# Monte Carlo Simulations of Electrons in $\mathrm{Al}_{4} \mathrm{SiC}_{4}$ Ternary Carbide 

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An $\mathrm{Al}_{4} \mathrm{SiC}_{4}$ ternary carbide has become a promising wide band-gap semiconductor for the semiconductor industry over the last decade because of its emerging properties [1]. A crystal structure of $\mathrm{Al}_{4} \mathrm{SiC}_{4}$ is illustrated in Fig. 1. The $\mathrm{Al}_{4} \mathrm{SiC}_{4}$ band-gap has been calculated to be $2.48 \mathrm{eV}[2,3]$ thus opening a possibility for the design of carbide heterostructure devices in a combination with $4 \mathrm{H}-\mathrm{SiC}$ or $3 \mathrm{C}-\mathrm{SiC}$. These heterostructure carbide devices could potentially resolve issues with the large interface density of states at the semiconductor interface with a dielectric layer leading to a low inversion carrier mobility in SiC MOSFETs [4]. Other remarkable properties include superior oxidation resistance [5], superior wear resistance, low weight, high strength, and high thermal conductivity [6].

In this work, an ensemble Monte Carlo (MC) simulation code is developed to investigate the electron transport in bulk $\mathrm{Al}_{4} \mathrm{SiC}_{4} . \mathrm{Al}_{4} \mathrm{SiC}_{4}$ has a wurzite lattice $[2,3]$ as shown in Fig. 2. We assume that the two lowest valleys will play a role in electron transport. The $M$-valley has also six locations contributing one-half (a total of 3 equivalent valleys). The $K$-valley has six locations contributing one-third to the 1 st Brillouin zone (a total of 2) as shown in Figs. 3 ad 4. Therefore, a two-valley nonparabolic anisotropic bandstructure model is employed with the $M$-valley to be a minimum and the second $K$-valley to be 0.52 eV above as illustrated in Fig. 5. The electron interactions with polar and non-polar phonons within and between $M$ and $K$-valleys are listed in Table 1. The material parameters in Table 2 use a mix of experimental and theoretical sources like optical phonon energies extracted from IR/Raman spectroscopy [3].

| Valley | Transition | Scattering Type |
| :---: | :---: | :---: |
| $M_{1}$ | $M_{1} \rightarrow M_{1}$ | Intra Polar |
|  | $M_{1} \rightarrow M_{2,3}$ | Inter Non-Polar |
|  | $M_{1} \rightarrow K$ | Inter Non-Polar |
| $M_{2}$ | $M_{2} \rightarrow M_{2}$ | Intra Polar |
|  | $M_{2} \rightarrow M_{1,3}$ | Inter Non-Polar |
|  | $M_{2} \rightarrow K$ | Inter Non-Polar |
| $M_{3}$ | $M_{3} \rightarrow M_{3}$ | Intra Polar |
|  | $M_{3} \rightarrow M_{1,2}$ | Inter Non-Polar |
|  | $M_{3} \rightarrow K$ | Inter Non-Polar |
| $K$ | $K \rightarrow M_{1}$ | Inter Non-Polar |
|  | $K \rightarrow M_{2}$ | Inter Non-Polar |
|  | $K \rightarrow M_{3}$ | Inter Non-Polar |

Table 1: Electron-phonon scattering transitions considered in the MC model.

Finally, $M$-valley k-vector (inverse) transformations to a spherical space (denoted by ${ }^{*}$ ) within the anisotropic analyt-

Table 2: $\mathrm{Al}_{4} \mathrm{SiC}_{4}$ material parameters considered in the MC simulations.

| Parameter [Unit] | Value |
| :--- | :---: |
| Mass Density [g/cm ${ }^{3}$ ] | $3.03^{\mathrm{a}}$ |
| Lattice Const. [Å] | $3.28^{\mathrm{a}}$ |
| Piezoelectric Const. [C/m²] | $0.47^{\mathrm{a}}$ |
| Longitudinal Acoustic Velo. [m/s] | $10577^{\mathrm{a}}$ |
| Transverse Acoustic Velo. [m/s] | $6431^{\mathrm{a}}$ |
| Polar Opt. Phon. Energy [meV] | $67.32^{\mathrm{b}}, 107.24^{\mathrm{b}}$ |
| Non-Polar Opt. Phon. Energy [meV] | $85.55^{\mathrm{b}}$ |
| Acoustic Def. Potential [eV] | $11.4^{\mathrm{c}}$ |
| Indirect Band Gap for the $M$-valley | $\mathrm{E}_{G}^{(M)}=2.78$ |
| (M) \& the $K$-valley (K) [eV] | $\mathrm{E}_{G}^{(K)}=3.30^{\mathrm{a}}$ |
| Electron Effective Masses $\left[m_{e}\right]$ | $\mathrm{m}_{l}^{*(M)}=0.568^{\mathrm{d}}$ |
|  | $\mathrm{m}_{t}^{*(M)}=0.695^{\mathrm{d}}$ |
|  | $\mathrm{m}_{l}^{*(K)}=1.057^{\mathrm{d}}$ |
|  | $\mathrm{m}_{t}^{*(K)}=0.936^{\mathrm{d}}$ |

${ }^{\text {a }}$ Ref. [2]. ${ }^{\mathrm{b}}$ Ref. [3]. ${ }^{\mathrm{c}}$ Average taken from [7]. ${ }^{\mathrm{d}}$ Extracted value from DFT calculations [2]. $m_{e}$ is the rest mass of an electron.
ical model use a combination of Herring-Vogt and rotational transformations [8] as:

$$
\begin{gather*}
k_{x}^{*}\left(k_{x}\right)=k_{x}\left(k_{x}^{*}\right) \cos (\theta)-(+) k_{y}\left(k_{x}^{*}\right) \sin (\theta)  \tag{1}\\
k_{y}^{*}\left(k_{y}\right)=k_{y}\left(k_{y}^{*}\right) \cos (\theta)+(-) k_{x}\left(k_{y}^{*}\right) \sin (\theta)  \tag{2}\\
k_{z}^{*}\left(k_{z}\right)=k_{z}\left(k_{z}^{*}\right) \tag{3}
\end{gather*}
$$

The MC simulations in Figs. 6 and 7 predict that $\mathrm{Al}_{4} \mathrm{SiC}_{4}$ will have a maximum electron drift velocity of $1.35 \times 10^{7} \mathrm{cms}^{-1}$ at an electric field of $1400 \mathrm{kVcm}^{-1}$ and a maximum electron mobility of $82.9 \mathrm{~cm}^{2} \mathrm{~V}^{-1} \mathrm{~s}^{-1}$. Fig. 8 shows the electron mobility dependence on ionized impurity concentration. The average electron kinetic energy and valley occupation are plotted in Figs. 9 and 10, respectively.

## References

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Fig. 1: A crystal structure of Fig. 2: A schematic of hexagonal bandstruc- Fig. 3: The hexagonal (0001) Fig. 4: Detail of locations of the $M$ $\mathrm{Al}_{4} \mathrm{SiC}_{4}$. The blue, yellow, and ture of $\mathrm{Al}_{4} \mathrm{SiC}_{4}$ in the $\mathbf{k}$-space showing a lo- $\mathbf{k}_{\mathbf{x}}-\mathbf{k}_{\mathbf{y}}$ plane of $\mathrm{Al}_{4} \mathrm{SiC}_{4}$ showing a black spheres represent $\mathrm{Al}, \mathrm{Si}$, and cation of principal valleys. C atoms, respectively.
$\mathbf{k}_{\mathbf{x}}-\mathbf{k}_{\mathbf{y}}$ plane of $\mathrm{Al}_{4} \mathrm{SiC}_{4}$ sho
location of principal valleys.
valleys in the ( 0001 ) plane within the $\mathrm{Al}_{4} \mathrm{SiC}_{4}$ hexagonal structure, where $\theta= \pm 60^{\circ}$ or $\pi / 3$.


Fig. 5: A schematic of conduction band minimum val- Fig. 6: Electron drift velocity as a function of applied leys for $\mathrm{Al}_{4} \mathrm{SiC}_{4}$ showing details of the number of electric field in a bulk $\mathrm{Al}_{4} \mathrm{SiC}_{4}$. The velocity obtained equivalent $M$ - and $K$ - valleys, the separation between valleys, and the electron-non-polar phonon interactions considered in the transport model.
 assuming an anisotropic (red squares) and a simpler isotropic (blue circles) bandstructure are shown.


Fig. 9: Average kinetic energy as a function of applied electric field in a bulk $\mathrm{Al}_{4} \mathrm{SiC}_{4}$.


Fig. 7: Electron mobility as a function of applied elecric field in a bulk $\mathrm{Al}_{4} \mathrm{SiC}_{4}$. The mobility obtained assuming an anisotropic (red squares) and a simpler isotropic (blue circles) bandstructure are plotted.


Fig. 8: Electron mobility as a function of ionized impurity concentration in a bulk $\mathrm{Al}_{4} \mathrm{SiC}_{4}$.


Fig. 10: Valley occupancy of electrons in the $M$ - and $K$-valleys vs. applied electric field in bulk $\mathrm{Al}_{4} \mathrm{SiC}_{4}$.

