t_{IL}=3.9nm$ ,  $L_G=0.8\mu m$  and  $W_G=50\mu m$ . Measurement frequency is 1MHz.

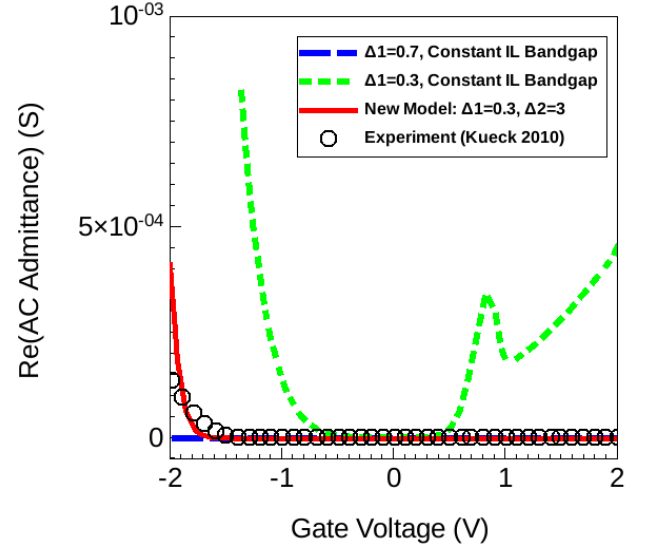


Fig.4. Experimental [4] and simulated real part of AC admittance measurement under the same condition as in Fig. 3.

## III. SIMULATION SETUP

In the simulations, Al gate has work function (WF) of  $4.1eV$ . The 2DHG is induced by delta doping at the IL/Diamond interface to give areal concentration of the values given by the corresponding experimental data. In order to match the experimental threshold voltage, fixed charge ( $Q_F$ ) is added at the IL/Diamond interface under the Al gate. Since the IL/Diamond under and outside Al gate underwent different processing steps, it is reasonable to assume they have different amount of fixed charge.

The diamond buffers are boron doped with concentration based on the values given in the experiment with activation energy =  $0.36eV$  and is simulated with incomplete ionization model.

Fermi-Dirac statistics and mobility high field saturation models calibrated to experimental data were used [10]. Poisson equation, electron and hole continuity equations and hole tunneling (using WKB) at gate are solved self-consistently. Hole tunneling mass =  $0.8m_0$  was used in all simulations, where  $m_0$  is the free electron mass.

## IV. SIMULATION RESULTS

We firstly simulated the MESFET reported by Kueck et al. which has  $L_G = 0.8\mu m$  and  $W_G = 50\mu m$  [4]. Fig. 3 shows the experimental and simulated CV curves, with 2DHG areal density  $p_s = 1.1 \times 10^{13} cm^{-2}$ , IL dielectric thickness  $t_{IL} = 3.9nm$  and measurement frequency  $f = 1MHz$ .  $p_s$  and  $t_{IL}$  values were based on the data given in [4]. IL dielectric constant is found to be 6.5 to match the

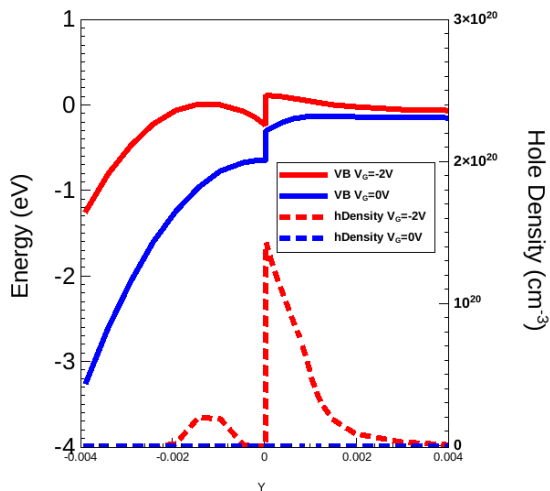


Fig.5. Valence band edge profiles and hole carrier density distributions at  $V_G=0V$  or  $V_G=-2V$  using the new model. IL/Diamond interface is at  $Y=0$ .

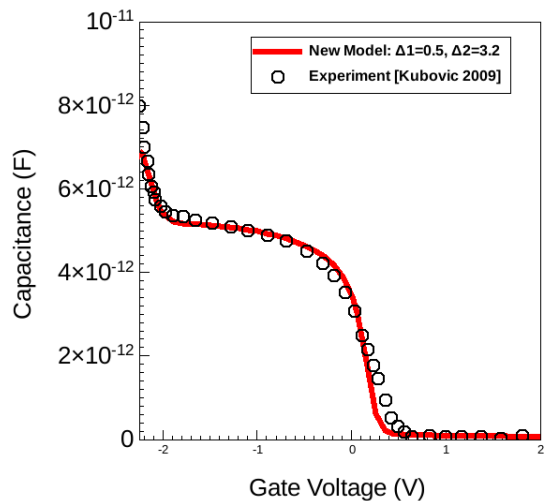


Fig. 7. Experimental [5] and simulated CV curves with  $p_s=1.1 \times 10^{13} \text{cm}^{-2}$ ,  $t_{IL}=5 \text{nm}$ ,  $L_G=5 \mu\text{m}$  and  $W_G=100 \mu\text{m}$ . Measurement frequency is 1MHz.

experimental capacitance value.  $Q_F = 1.35 \times 10^{13} \text{cm}^{-2}$  was used to match the threshold voltage. With the proposed model, the plateau and the abrupt increase at  $V_G \sim -1.5V$  can be properly captured with  $\Delta_1 = 0.3 \text{eV}$  and  $\Delta_2 = 3 \text{eV}$ . On the other hand, if homogenous IL is assumed, either the abrupt increase was not captured ( $\Delta_1 = 0.7 \text{eV}$  case) or the leakage is too large and gives wrong C-value (which is expected to also appear in experiment when leakage is large) at high negative gate bias ( $\Delta_1 = 0.3 \text{eV}$  case).

Fig. 4 shows the real part of simulated and experimental [4] admittance in a CV measurement, which is a strong function of gate leakage. The new model properly captures the abrupt increase of the AC conductance at  $V_G \sim -1.7V$  as in experiment. However, for constant IL bandgap, when  $\Delta_1 = 0.3 \text{eV}$  is used, the

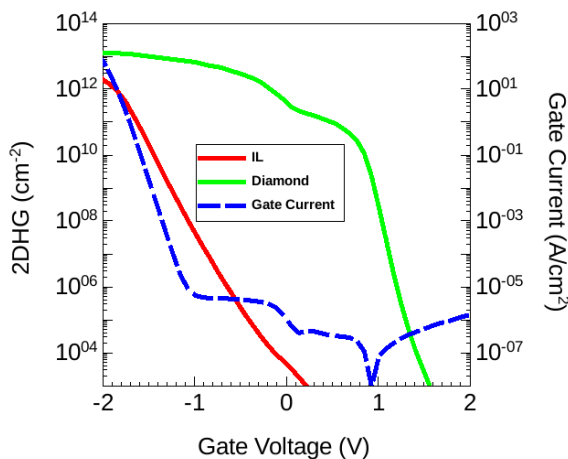


Fig. 6. Simulated gate leakage current and 2DHG in IL and Diamond as a function of gate voltage of the structure in [4].

leakage current is too large resulting in erroneous AC admittance.

Fig. 5 shows the VB profiles and hole distributions at  $V_G = 0V$  and  $-2V$ . It is evident that holes accumulate inside IL at high bias due to the inhomogeneity of IL bandgap, and a second conduction channel is formed, which is the origin of the abrupt increase in capacitance in CV measurements. From the band diagram, it can be seen that the formation of the second conduction channel is due to the quadratic increase in the bandgap of IL. Moreover, if  $\Delta_1$  and  $\Delta_2$  are too small, the gate leakage will be excessive. When  $\Delta_1$  is too large, there will not be enough hole accumulating in the IL second channel and thus there will be no abrupt increase in capacitance.

Fig. 6 shows the simulated gate leakage and the 2DHG areal densities in diamond and IL as a function of gate voltage. It can be seen that hole starts populating in IL after  $V_G = -1.2V$  and reaches almost  $2 \times 10^{12} \text{cm}^{-2}$  at  $V_G = -2V$ .

## V. ADDITIONAL VERIFICATIONS

By using the same model, experimental data [5] from another research group was fitted, in which  $p_s = 1.1 \times 10^{13} \text{cm}^{-2}$ ,  $t_{IL} = 5 \text{nm}$ ,  $L_G = 5 \mu\text{m}$ ,  $W_G = 100 \mu\text{m}$  and  $f = 1 \text{MHz}$ . IL dielectric constant = 6.5 was used which gives the best fitting. This is different from the extracted value (5.7) in the original paper because of different extraction methodologies.  $Q_F = 1.1 \times 10^{13} \text{cm}^{-2}$  was used to match the threshold voltage. Fig. 7 shows the experimental [5] and simulated CV. By adjusting only  $\Delta_1$  to 0.5eV and  $\Delta_2$  to 3.2eV, the simulated result matches not only the CV but also the DC gate leakage well (Fig. 8) in the regime of

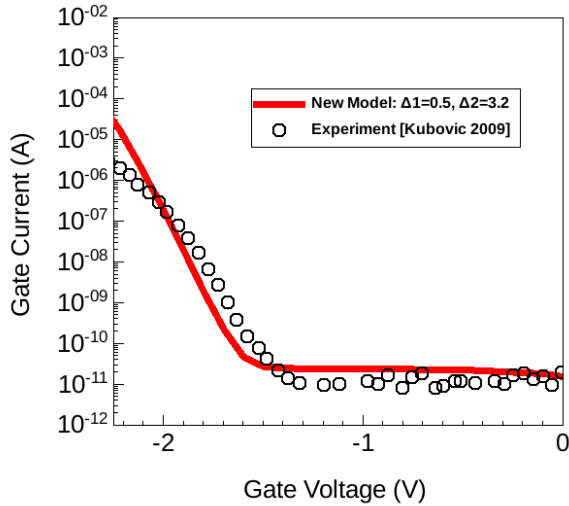


Fig. 8. Experimental [5] and simulated DC gate leakage using the same structure in Fig. 7

interest. Since IL is unintentionally formed without well-controlled conditions, it is reasonable to have different  $\Delta_1$  and  $\Delta_2$  in different fabrication processes.

It is important to verify that the proposed model has no unphysical effect to other type of simulations. By including the proposed model, the  $I_D$ - $V_G$  curve of H-Diamond MESFET from yet another research group [2] is fitted successfully as showed in Fig. 9.

The device has  $L_G=0.5\mu\text{m}$ ,  $L_{GS}=1\mu\text{m}$ ,  $L_{DS}=1.5\mu\text{m}$ ,  $W=50\mu\text{m}$  and  $V_D=-6\text{V}$ .  $\Delta_1$  and  $\Delta_2$  are 0.3eV and 3eV respectively.  $Q_F = 0.8 \times 10^{13}\text{cm}^{-2}$  was used to match the threshold voltage. Low hole field mobility is  $80\text{cm}^2/\text{V}\cdot\text{s}$  with saturation velocity =  $1.1 \times 10^7\text{cm/s}$ .

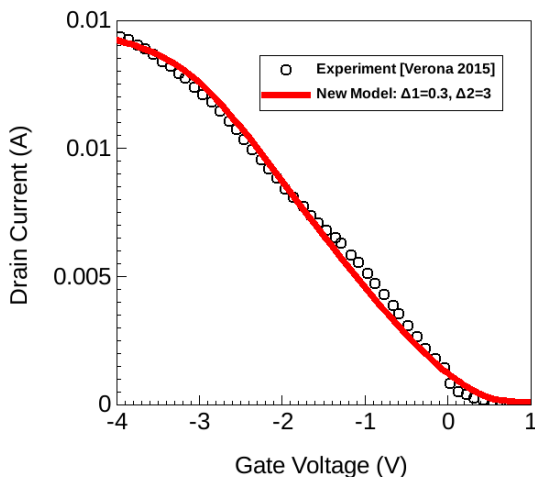


Fig. 9. Experimental [2] and simulated  $I_D$ - $V_G$ .  $L_G=0.5\mu\text{m}$ ,  $W=50\mu\text{m}$ ,  $L_{GS}=1\mu\text{m}$ ,  $L_{DS}=1.5\mu\text{m}$  and  $V_D=-6\text{V}$ .

## VI. CONCLUSION

A new model for H-Diamond Schottky contact is proposed and verified with TCAD simulations. It is assumed that the IL has bandgap increasing quadratically from about 4.5eV at the IL/diamond interface to about 7.2eV (that of  $\text{Al}_2\text{O}_3$ ) at the Al-gate/IL interface. The main fitting parameters are the VB offset between IL and Diamond ( $\Delta_1$ ) and the offset between the fermi level of Al gate and IL VB ( $\Delta_2$ ). By varying  $\Delta_1$  and  $\Delta_2$ , it can explain and fit experimental results of the abnormal behaviors in CV curves from several different research groups. It also predicts the formation of dual channels, which can be used to explain the double-bumps in some  $g_m$  measurements based on the speculation in [2].

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