

# NEMO5: Predicting $MoS_2$ Heterojunctions

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**Abstract**—Molybdenum disulfide ( $MoS_2$ ) is a promising 2D material since it has a finite band gap, and its electronic band structure depends on the layer thickness. The tunability of the gate voltage on band alignment of different  $MoS_2$  layers is analyzed. For this purpose, the multipurpose nanodevice simulation tool NEMO5 was altered by several new features: electronic bandstructure calculations in maximally localized Wannier function (MLWF) representation and self-consistent charge calculations with subatomic electrostatic resolution.

## I. INTRODUCTION

$MoS_2$  is a promising material candidate for next generation nano-devices due to few-atom device thickness and best imaginable gate control. In contrast to e.g. graphene,  $MoS_2$  offers a finite band gap which makes it suitable for transistors [1] to have a smaller off-current.

It is also shown to have a wide range of material characteristics depending on the layer thickness. The band-gap decrease from 1.8 eV to 0.9 eV [2] from bulk to mono-layer(ML) and is indirect for two-layer and the thicker films while direct for a single layer  $MoS_2$ . A geometry induced heterojunction [3] can therefore be formed with a discontinuity of layer thickness. The wide range of band structure characteristics provided are interesting for thin and flexible optoelectronics e.g. photodetectors[4] [5], LED [6]. To contribute to both applications, the band offset variation with different layer thicknesses and its gate field dependence is assessed here. With the calculation of electron affinity, the junction type can be determined and predict the transport across the junction.

To represent the hamiltonian of the system, empirical tight binding model(ETB) [7] and density functional theory(DFT) [8] are available. The intensive computation requirement for DFT and non-transferability of van der Waals induced interlayer coupling with ETB have lowered the application to device simulation and prediction. In this paper, Wannier functions are used and transferability is shown for different layer of Van der Waals layered structures.

### A. Methodology

$MoS_2$  with trigonal prismatic polytype, known as "2H" symmetry, in the bulk form is used. The electronic Hamiltonian and basis functions of density functional theory simulations (GGA with PBE [9]) are calculated with the VASP software [10]. The structure is relaxed with a convergence

criterion of  $10^{-8}$ eV in between iterations of a self-consistent electronic calculation.

With the Wannier90 software[11], the electronic Hamiltonian of a  $MoS_2$  unit cell in MLWF representation is transformed from the original basis. The initial projection basis includes d orbitals with Molybdenum and p orbitals for Sulfur atoms. The spreading of the Wannier function, defined in Eq. 1, is reduced iteratively and  $2 \text{ \AA}^2$  achieved in the end.

$$\Omega = [\langle r^2 \rangle - \bar{r}^2] \quad (1)$$

This unit cell Hamiltonian is repeated in the NEMO5 software [12] to generate a real-size device Hamiltonian and solve for electronic densities. For charge self-consistency, the Schrödinger equation is iteratively solved with the Poisson equation in sub-atomic resolution, simplifying the spatial shape of Wannier functions with Gaussian functions with a standard deviation of 0.14 nm( square root of the  $\Omega$  in Eq. 1). All assessed structures are close to the experimental setups of Ref. [3] with a consistently scaled oxide and gate voltage.

A metal oxide semiconductor structure is set up as shown in Fig. 1. The equivalent oxide thickness is set to be 12.36 nm. The channel consists of one to ten layers of  $MoS_2$  with donor doping  $1.5 \times 10^{18} \text{ cm}^{-3}$ . On the gate side, a Dirichlet boundary condition is used to set the gate voltage value. The top side is exposed to vacuum and a Neumann boundary condition of flat band (zero electric field) is applied. Bandstructure is calculated and the response to different gate voltages, doping and layer thickness are analyzed.

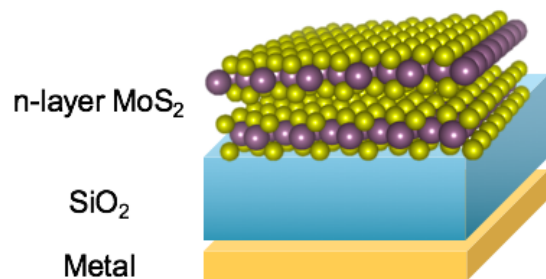


Fig. 1. Schematics of the metal-oxide semiconductor structure simulated in this paper. Gate is on the bottom and vacuum on the top with varying thickness of  $MoS_2$  layer.

### B. Discussion

After a structure relaxation, a lattice constant with  $a = 3.18 \text{ \AA}$  and  $c = 12.48 \text{ \AA}$  are used [13]. The MLWF-calculated bandstructures are found to agree very well with ab-initio calculations, Fig.2 and Fig.3, for any  $MoS_2$  layer thickness with an interaction radius of 1.2 nm. The conduction band minimum in 2D momentum space is shown in Fig.4. Two inequivalent valleys exist with non-isotropic and layer dependent effective masses (in agreement with Ref.[14]). The K valley has an average DOS effective mass of  $0.62 m_0$  increasing with the layer thickness, while the  $\Lambda$  valley is about  $0.52 m_0$  decreasing with the number of layers.

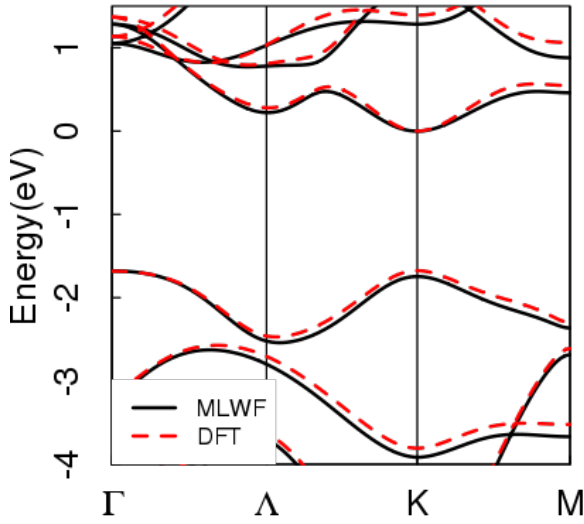


Fig. 2. Comparison of mono-layer  $MoS_2$  band diagrams solved with MLWF in NEMO5 and the DFT functionality of VASP. The agreement confirms the transferability of the MLWF parameters.

The conduction band edge relative to the Fermi level as a function of numbers of layers and gate bias is given in Fig.5. Both the layer thickness and the applied bias show a significant impact on the conduction band edge. The energy difference of the K and  $\Lambda$  valleys is shown in Fig. 6. The crossover of the conduction band minimum from K to  $\Lambda$  valley illustrated in this figure is the reason for the minimum of conduction band edge in Fig.5 .

Figure 7 compares the band structure of a 5-layer thick  $MoS_2$  system with and without a finite gate bias applied. Without electric fields, the K-valley is 5-fold degenerate. The finite gate field lifts this degeneracy and reduces the energy separation between K and  $\Lambda$  valleys. The wave function is spatially resolved for the three lowest eigenstates at these two valleys in Fig. 8 and Fig.9. The  $\Lambda$  valley states are shown to have a strong interlayer coupling and extend across the layers. In contrast, K valley wavefunctions are very localized.

### II. CONCLUSION

The nanodevice simulation tool NEMO5 was altered to accurately describe  $MoS_2$  electrons in MLWF representation.

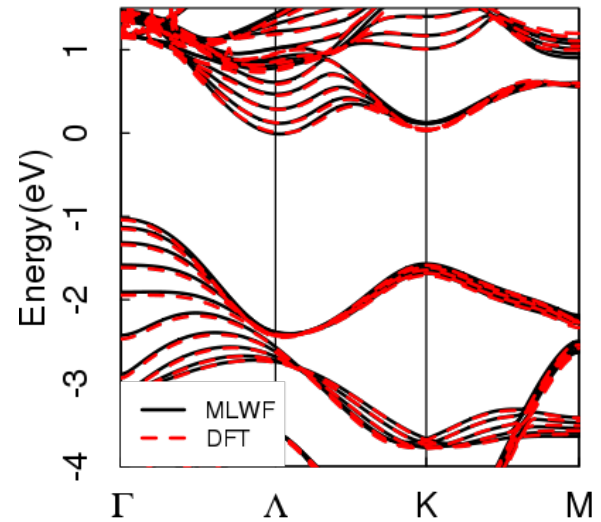


Fig. 3. Comparison of five-layer  $MoS_2$  band diagrams solved with MLWF in NEMO5 and the DFT functionality of VASP. The agreement confirms the transferability of the MLWF parameters.

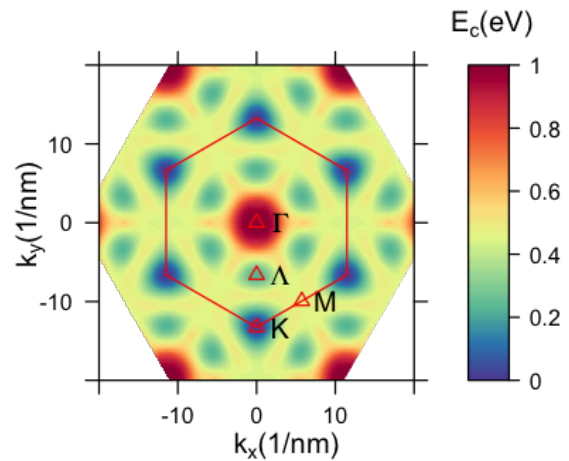


Fig. 4. Contour graph of the minimum of the conduction band for mono-layer  $MoS_2$ . The red hexagon depicts the first Brillouin zone.

All MLWF Hamiltonians are verified against ab-initio models. Hetero-junctions of different layers show band misalignment as a function of gate voltages. Distinct valleys are tunable with layer thickness and applied gate voltage.

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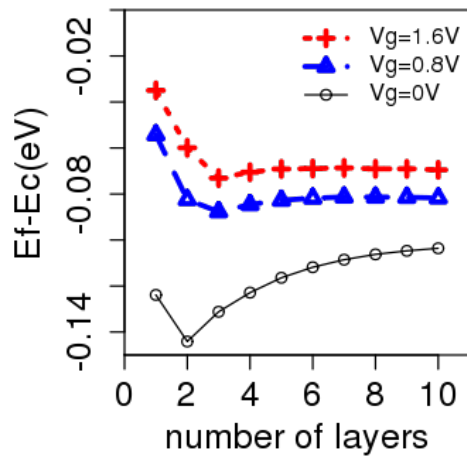


Fig. 5. Conduction band minima of  $MoS_2$  structures with different layer thicknesses as shown in Fig. 1 and different gate bias voltages.

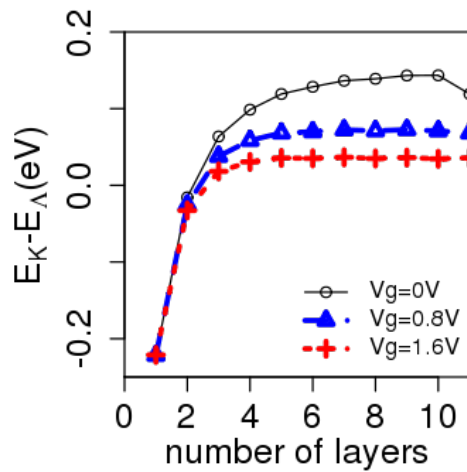


Fig. 6. Energy difference between the K and  $\Lambda$  valleys of a  $MoS_2$  structure with different number of layers as shown in Fig. 1 and for different gate bias voltages.

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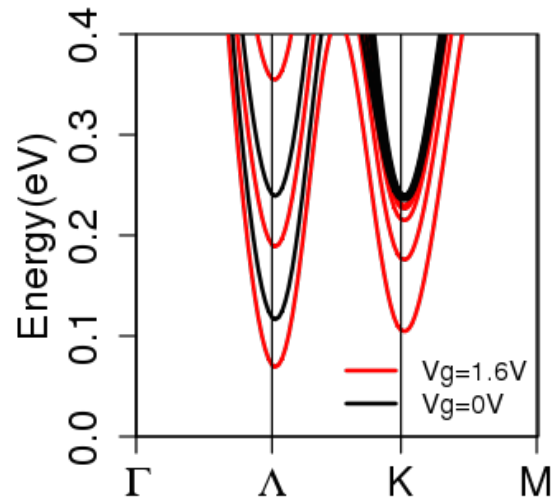


Fig. 7. Bandstructures of a 5-layer  $MoS_2$  system as shown in Fig. 1 with and without finite gate voltage. The gate field lifts the degeneracy at the K valley while it only shifts  $\Lambda$  valley energies.

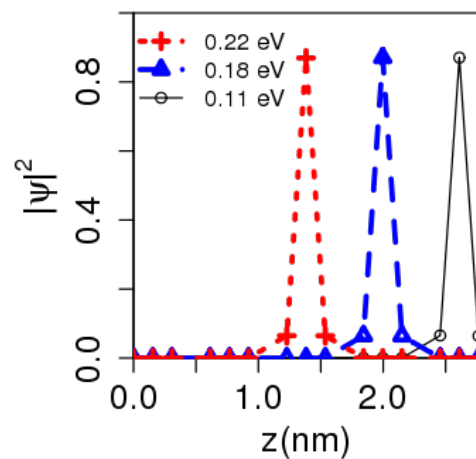


Fig. 8. Absolute squared wave function of the three lowest K valley eigenstates of the 5 layer  $MoS_2$  device of Fig. 7 at  $V_g = 1.6V$ .

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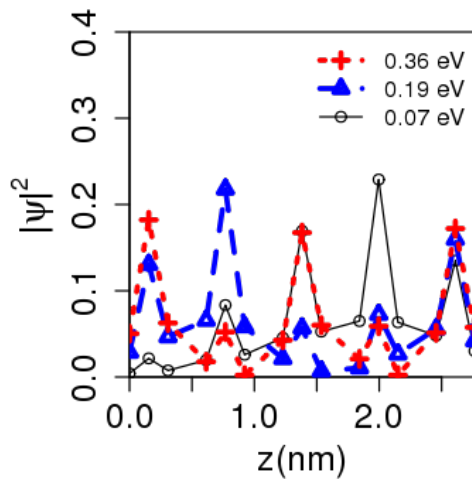


Fig. 9. Absolute squared wave function of the three lowest  $\Lambda$  valley eigenstates of the 5 layer MoS2 device of Fig. 7 at  $V_g = 1.6V$ .

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