

Electronic and Transport Properties of Graphene

Tsuneya ANDO

1. Introduction

- Weyl's equation for neutrino
- Berry's phase and topological anomaly

2. Singular diamagnetic susceptibility

- Band-gap effect
- Spatially varying magnetic field

3. Transport properties graphene

- Singularities at Dirac point
- Long-range scatterers

4. Multi-layer graphene

- Bilayer graphene
- Hamiltonian decomposition

5. Future outlook

- Gauge fields
- Chiral edge states

6. Summary

Nara, Jun 4 (Tue) 2013



Collaborators

Takeshi NAKANISHI (AIST)

Masaki NORO (TITech)

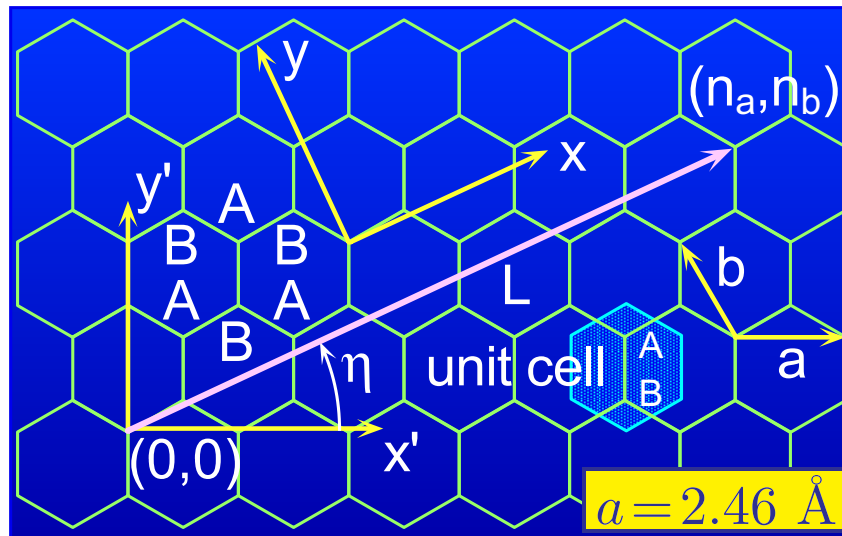
Mikito KOSHINO (Tohoku Univ)



16th International Workshop on Computational Electronics, Nara Prefectural New Public Hall
Jun 4 (Tue) – 7 (Fri) 2013 [13:30–14:30 (50+10)]

Effective-Mass Description: Neutrino or Massless Dirac Electron

Graphene (Triangular antidot lattice)



Weyl's equation for **neutrino** (K)

$$\Leftrightarrow \gamma(\boldsymbol{\sigma} \cdot \hat{\mathbf{k}}) \mathbf{F}(\mathbf{r}) = \varepsilon \mathbf{F}(\mathbf{r})$$

$$\Leftrightarrow \gamma(\sigma_x \hat{k}_x + \sigma_y \hat{k}_y) \mathbf{F}(\mathbf{r}) = \varepsilon \mathbf{F}(\mathbf{r})$$

$$\begin{pmatrix} 0 & \gamma(\hat{k}_x - i\hat{k}_y) \\ \gamma(\hat{k}_x + i\hat{k}_y) & 0 \end{pmatrix} \begin{pmatrix} F^A(\mathbf{r}) \\ F^B(\mathbf{r}) \end{pmatrix} = \varepsilon \begin{pmatrix} F^A(\mathbf{r}) \\ F^B(\mathbf{r}) \end{pmatrix}$$

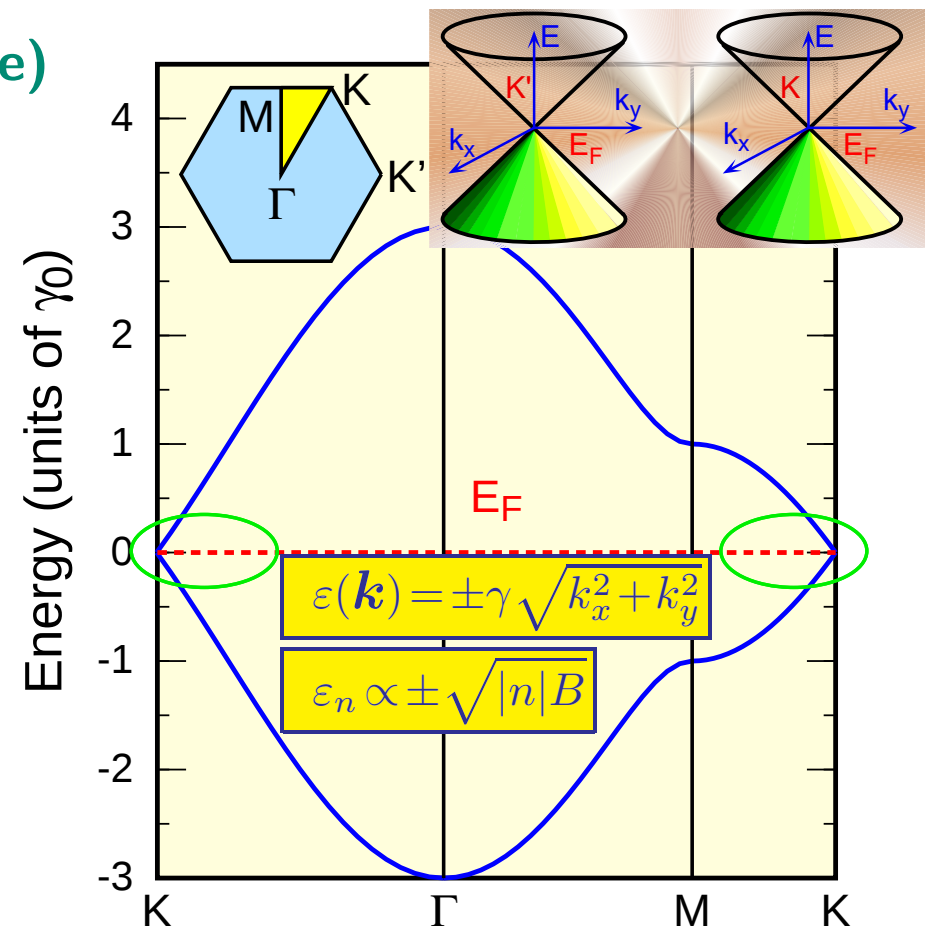
Massless (Dirac)

Constant velocity ($\sim$ light, cannot stop)

Topological anomaly

$$v_F \sim c/300 \quad (\gamma_0 \sim 3 \text{ eV})$$

$$\gamma = \sqrt{3} \gamma_0 a / 2 \quad (\gamma_0: \text{Hopping integral})$$



Wave Vector

$$\hat{\mathbf{k}} = -i \vec{\nabla}$$

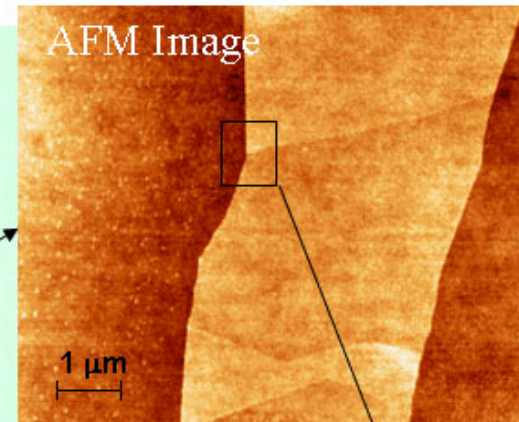
Velocity: $v_F = \gamma / \hbar$

K': $\sigma \rightarrow \sigma^*$

A.K. Geim & K.S. Novoselov
– Nobel Prize Physics 2010 –

*K.S. Novoselov et al.,
Science 306, 666 (2004)*

Mechanical exfoliation
(Repeated peeling
of graphite (HOPG)
using scotch tape)

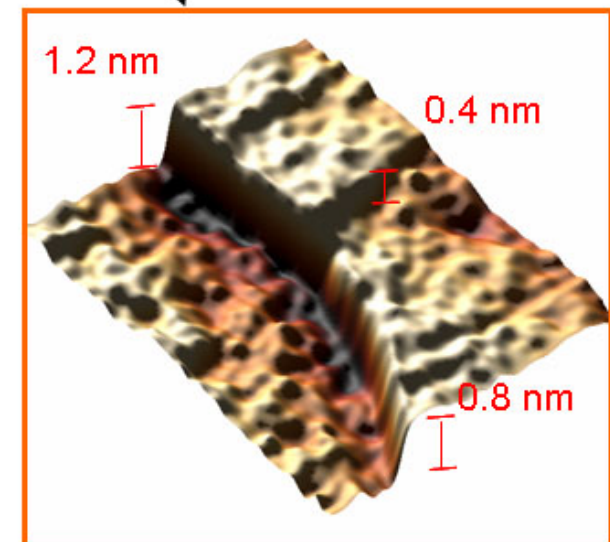
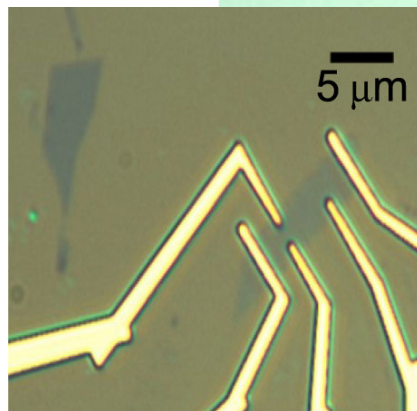


SiO₂ substrate

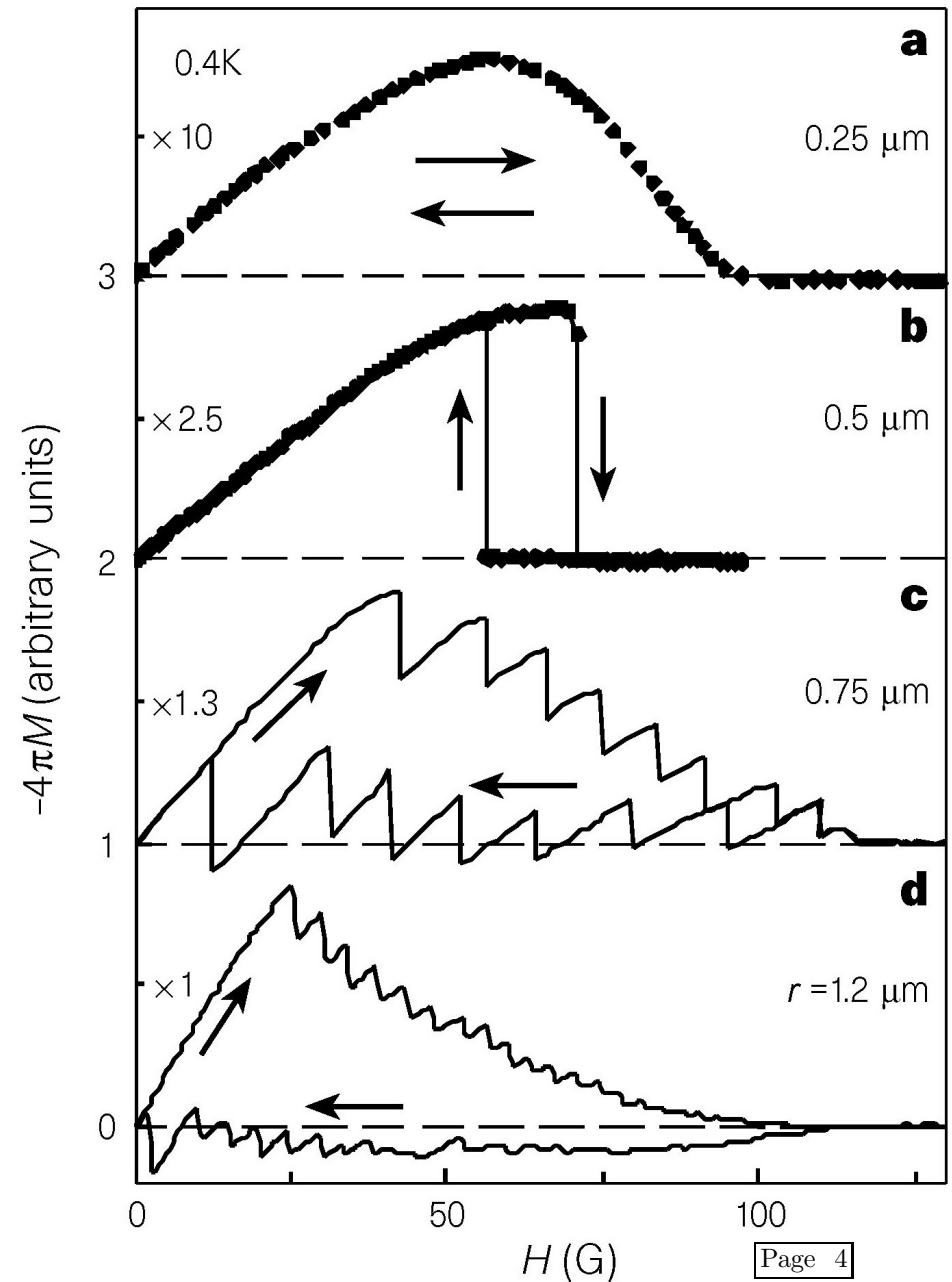
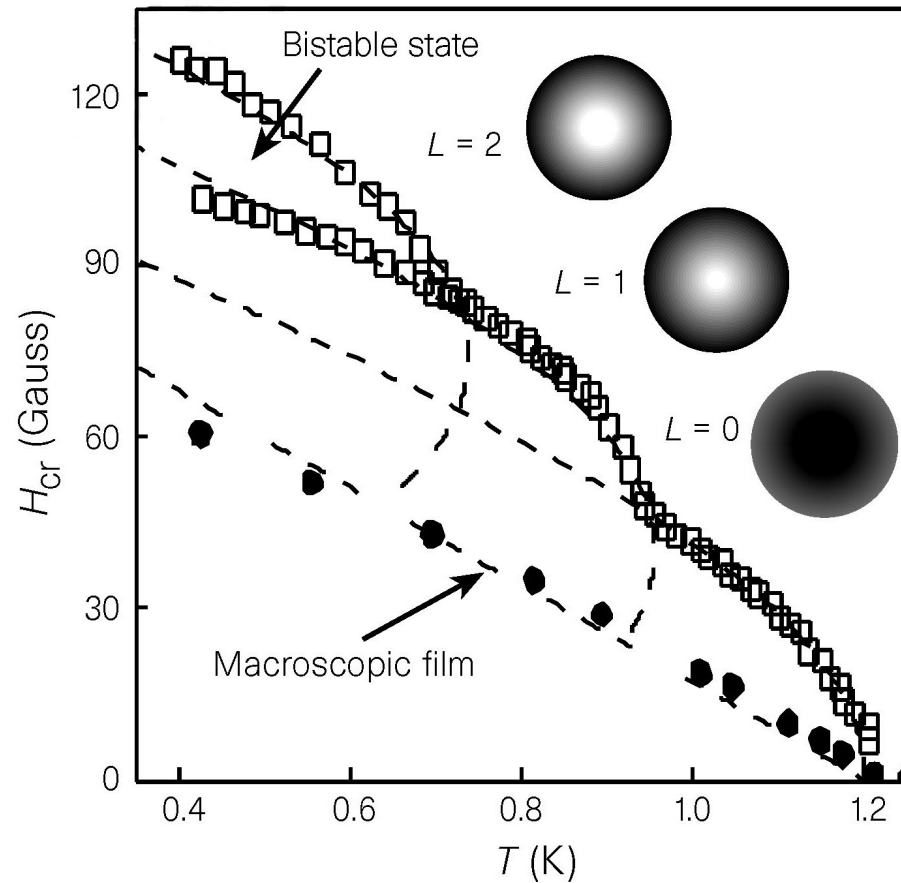
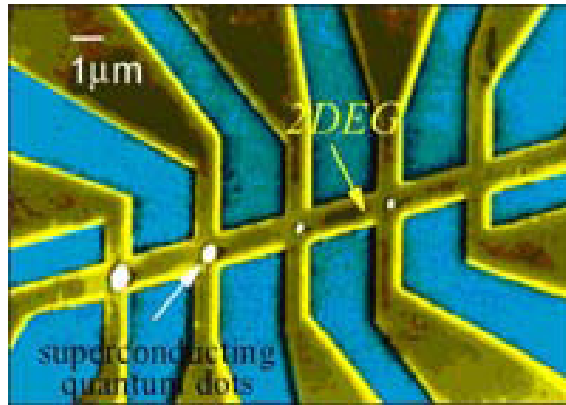
Optical microscope image

20 μm

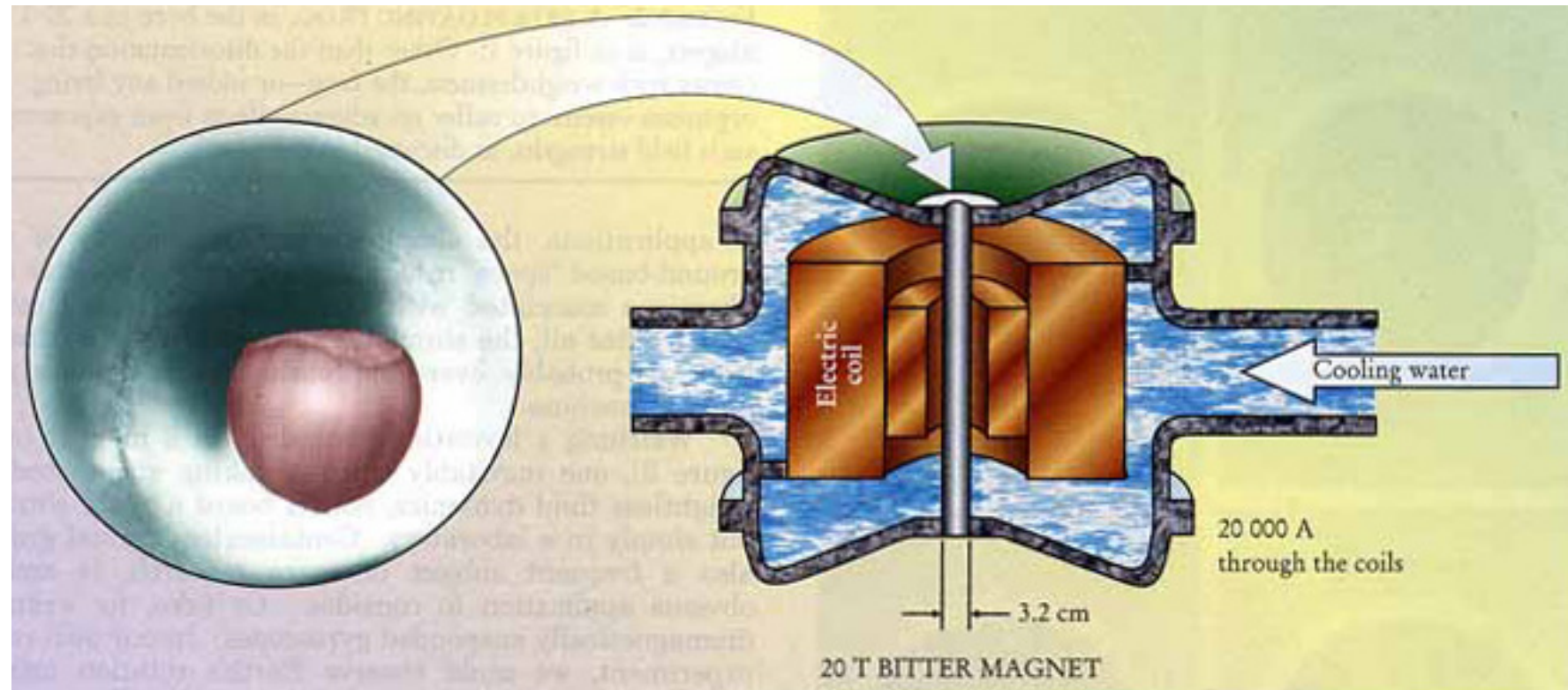
FET: Field Effect Transistor
($n_s < 5 \times 10^{13} \text{ cm}^{-2}$)



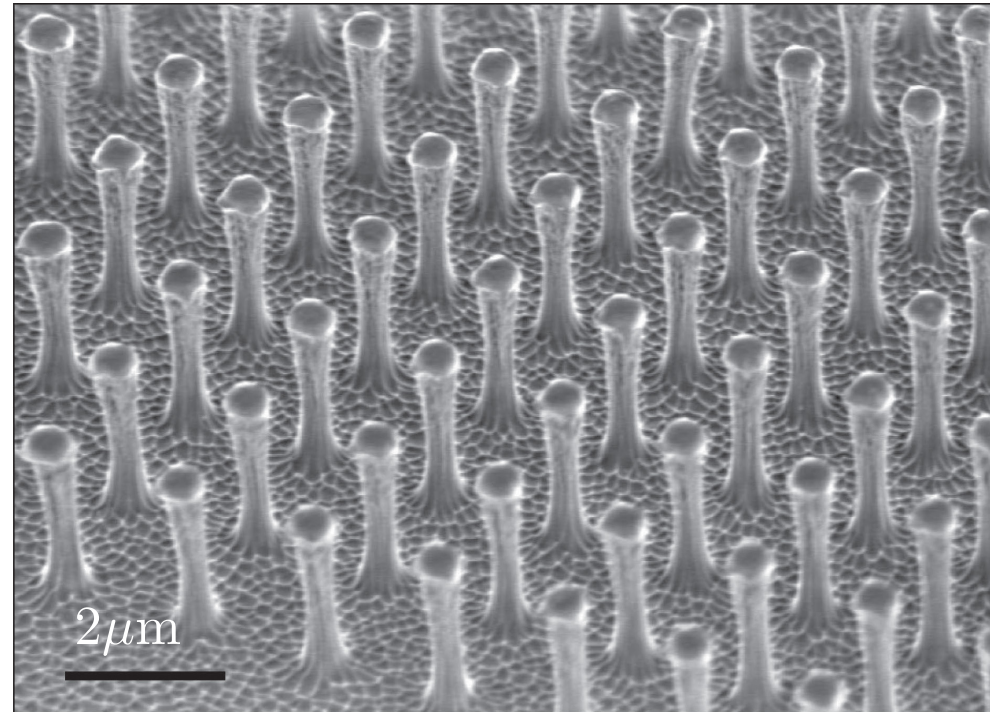
'Phase Transitions' in Small Superconductors (Nijmegen, ~1997)



Magnetic Levitation (Nijmegen, ~1998)



Gecko Tape (Manchester, ~2003)



Gecko (やもり)

Topological Anomaly and Berry's Phase

Weyl's equation : **Neutrino** $\Leftrightarrow$ **Helicity** ($\sigma \leftrightarrow k$)

$$R(\theta) = \begin{pmatrix} e^{-i\theta/2} & 0 \\ 0 & e^{i\theta/2} \end{pmatrix}$$

$$\gamma(\sigma \cdot \hat{k}) F_{s\mathbf{k}}(\mathbf{r}) = \varepsilon_s(\mathbf{k}) F_{s\mathbf{k}}(\mathbf{r}) \quad F_{s\mathbf{k}}(\mathbf{r}) = \frac{e^{i\varphi_{\mathbf{k}}}}{\sqrt{2L^2}} \exp(i\mathbf{k} \cdot \mathbf{r}) R^{-1}[\theta(s\mathbf{k})] |s\rangle$$

$$R(\theta \pm 2\pi) = -R(\theta) \quad R(-\pi) = -R(+\pi)$$

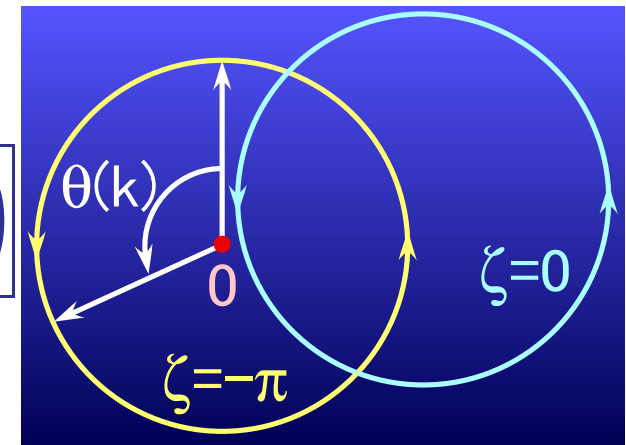
$$\varepsilon_s(\mathbf{k}) = s\gamma|\mathbf{k}| \quad s = \pm 1$$

$$\psi(T) = e^{-i\zeta} \psi(0)$$

Pseudo spin $\Rightarrow$ **Berry's phase**

$$\zeta = -i \int_0^T dt \left\langle s\mathbf{k}(t) \left| \frac{d}{dt} \right| s\mathbf{k}(t) \right\rangle = -\pi$$

$$\frac{1}{\sqrt{2}} \begin{pmatrix} s \\ e^{i\theta} \end{pmatrix}$$



Landau levels at $\varepsilon = 0$ [J.W. McClure, PR 104, 666 (1956)]

$$\chi = -\frac{g_v g_s \gamma^2}{6\pi} \left(\frac{e}{c\hbar} \right)^2 \int \left(-\frac{\partial f(\varepsilon)}{\partial \varepsilon} \right) \delta(\varepsilon) d\varepsilon$$

$$\varepsilon_n \propto \pm \sqrt{|n|B}$$

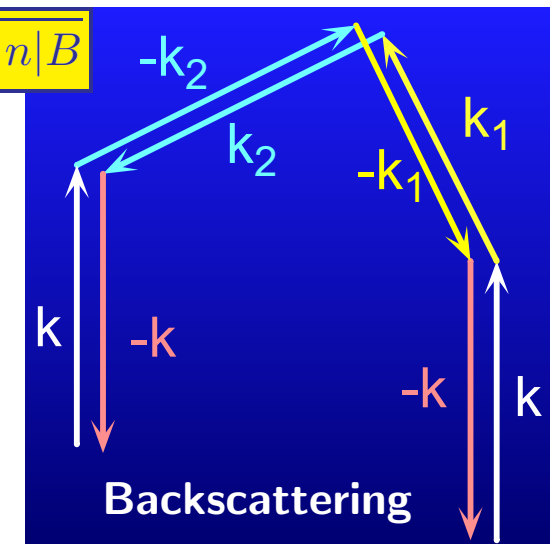
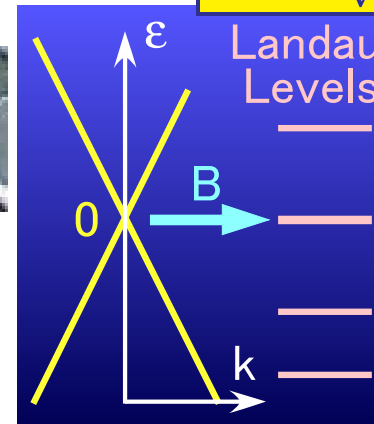


Absence of backscattering

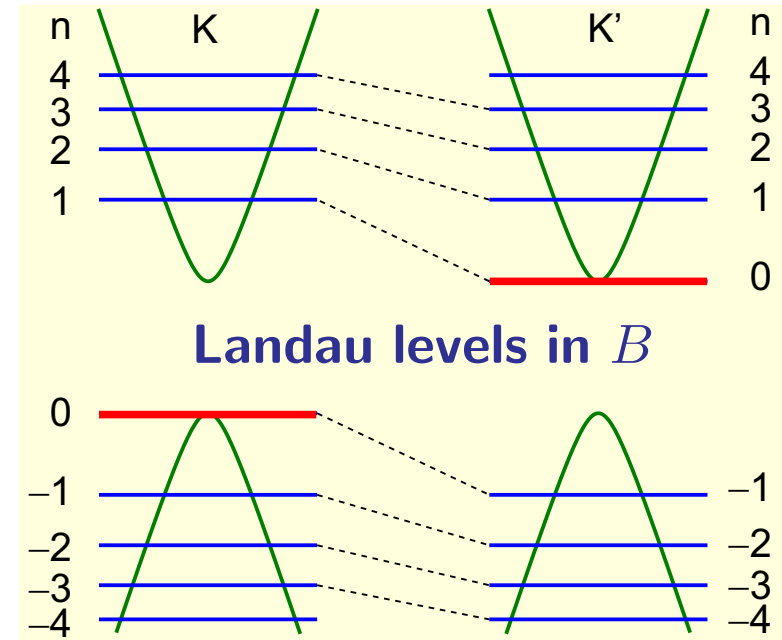
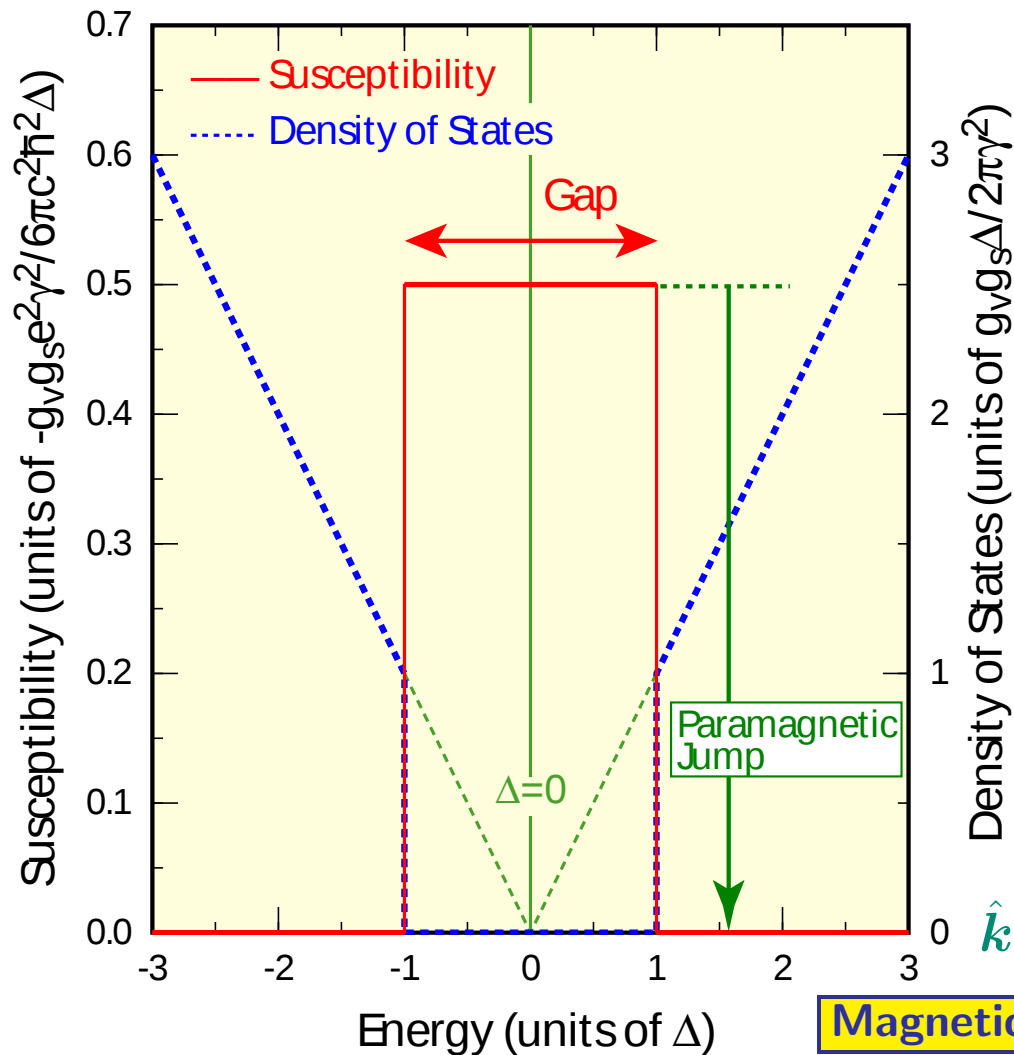
Metallic CN with scatterers

$\Rightarrow$ **Perfect conductor**

T. Ando & T. Nakanishi, JPSJ 67, 1704 (1998)



Band-Gap Effect [M. Koshino & T. Ando, PRB 81, 195431 (2010)]



Hamiltonian at band edge

$$\mathcal{H} = \frac{\hbar^2 \hat{k}^2}{2m^*} \pm \frac{e\hbar}{2m^*c} B \quad m^* = \frac{\hbar^2 \Delta}{\gamma^2}$$

$$\hat{\mathbf{k}} = \frac{\vec{\nabla}}{i} + \frac{e\mathbf{A}}{c\hbar}$$

$$\chi_P(\epsilon) = + \left(\frac{e\hbar}{2m^*c} \right)^2 D(\epsilon)$$

$$\chi_L(\epsilon) = - \frac{1}{3} \left(\frac{e\hbar}{2m^*c} \right)^2 D(\epsilon)$$

$$D(\epsilon) = \frac{g_s g_v m^*}{2\pi \hbar^2}$$

Graphene with a gap

$$\begin{pmatrix} \Delta & \gamma \hat{k}_- \\ \gamma \hat{k}_+ & -\Delta \end{pmatrix} \Rightarrow \chi(\epsilon) = - \frac{g_v g_s e^2 \gamma^2}{6\pi c^2 \gamma^2} \frac{\theta(\Delta - |\epsilon|)}{2\Delta}$$

$\rightarrow \delta(\epsilon)$

Magnetic moment

Pauli

Landau

Diamagnetic Susceptibility of Spatially Varying Field

Induced current $\hat{K}(\mathbf{r}) = (K_{\mu\nu}(\mathbf{r}))$

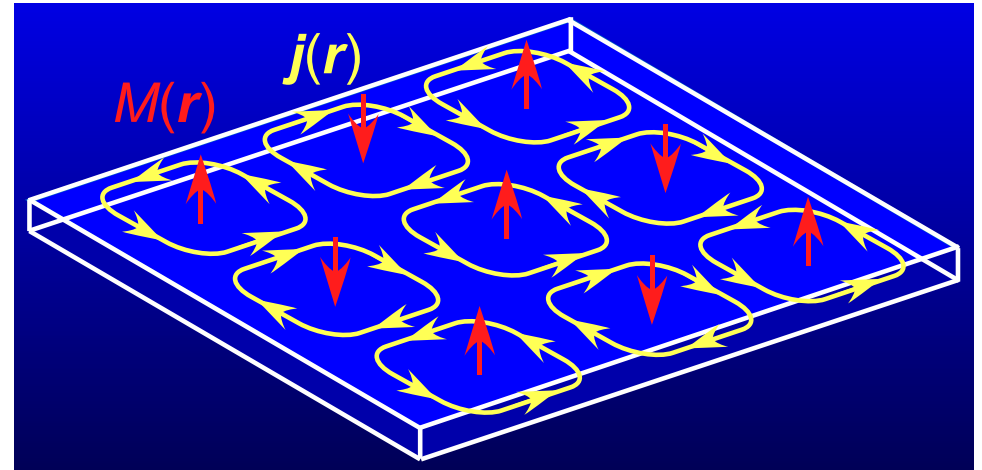
$$\mathbf{j}(\mathbf{r}) = \int \hat{K}(\mathbf{r} - \mathbf{r}') \mathbf{A}(\mathbf{r}') d\mathbf{r}'$$

$$\Leftrightarrow \mathbf{j}(\mathbf{q}) = \hat{K}(\mathbf{q}) \mathbf{A}(\mathbf{q})$$

Gauge invariance

Current conservation

$$K_{\mu\nu}(\mathbf{q}) = \mathbf{K}(\mathbf{q}) \left[\delta_{\mu\nu} - \frac{q_\mu q_\nu}{q^2} \right]$$



Magnetic moment: $M(\mathbf{r}) \Leftrightarrow \mathbf{j}(\mathbf{r}) = c (\nabla \times \hat{z}) M(\mathbf{r}) \quad H(\mathbf{r}) = \nabla \times \mathbf{A}(\mathbf{r})$

Diamagnetic susceptibility: $\chi(\mathbf{q}) = \frac{1}{cq^2} \mathbf{K}(\mathbf{q}) \Leftrightarrow M(\mathbf{q}) = \chi(\mathbf{q}) H(\mathbf{q})$

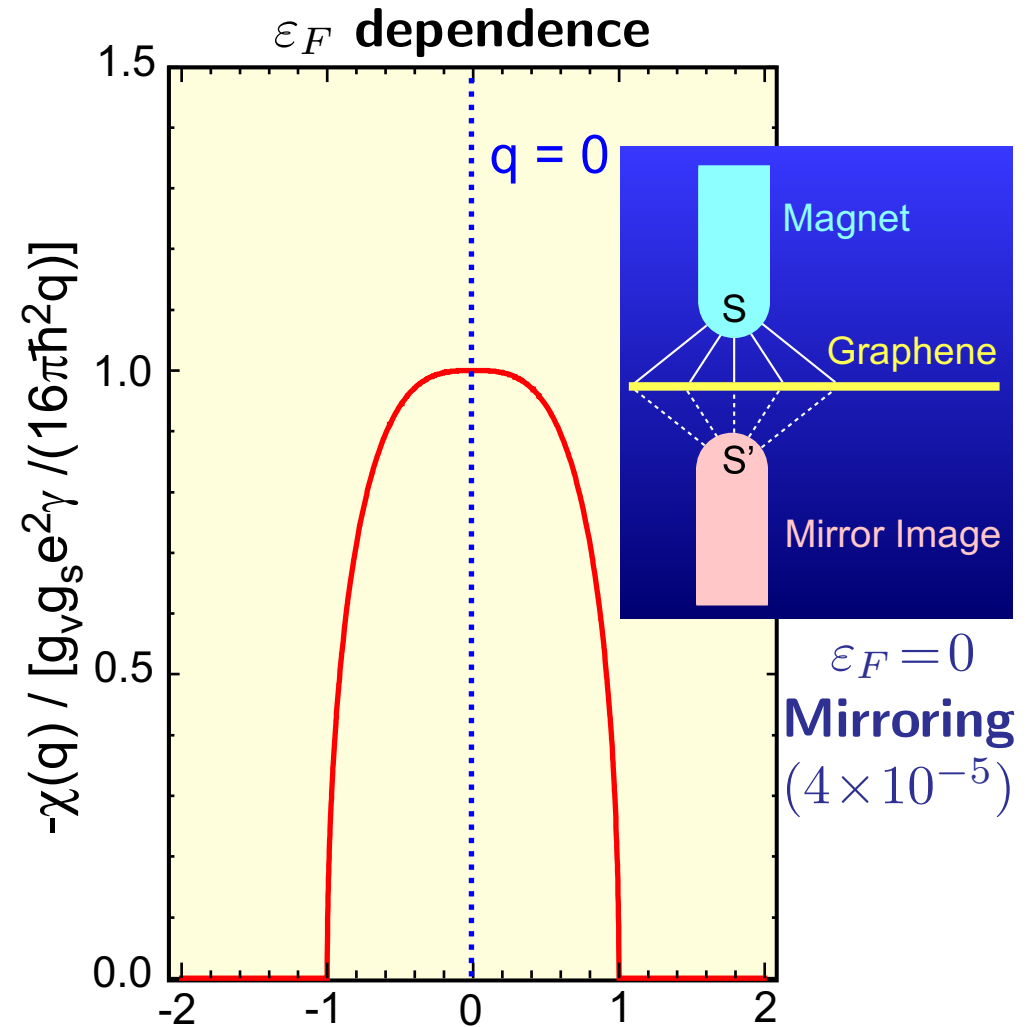
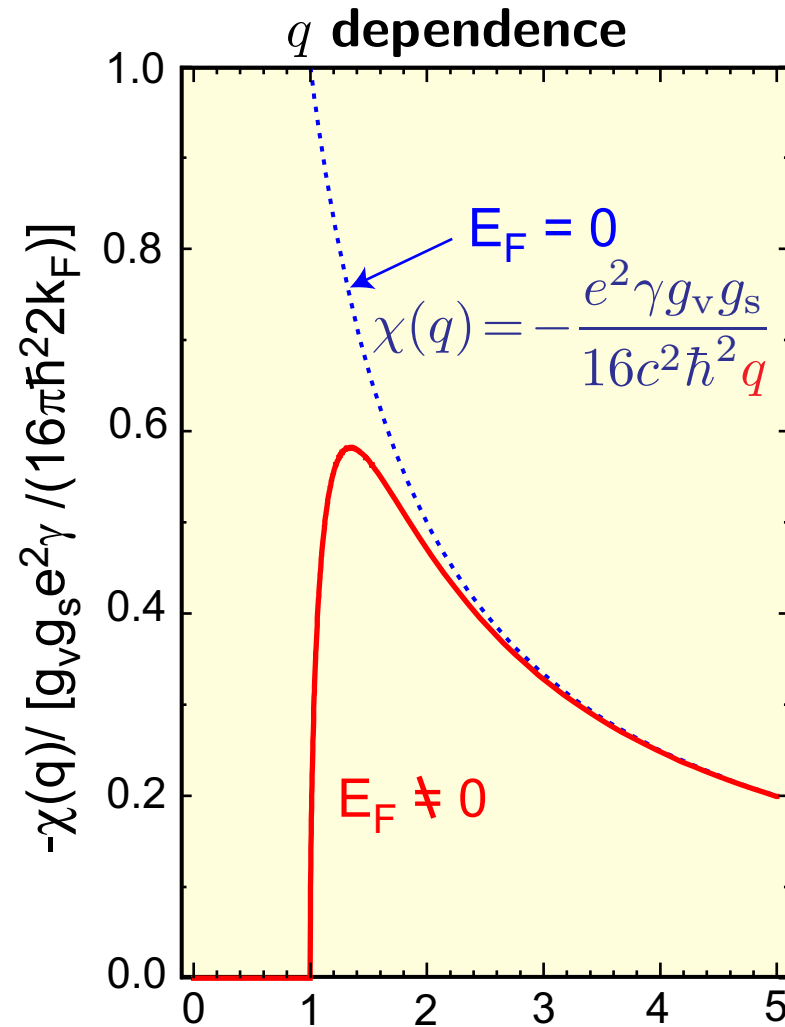
$$\chi(q) = -\frac{e^2 \gamma g_v g_s}{16c^2 \hbar^2} \frac{1}{q} \theta(q - 2k_F) \left[1 + \frac{2}{\pi} \frac{2k_F}{q} \sqrt{1 - \left(\frac{2k_F}{q} \right)^2} - \frac{2}{\pi} \sin^{-1} \left(\frac{2k_F}{q} \right) \right]$$

$$\Rightarrow \chi(q) = \begin{cases} 0 & (q < 2k_F) \\ \sim q^{-1} & (q > 2k_F) \end{cases} \quad \int_{-\infty}^{\infty} \chi(q; \varepsilon_F) d\varepsilon_F = -\frac{g_v g_s \gamma^2}{6\pi} \left(\frac{e}{c\hbar} \right)^2$$

$$\mathbf{2DEG:} \quad \chi(q) = -\frac{e^2 g_v g_s}{24\pi m^* c^2} \left(1 - \left[1 - \left(\frac{2k_F}{q} \right)^2 \right]^{3/2} \theta(q - 2k_F) \right)$$

Diamagnetic Susceptibility $\chi(q)$ of Spatially Varying Field

M. Koshino, Y. Arimura, and T. Ando, PRL 102, 177203 (2009)



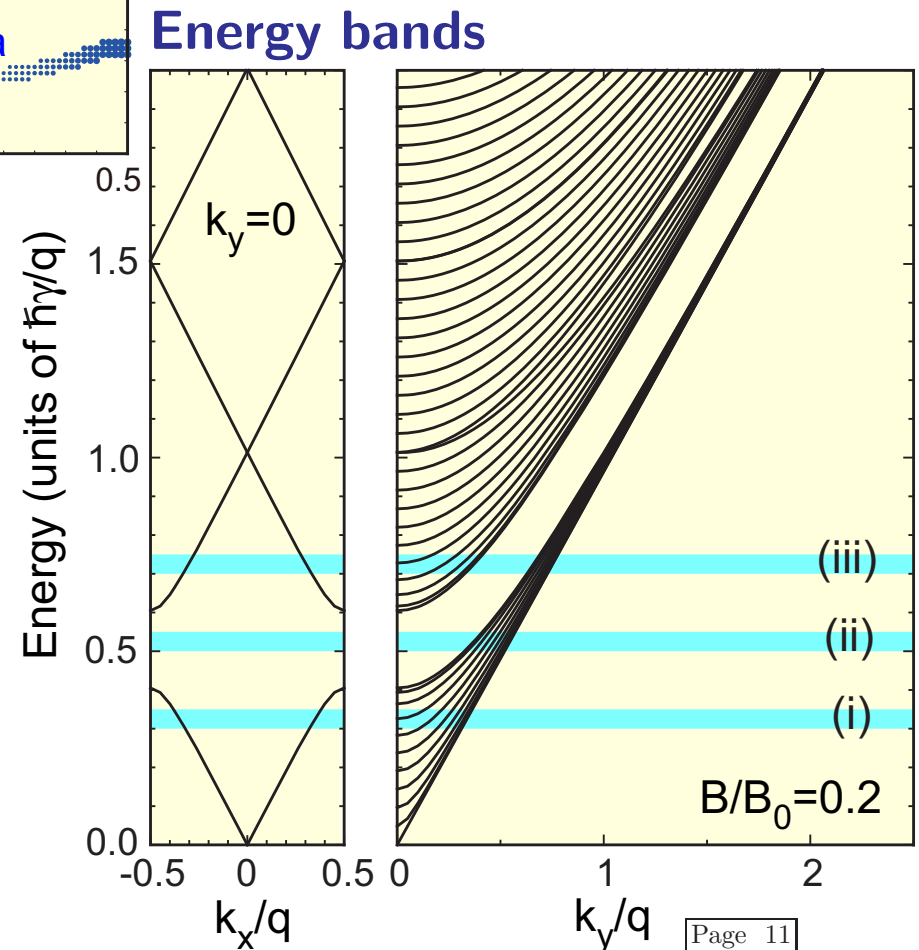
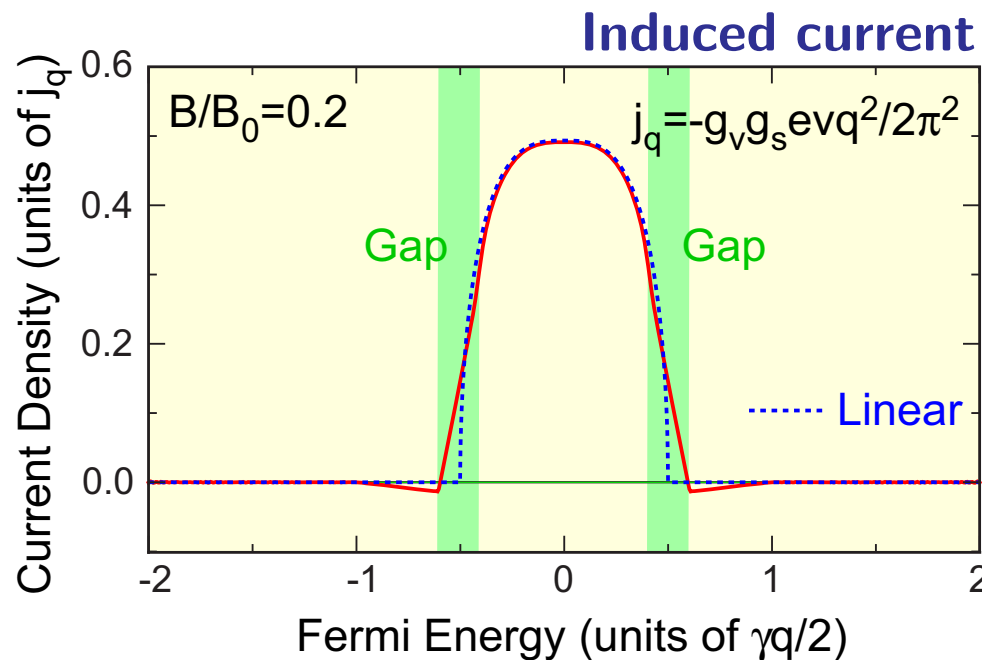
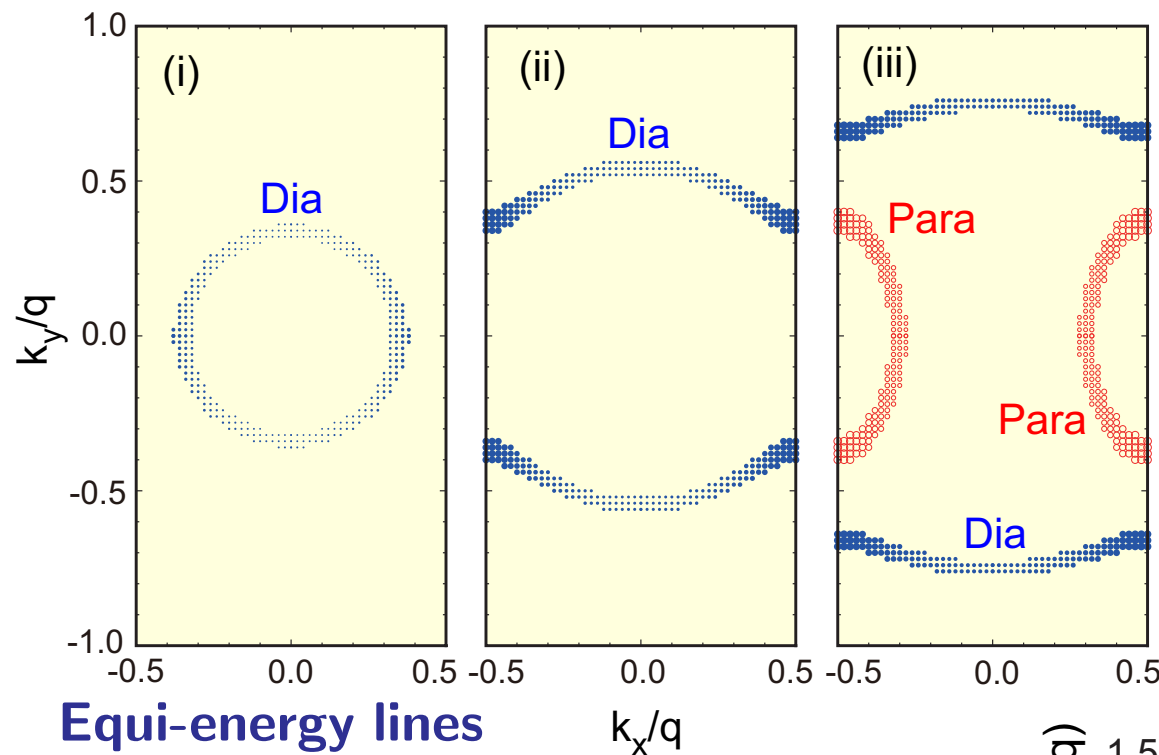
No response to
slowly-varying field $q/(2k_F)$

Uniform limit: $\lim_{q \rightarrow 0} \chi(q) = -\frac{g_v g_s \gamma^2}{6\pi} \left(\frac{e}{c\hbar}\right)^2 \frac{E_F}{(\gamma q/2)} \delta(\varepsilon_F)$

Magnetic Superlattice

*M. Koshino & T. Ando,
Solid State Commun.
151, 1054 (2011)*

*cf. C.H. Park et al., PRL
101, 126804 (2008), ...*



Zero-Mode Anomaly in Conductivity

Density of states: $D(\varepsilon) = \frac{|\varepsilon|}{2\pi\gamma^2} \Rightarrow$ **Zero-gap semiconductor**

Boltzmann conductivity

$$\sigma(\varepsilon_F) = g_s g_v e^2 D^* D(\varepsilon_F) = \frac{g_s g_v}{4} \frac{e^2}{\pi^2 \hbar} \frac{1}{W}$$

Einstein relation

$$D^* = v_F^2 \tau = \frac{\gamma^2}{\hbar^2} \tau \quad W = \frac{n_i u^2}{4\pi\gamma^2} \ll 1$$

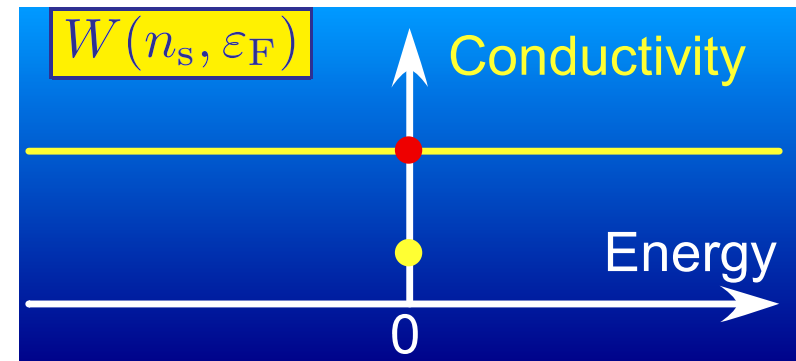
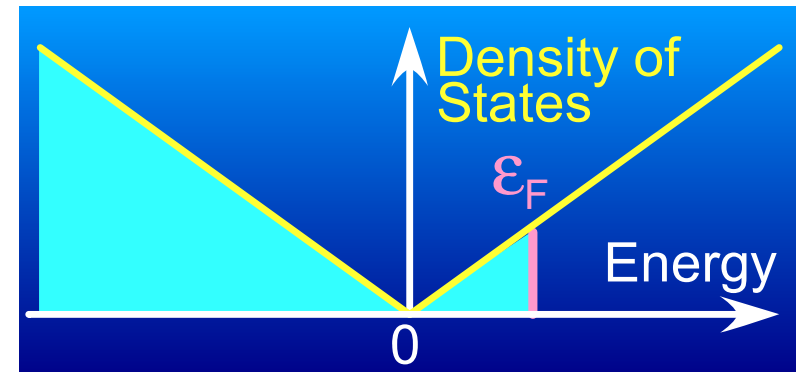
$$\frac{1}{\tau} = \frac{2\pi}{\hbar} n_i u^2 D(\varepsilon_F) \propto 2\pi W |\varepsilon_F| / \hbar$$

$$\tau \propto D(\varepsilon_F)^{-1} \quad \begin{array}{l} u \text{ Impurity strength} \\ n_i \text{ Impurity density} \end{array}$$

$\Rightarrow$ **Metal!**

Short-range ($k_F d \ll 1$)

$\Rightarrow \sigma(0)$ for $D(0)=0$ ($\Rightarrow 2e^2/\pi^2 \hbar$)



Singularity at the Dirac point ($\varepsilon_F = 0$) $\Leftrightarrow$ **Fermi energy scaling**

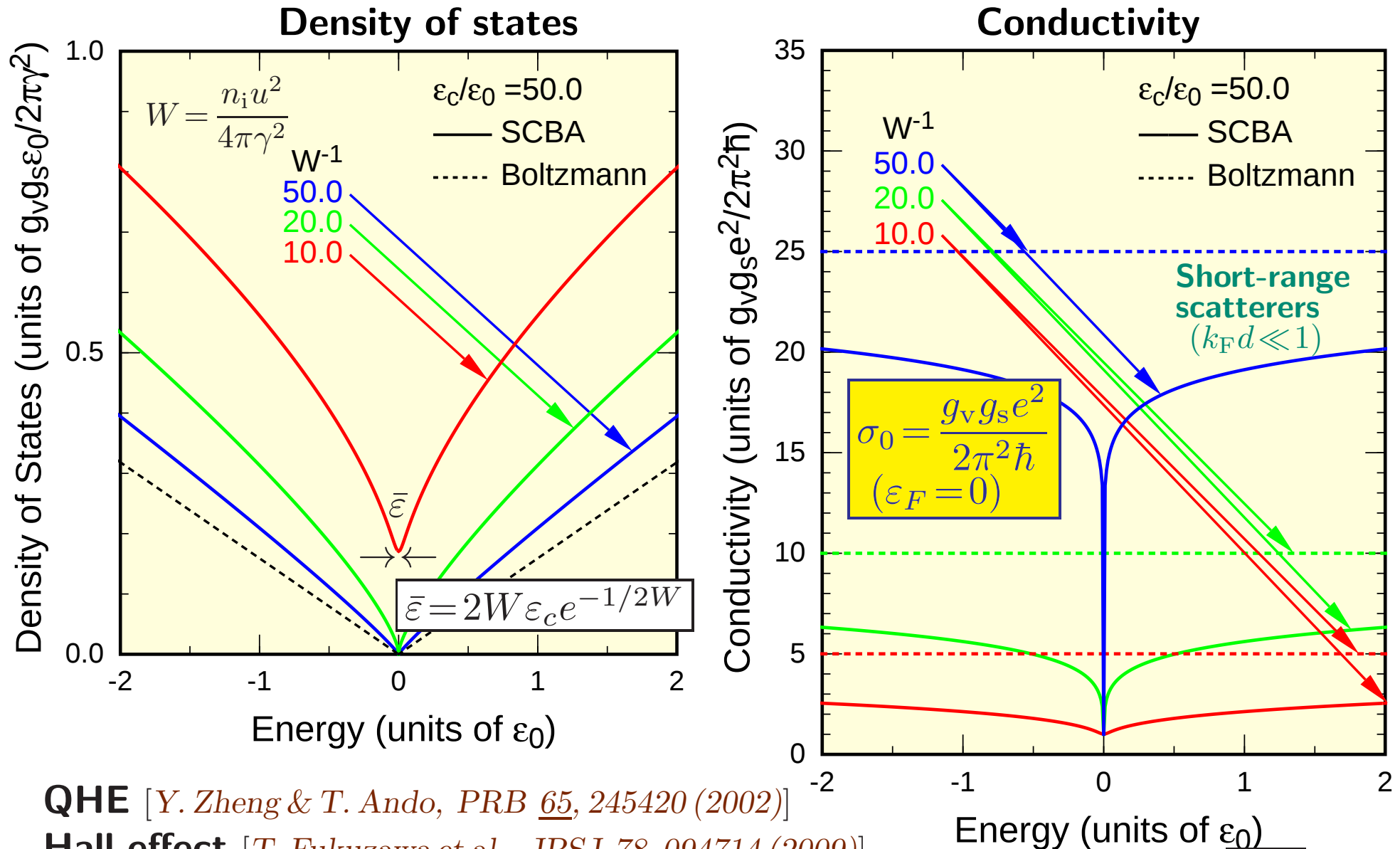
$$\sigma_{xx}(B) = \sigma_{xx}\left(\frac{\hbar\omega_B}{\varepsilon_F}\right) \quad \boxed{\hbar\omega_B \propto \sqrt{B}} \quad \sigma(\omega) = \sigma\left(\frac{\hbar\omega}{\varepsilon_F}\right) = \begin{cases} \frac{g_s g_v}{4} \frac{e^2}{2\pi^2 \hbar W} & (\omega = 0) \\ \frac{g_v g_s}{4} \frac{e^2}{4\hbar} & (\hbar\omega/\varepsilon_F \gg 1) \end{cases}$$

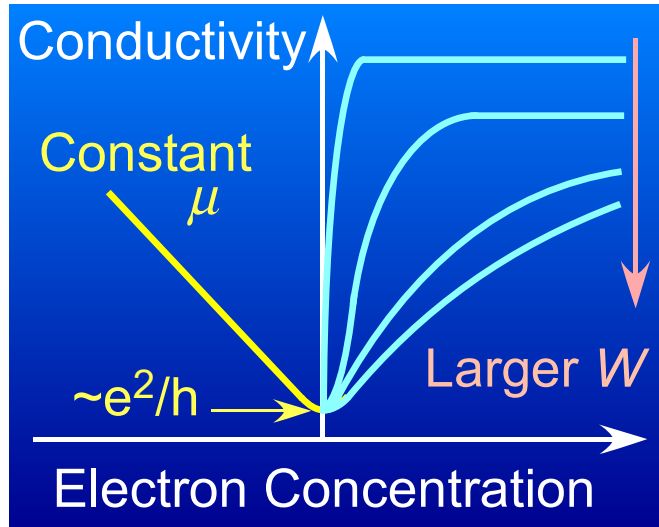
Magnetoconductivity **Intra-band**
Inter-band

Dynamical conductivity

Self-Consistent Born Approximation

[N.H. Shon and T. Ando, *J. Phys. Soc. Jpn.* **67**, 2421 (1998)]





Conductivity vs Carrier Concentration

K.S. Novoselov et al., Nature 438, 197 (2005)

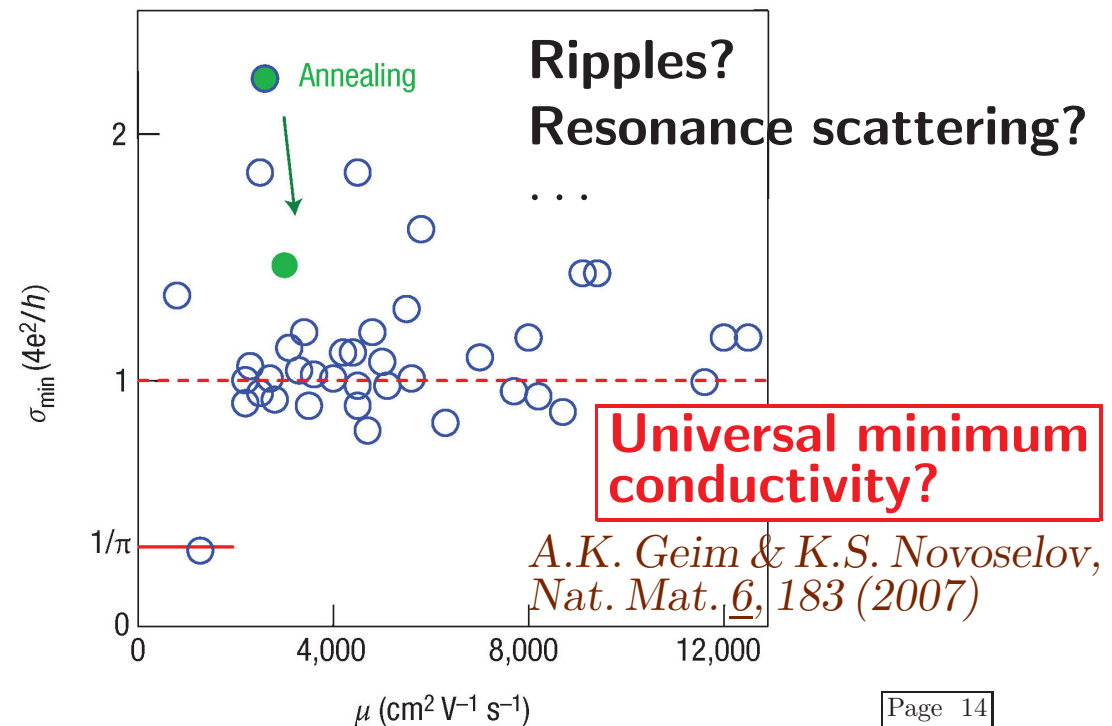
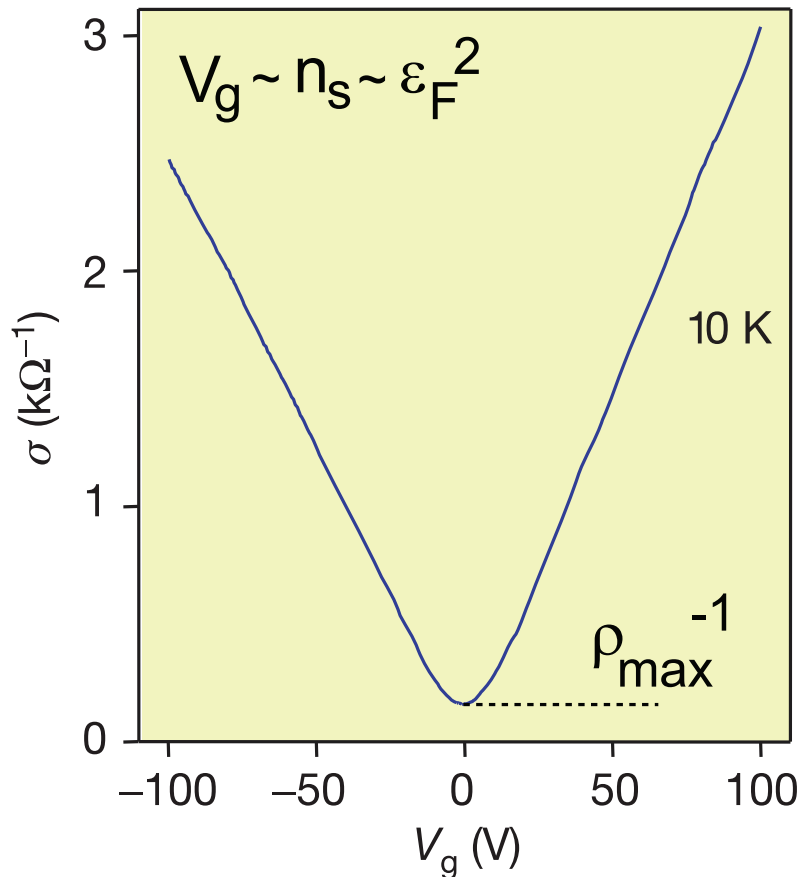
Scattering mechanisms

Boltzmann conductivity $\sigma(\varepsilon_F) = \frac{e^2}{\pi^2 \hbar} \frac{1}{4W}$

Constant mobility $\Leftrightarrow W \propto n_s^{-1}$

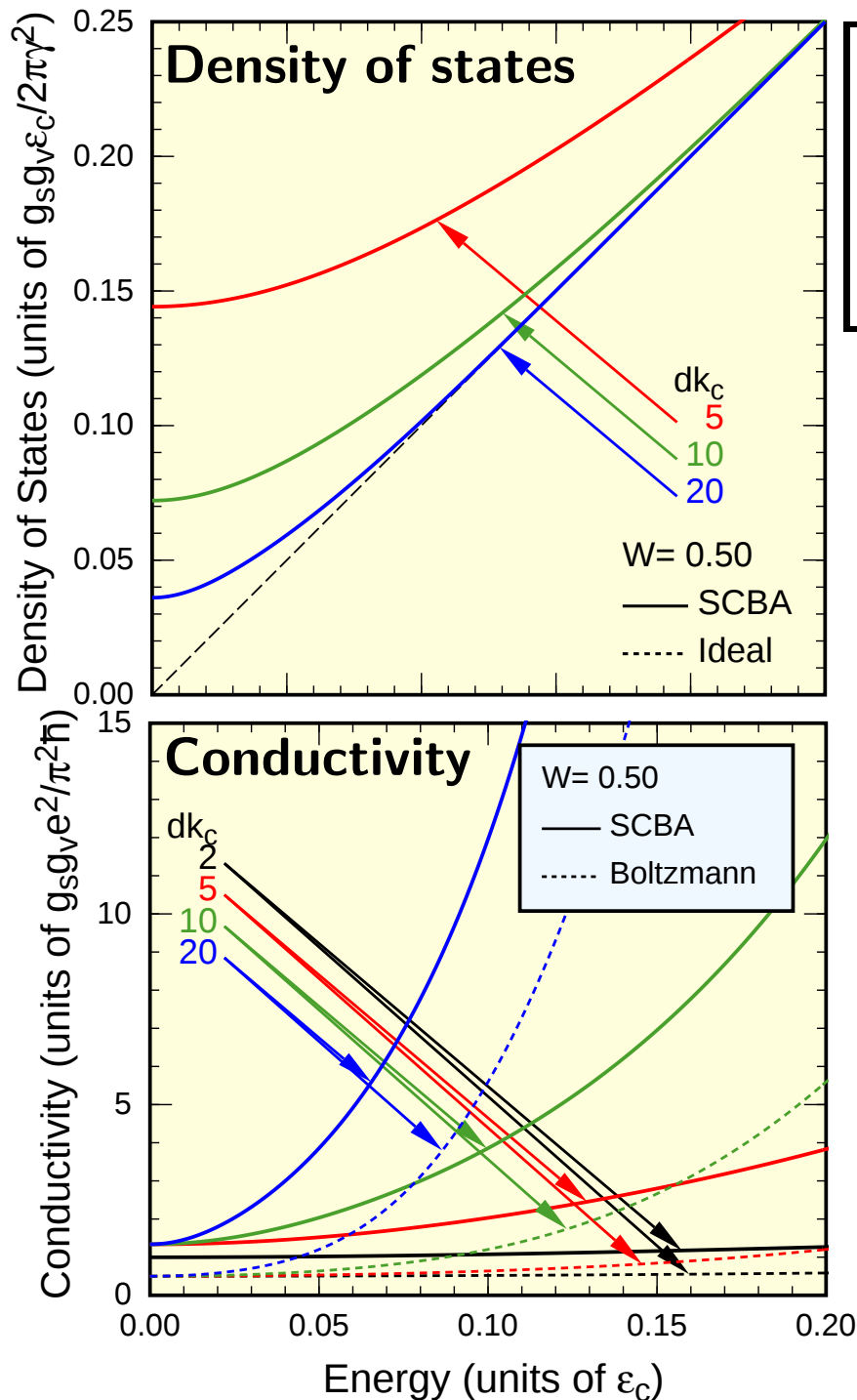
Charged impurity with screening

- *T. Ando, JPSJ 75, 074716 (2006)*
- *K. Nomura & A.H. MacDonald, PRL 96, 256602 (2006)*

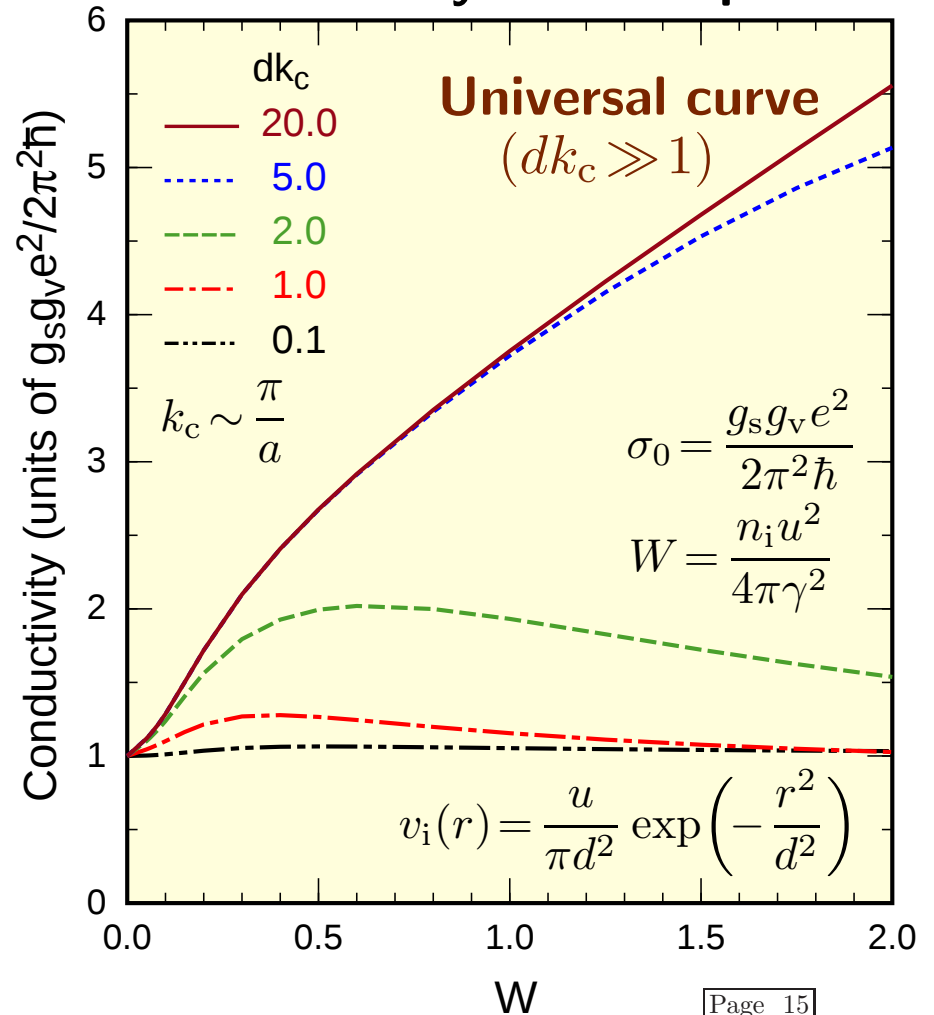


Scatterers with Gaussian Potential (Self-Consistent Born Approx.)

[M. Noro, M. Koshino & T. Ando
JPSJ 79, 094713 (2010)]



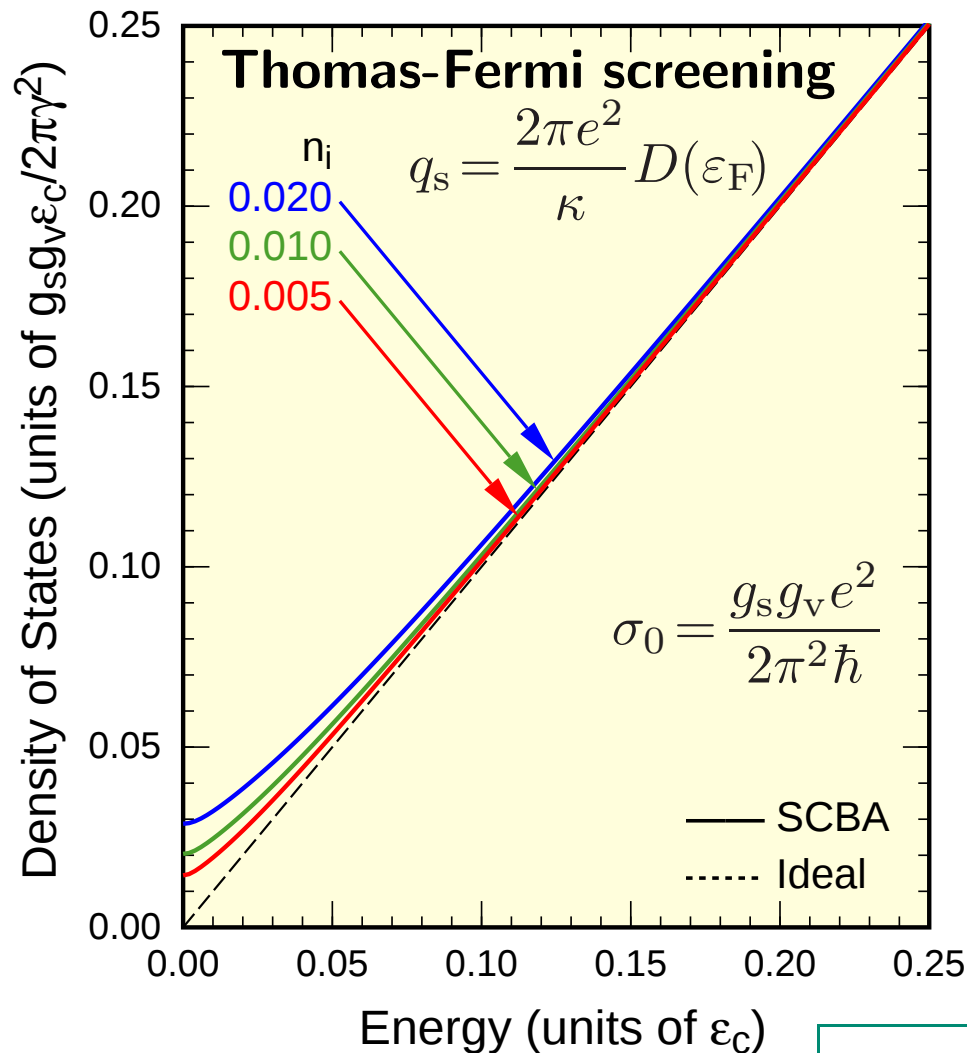
Conductivity at Dirac point



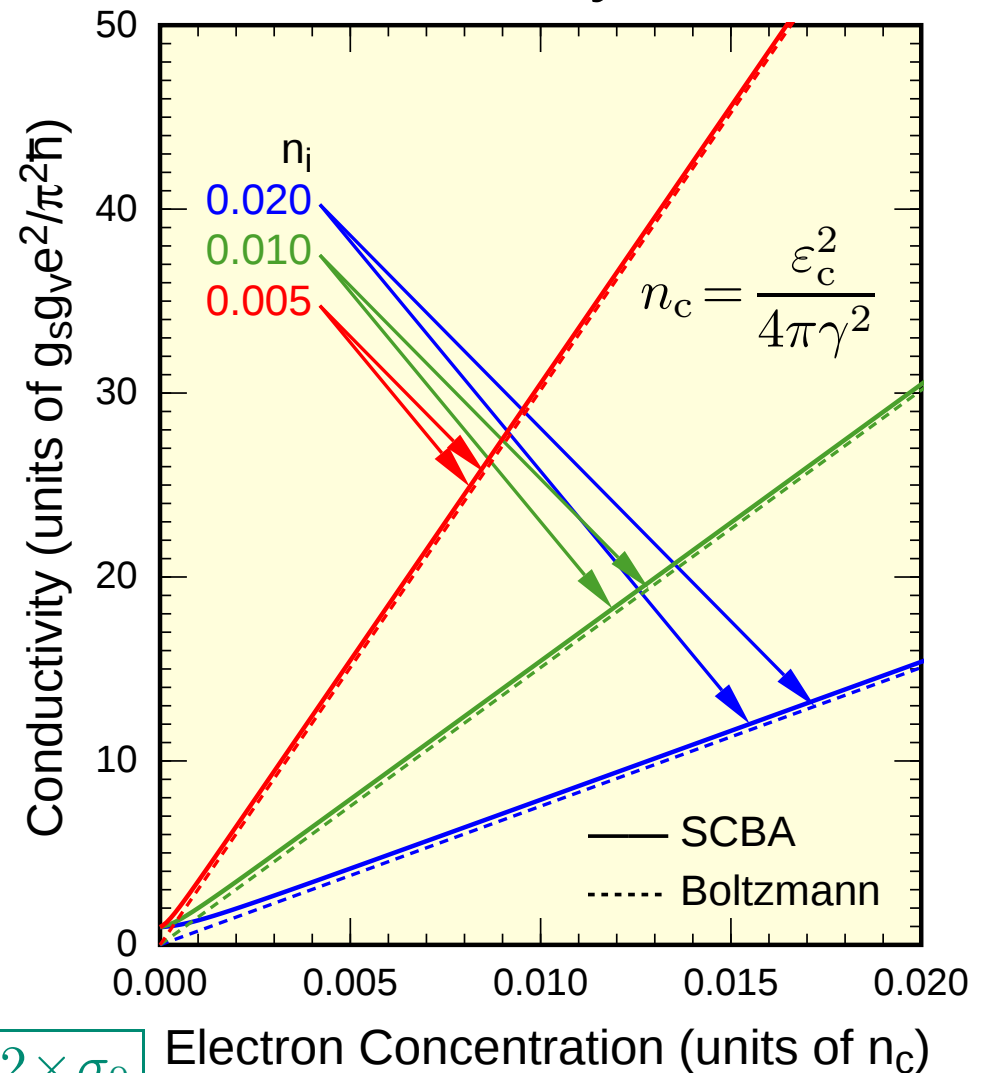
Charged Impurities (Self-Consistent Born Approximation)

[M. Noro, M. Koshino & T. Ando, *J. Phys. Soc. Jpn.* 79, 094713 (2010)]

Density of states



Conductivity vs n_s

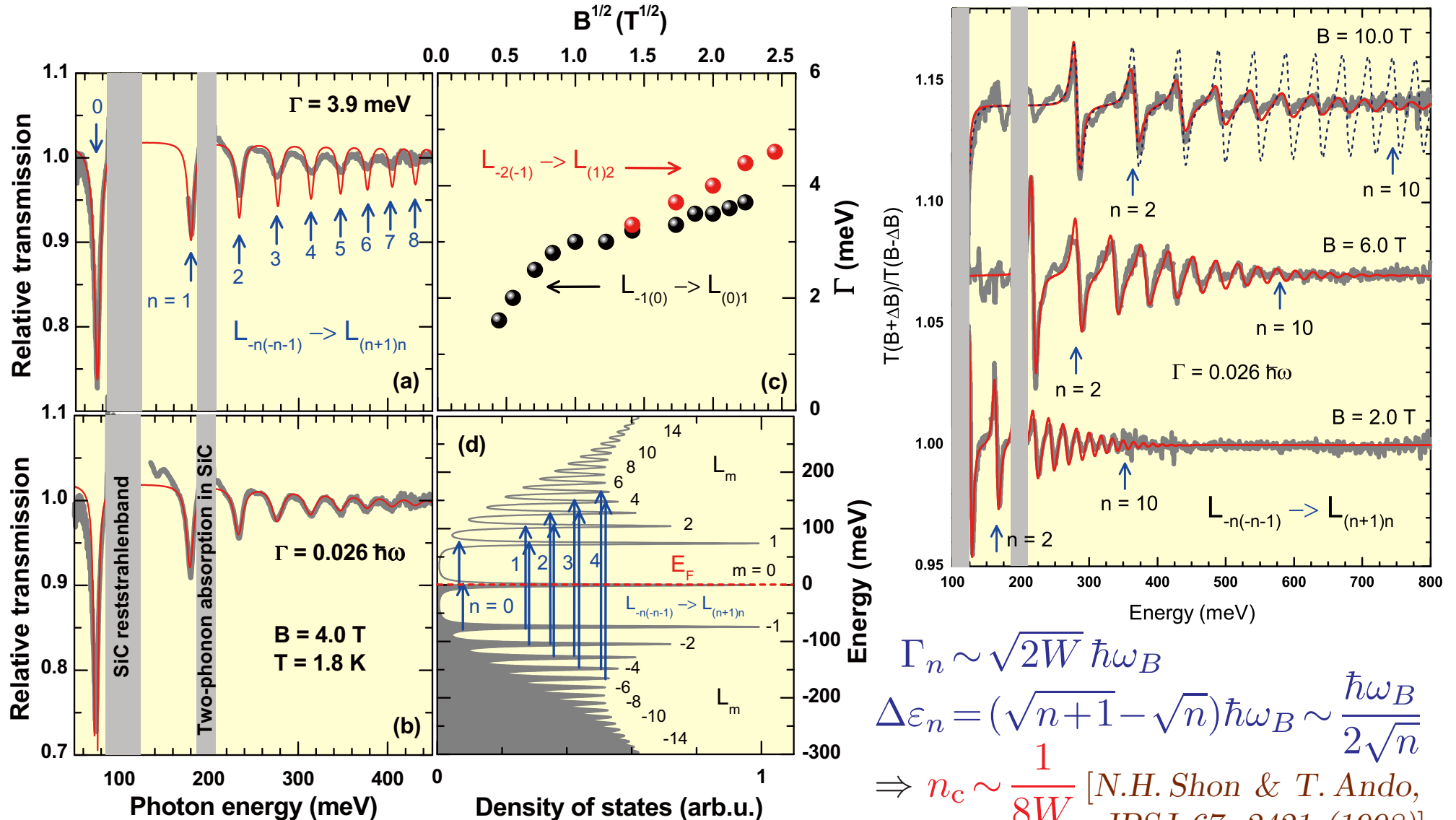


cf: X.-Z. Yan & C.S. Ting
Phys. Rev. B 80, 155423 (2009)

$$\sigma_{\min} \sim 2 \times \sigma_0$$

Interband Magnetoabsorption in Epitaxial Graphene

$\tau^{-1} \propto |\varepsilon|$ [M. Orlita et al., Phys. Rev. Lett. 107, 216603 (2011)]



Suspended, BN substrate, CVD grown, ...

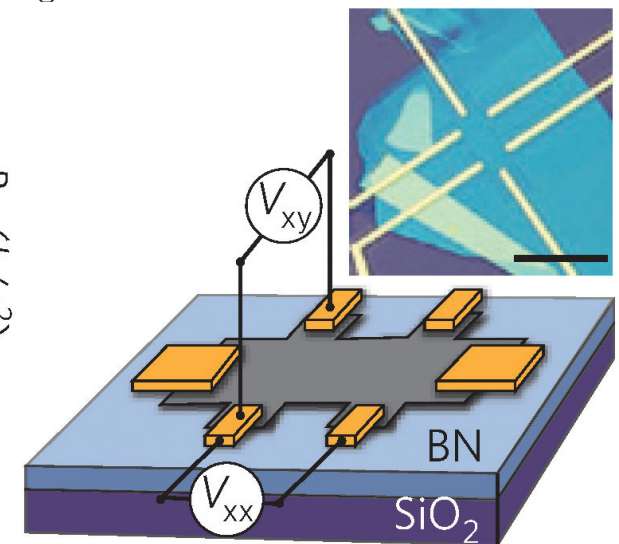
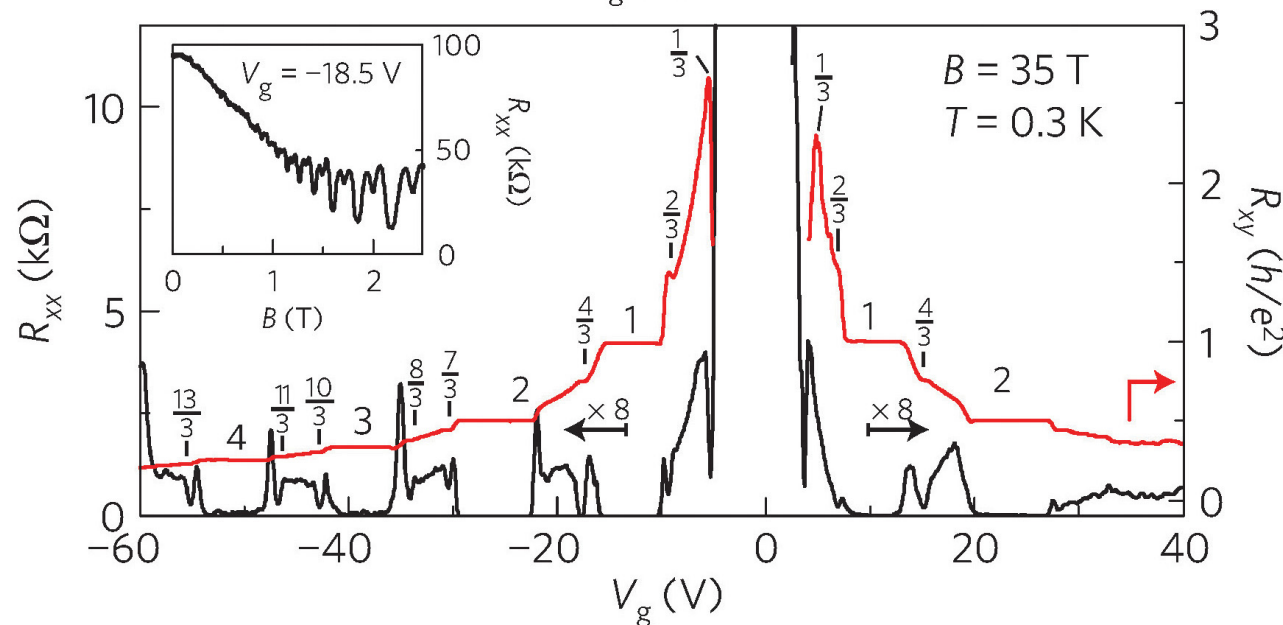
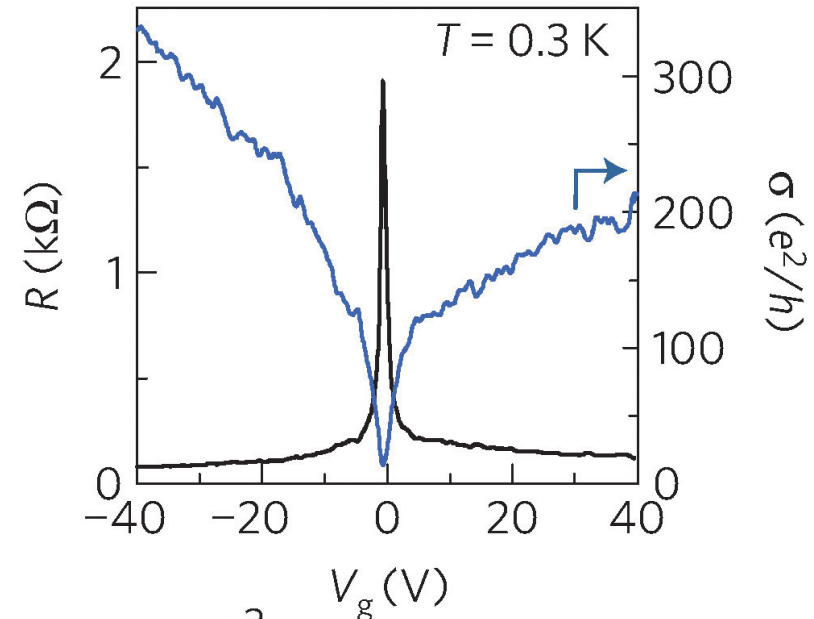
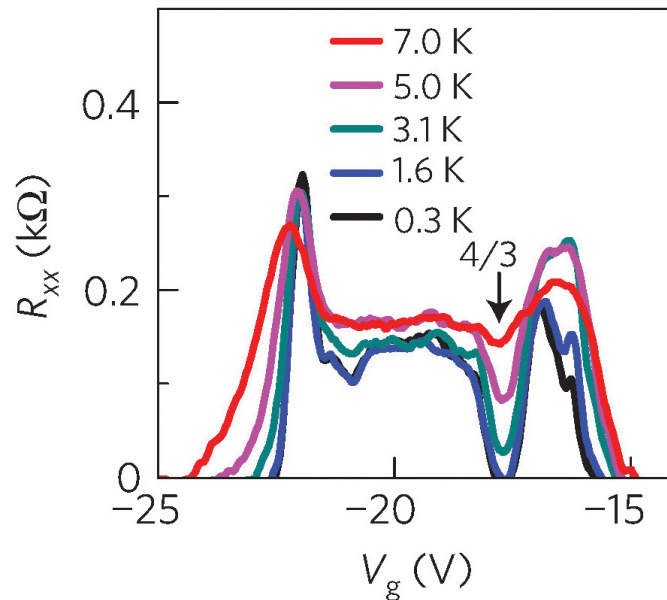
$$\Gamma_n \sim \sqrt{2W} \hbar\omega_B$$

$$\Delta\varepsilon_n = (\sqrt{n+1} - \sqrt{n}) \hbar\omega_B \sim \frac{\hbar\omega_B}{2\sqrt{n}}$$

$$\Rightarrow n_c \sim \frac{1}{8W} [N.H. Shon \& T. Ando, JPSJ \underline{67}, 2421 (1998)]$$

Fractional Quantum Hall Effect (BN Substrate)

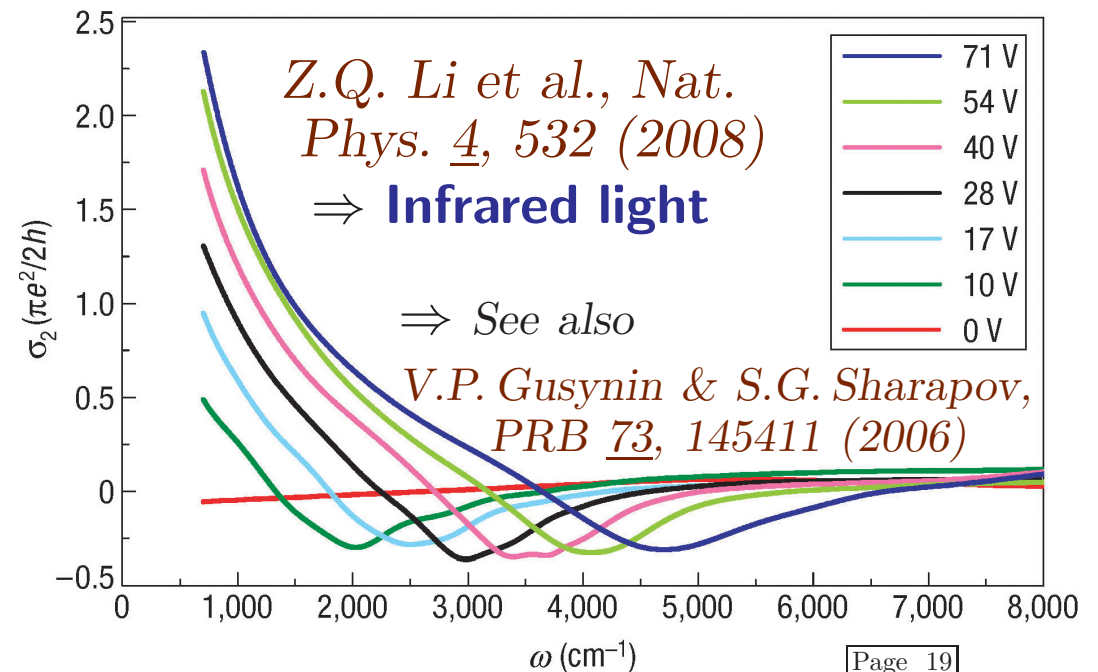
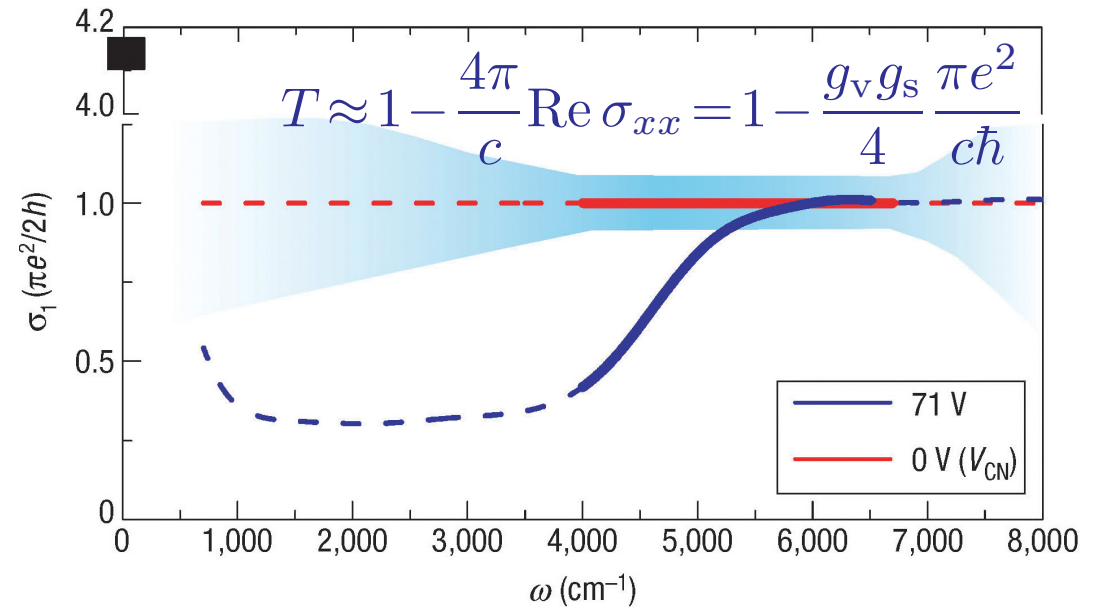
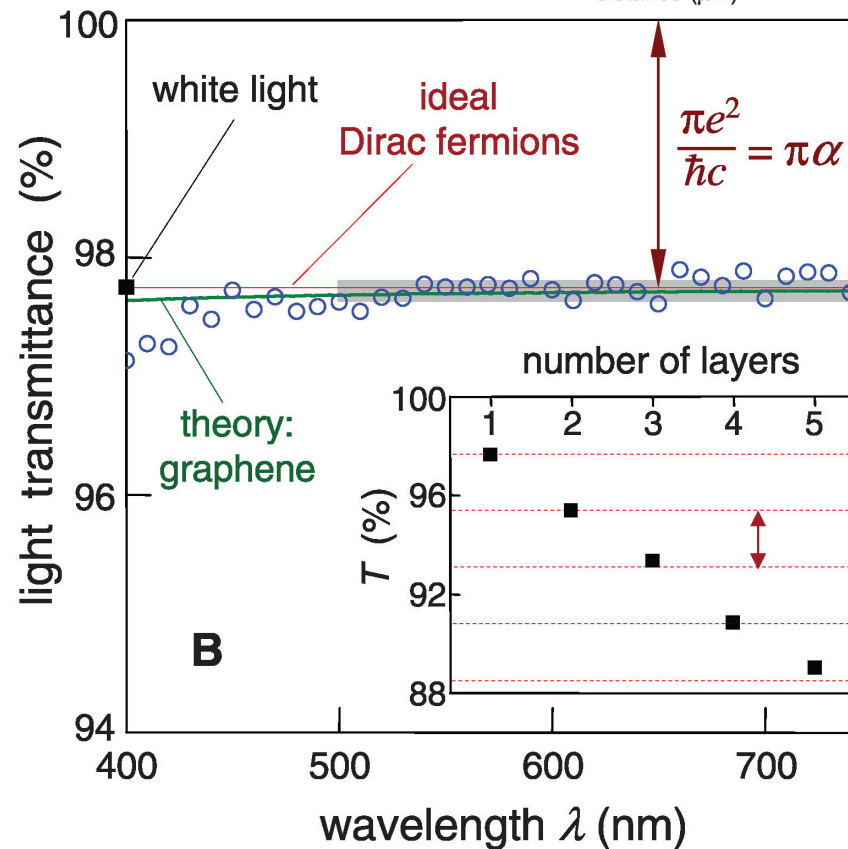
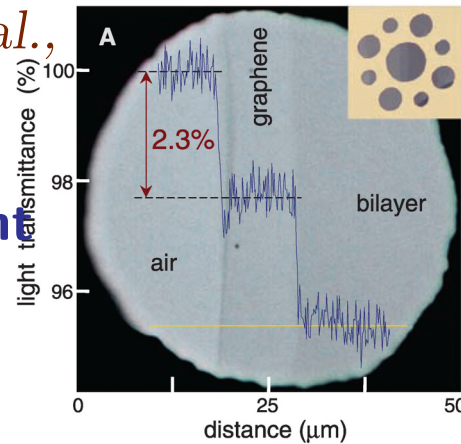
[C.R. Dean et al., *Nat. Phys.* 7, 693 (2011)]



Dynamical Conductivity: Experiments

R.R. Nair et al.,
Science **320**,
1308 (2008)

⇒ **Visible light**



Bilayer Graphene: AB Stacking (Bernal Stacking)

Quantum Hall effect in bilayer graphene

K.S. Novoselov et al., Nature 438, 197 (2005)

K.S. Novoselov et al., Nat. Phys. 2, 177 (2006)

ARPES [*T. Ohta et al., PRL 98, 206802 (2007)*]

Effective Hamiltonian in bilayer graphene

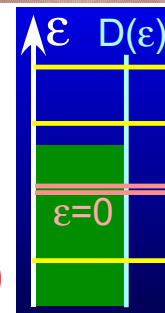
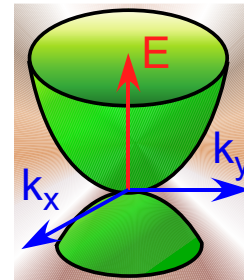
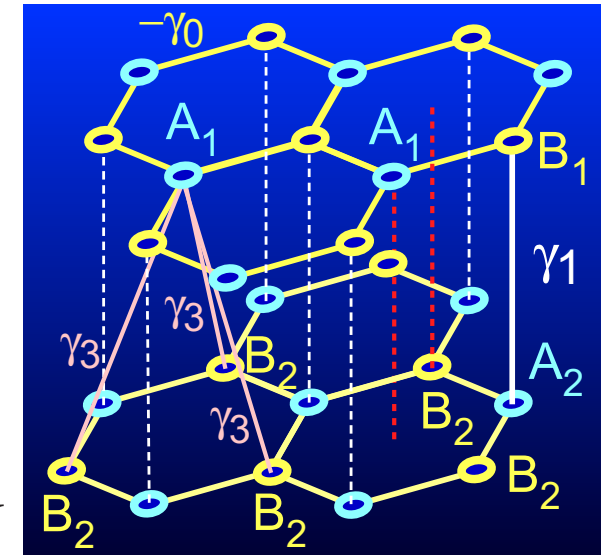
$$\mathcal{H} = \begin{pmatrix} A_1 & B_1 & A_2 & B_2 \\ 0 & \gamma \hat{k}_- & 0 & 0 \\ \gamma \hat{k}_+ & 0 & \Delta & 0 \\ 0 & \Delta & 0 & \gamma \hat{k}_- \\ 0 & 0 & \gamma \hat{k}_+ & 0 \end{pmatrix} \quad \begin{aligned} \hat{k}_{\pm} &= \hat{k}_x \pm i \hat{k}_y \\ \Delta &= \gamma_1 \approx 0.4 \text{ eV} \end{aligned}$$

$$\mathcal{H} \approx \frac{\hbar^2}{2m^*} \begin{pmatrix} 0 & \hat{k}_-^2 \\ \hat{k}_+^2 & 0 \end{pmatrix} \quad m^* = \frac{\hbar^2 \Delta}{2\gamma^2}$$

$$\varepsilon_n = \pm \hbar \omega_c \sqrt{n(n+1)} \quad (n=0, 1, \dots)$$

$\Rightarrow$ **Two Landau levels at $\varepsilon=0$**

Susceptibility $\Rightarrow \chi(\varepsilon) = -\frac{g_v g_s}{4\pi} \frac{e^2 \gamma}{c^2 \hbar^2} \ln \left| \frac{\Delta}{\varepsilon} \right|$ [S.A. Safran, *PRB* 30, 421 (1984)]



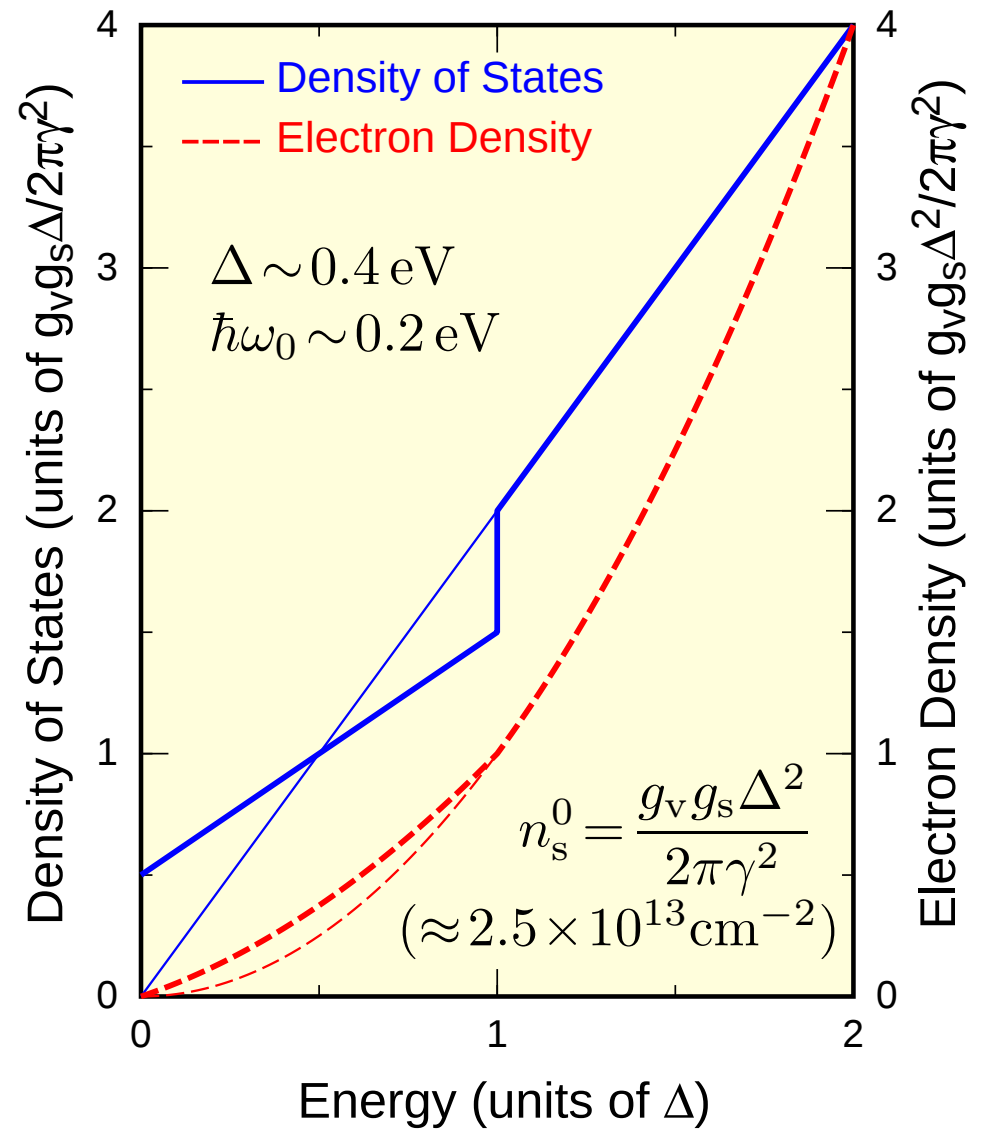
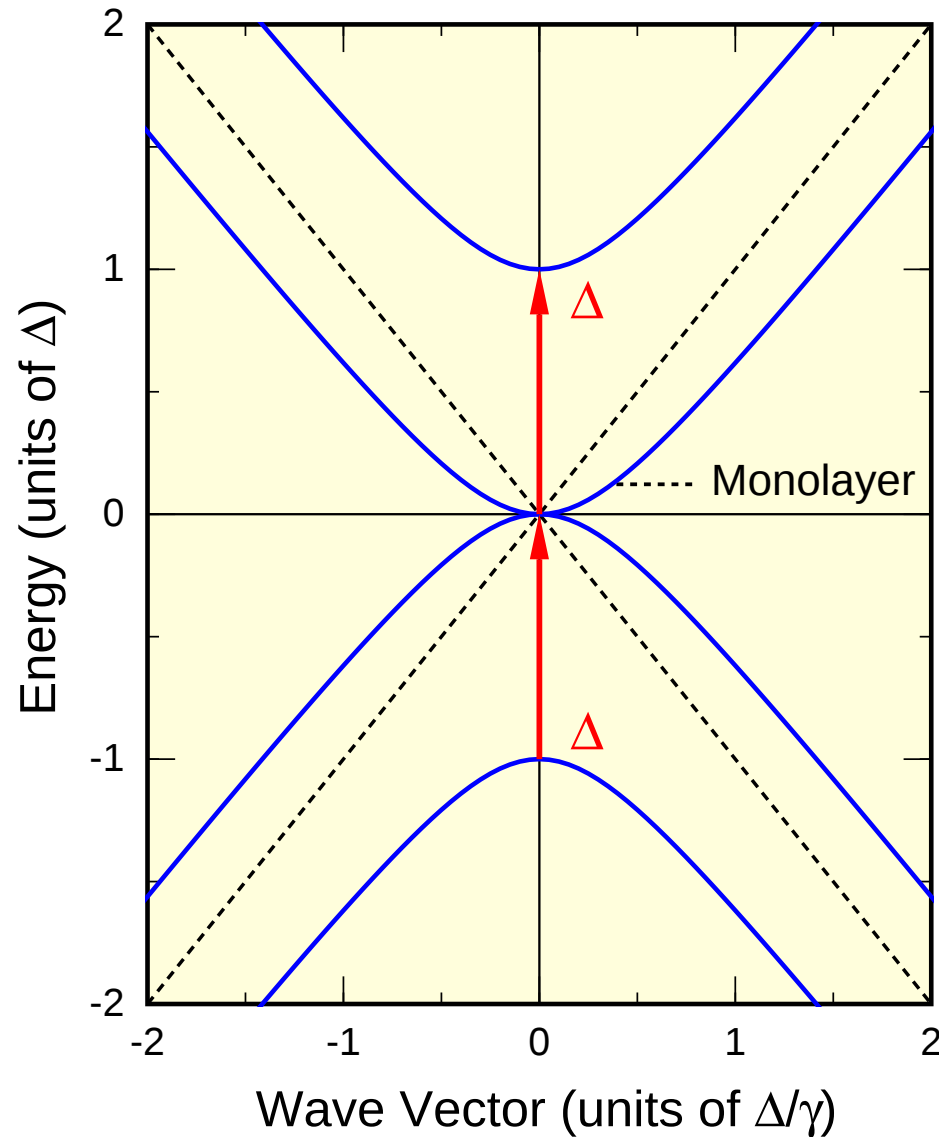
- *E. McCann and V.I. Falko, PRL 96, 086805 (2006)*
- *M. Koshino and T. Ando, PRB 73, 245403 (2006)*

Tight-binding models

- *S. Latil and L. Henrard, PRL 97, 036803 (2006)*
- *F. Guinea et al., PRB 73, 245426 (2006), ...*

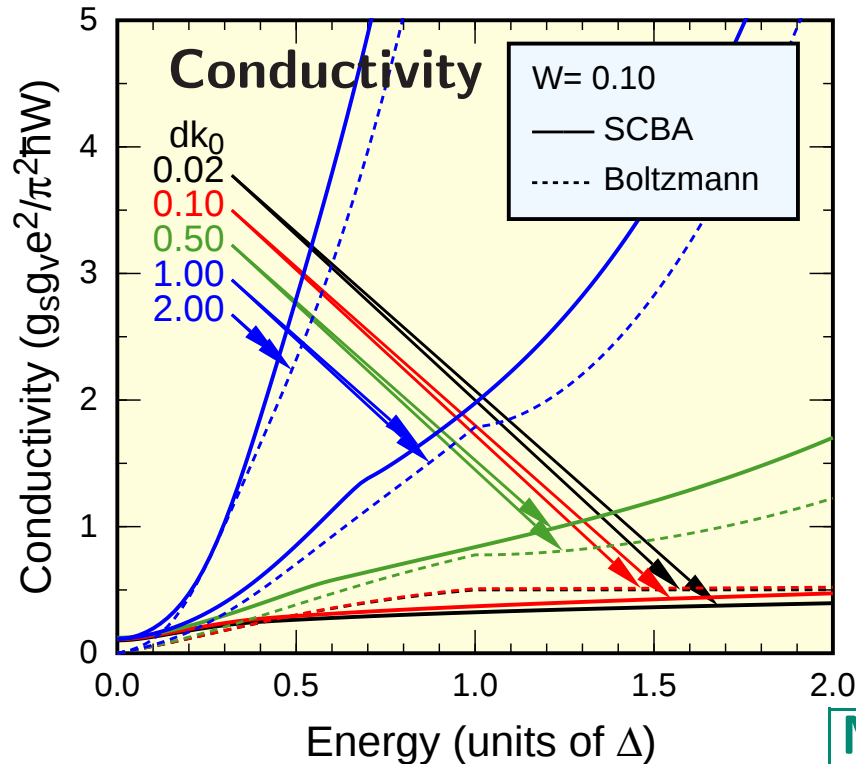
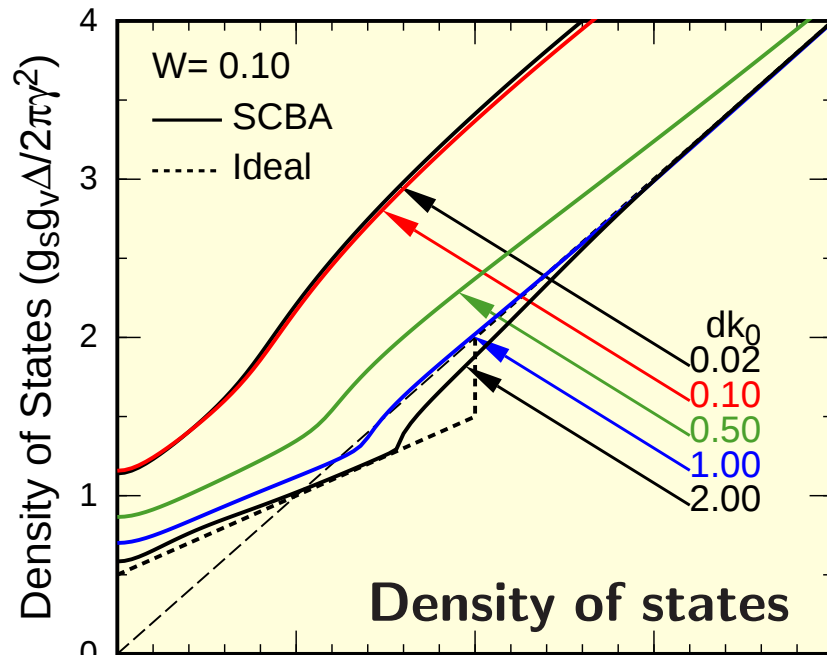
Energy Dispersion and Density of States of Bilayer Graphene

T. Ando, J. Phys. Soc. Jpn. 76, 104711 (2007)



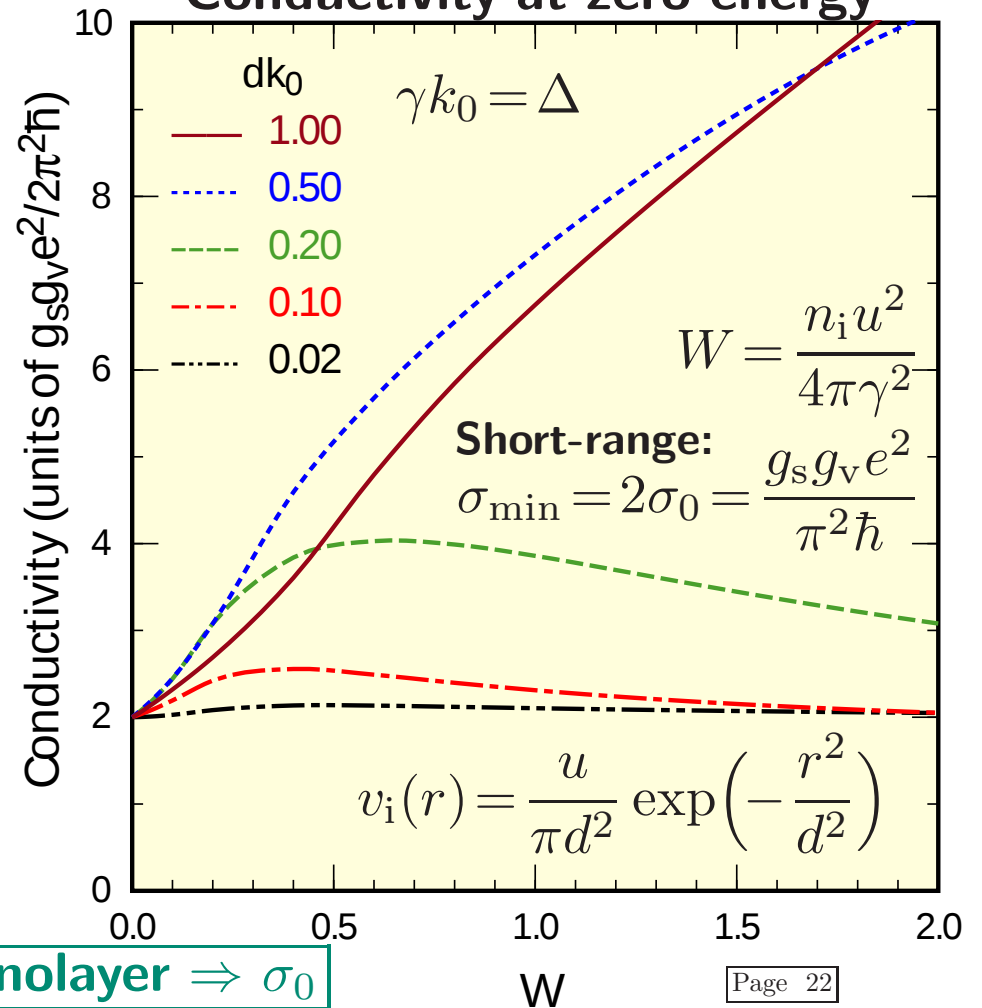
Scatterers with Gaussian Potential (Self-Consistent Born Approx.)

[T. Ando, *JPSJ* 80, 014707 (2011)]



Short-range: M. Koshino & T. Ando,
PRB 73, 245403 (2006)

Conductivity at zero energy

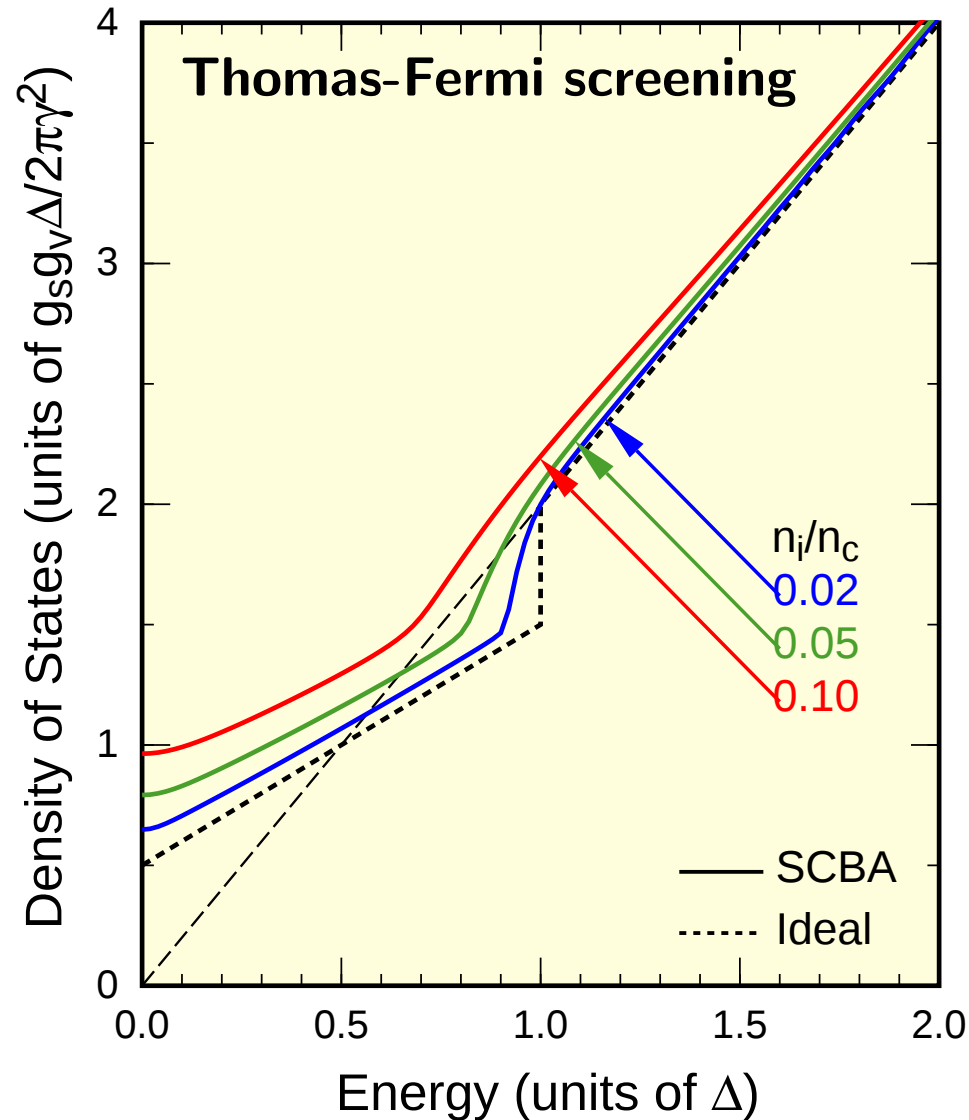


Monolayer $\Rightarrow \sigma_0$

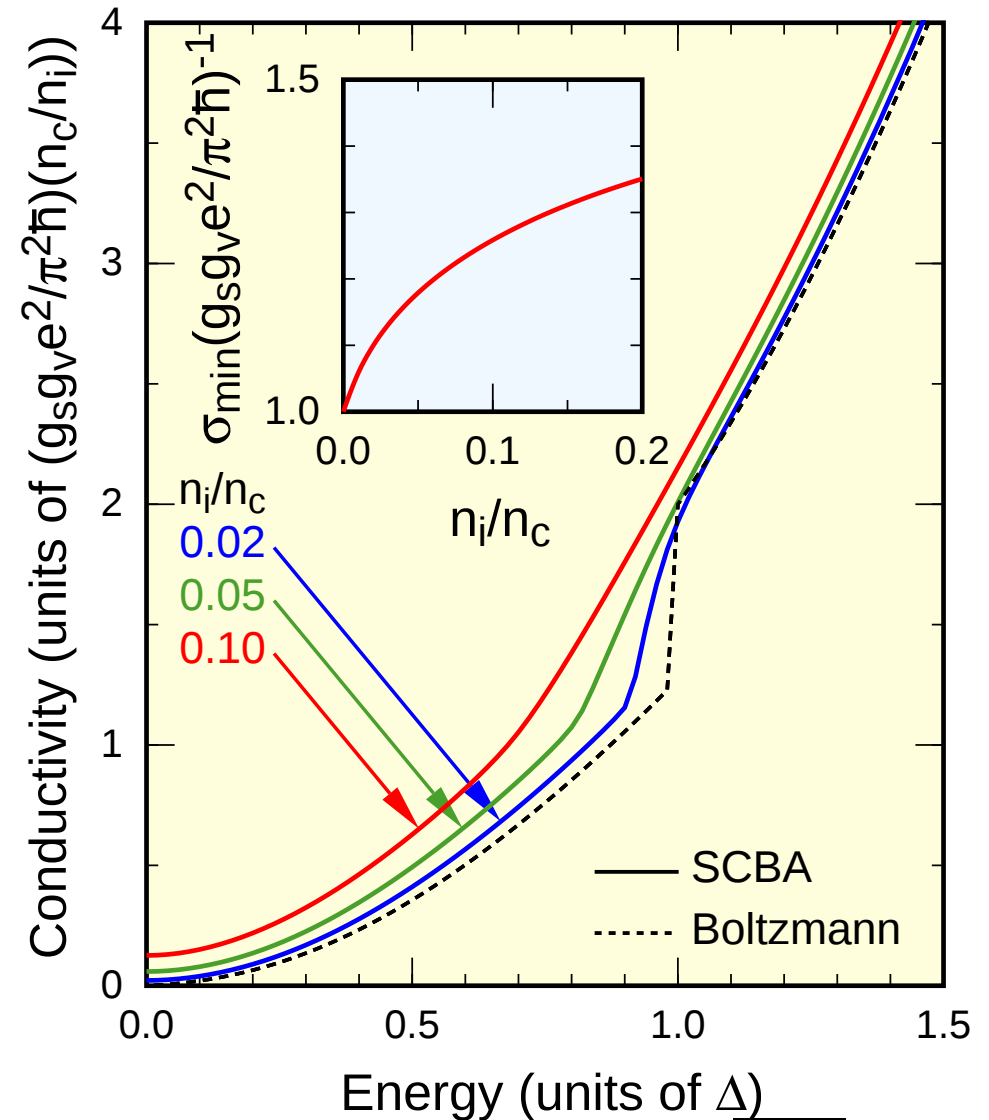
Charged Impurities (Self-Consistent Born Approximation)

[T. Ando, *J. Phys. Soc. Jpn.* 80, 014707 (2011)]

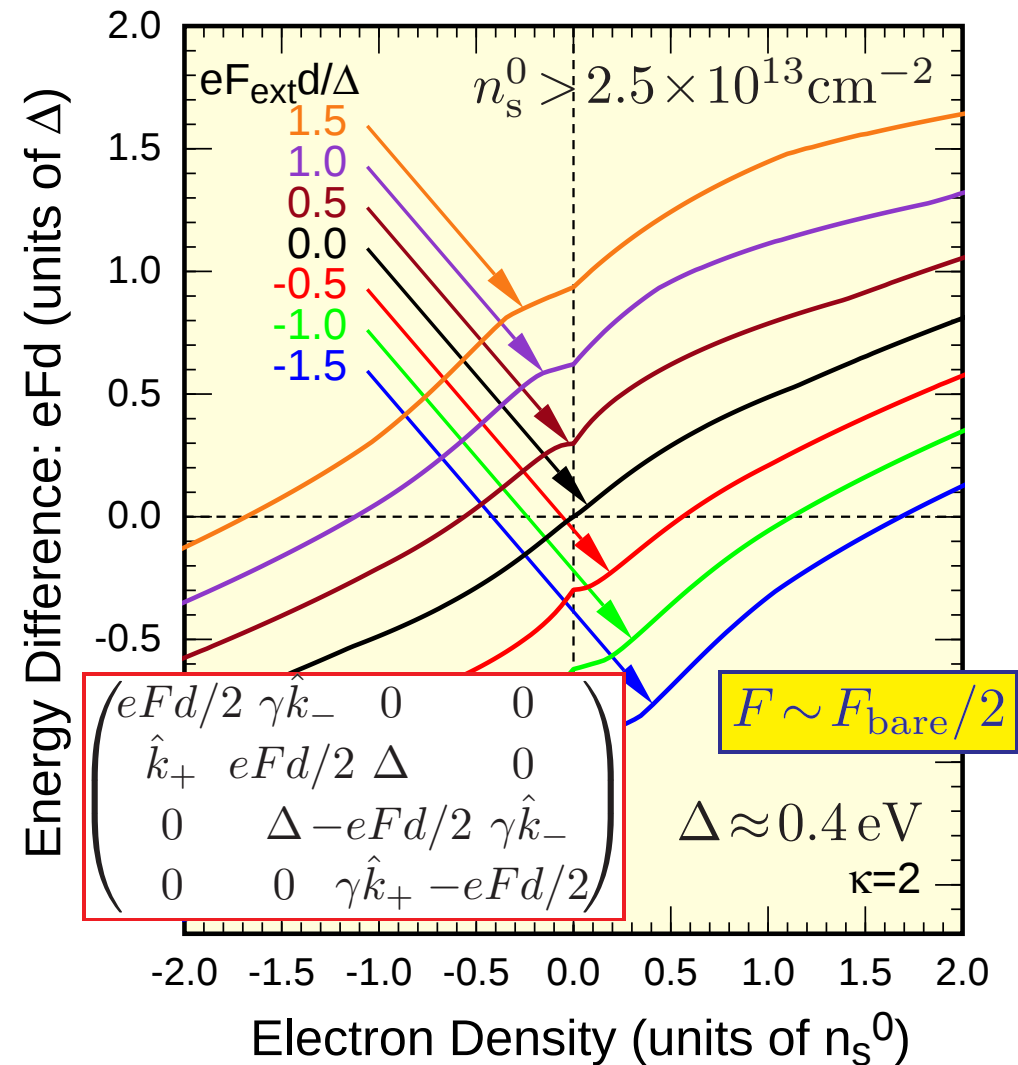
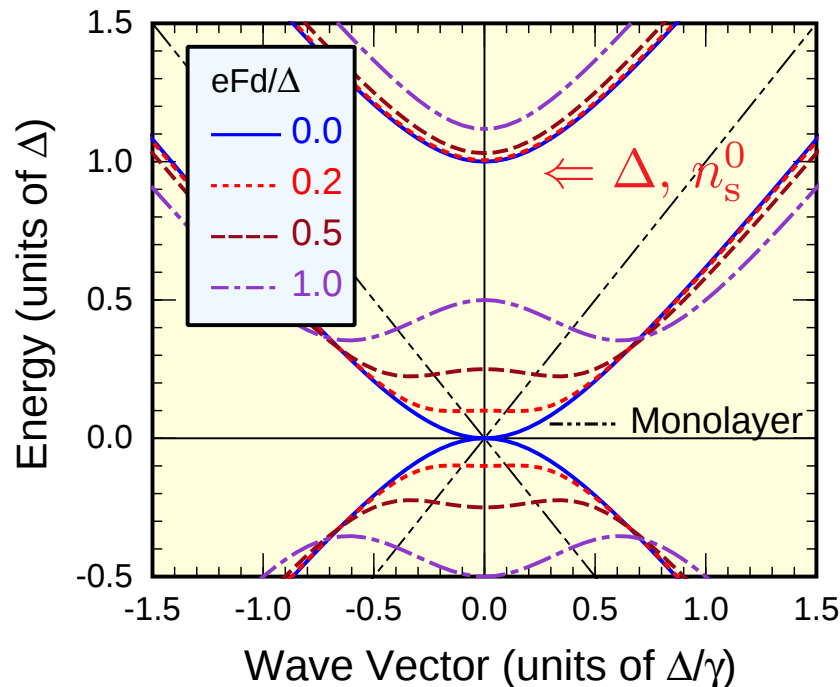
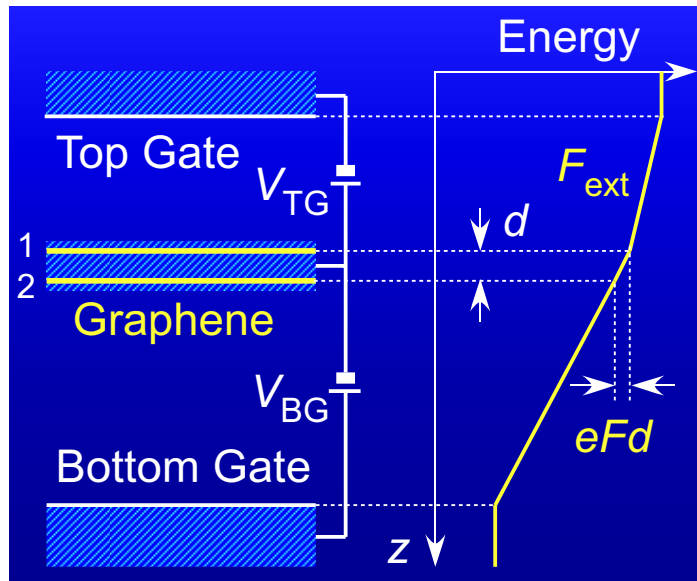
Density of states



Conductivity



Field Effect on Energy Spectrum in Bilayer Graphene



E. McCann, PRB 74, 161403 (2006)

H. Min et al., PRB 75, 155115 (2007)

T. Ando & M. Koshino, JPSJ 78, 034709 (2009)

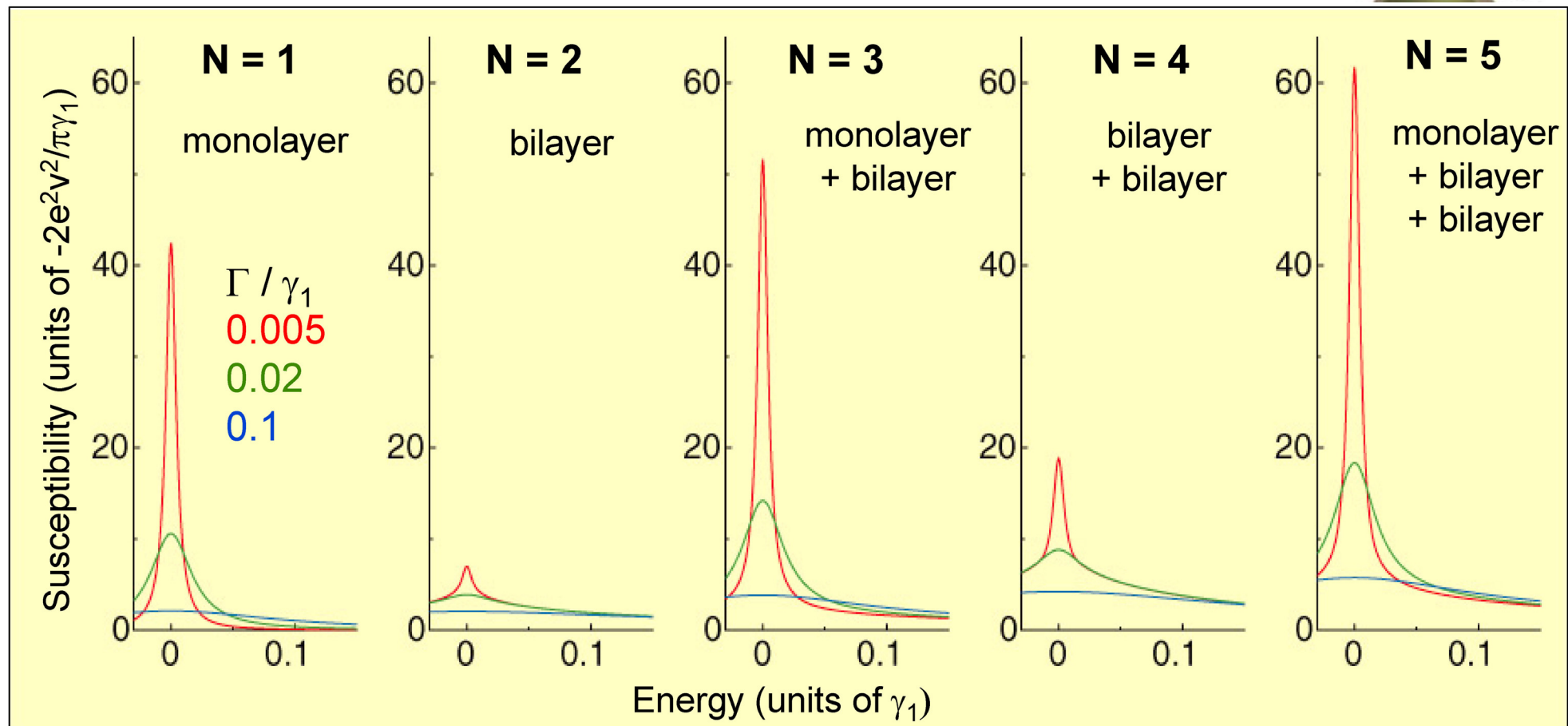
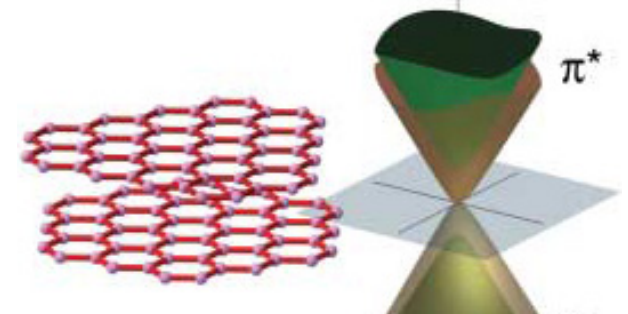
Multi-Layer Graphene [*M. Koshino & T. Ando, PRB 76, 085425 (2007)*]

Exact decomposition of effective Hamiltonian

$2M+1$ Layers = $1 \times$ Monolayer + $M \times$ Bilayers

$2M$ Layers = $0 \times$ Monolayer + $M \times$ Bilayers

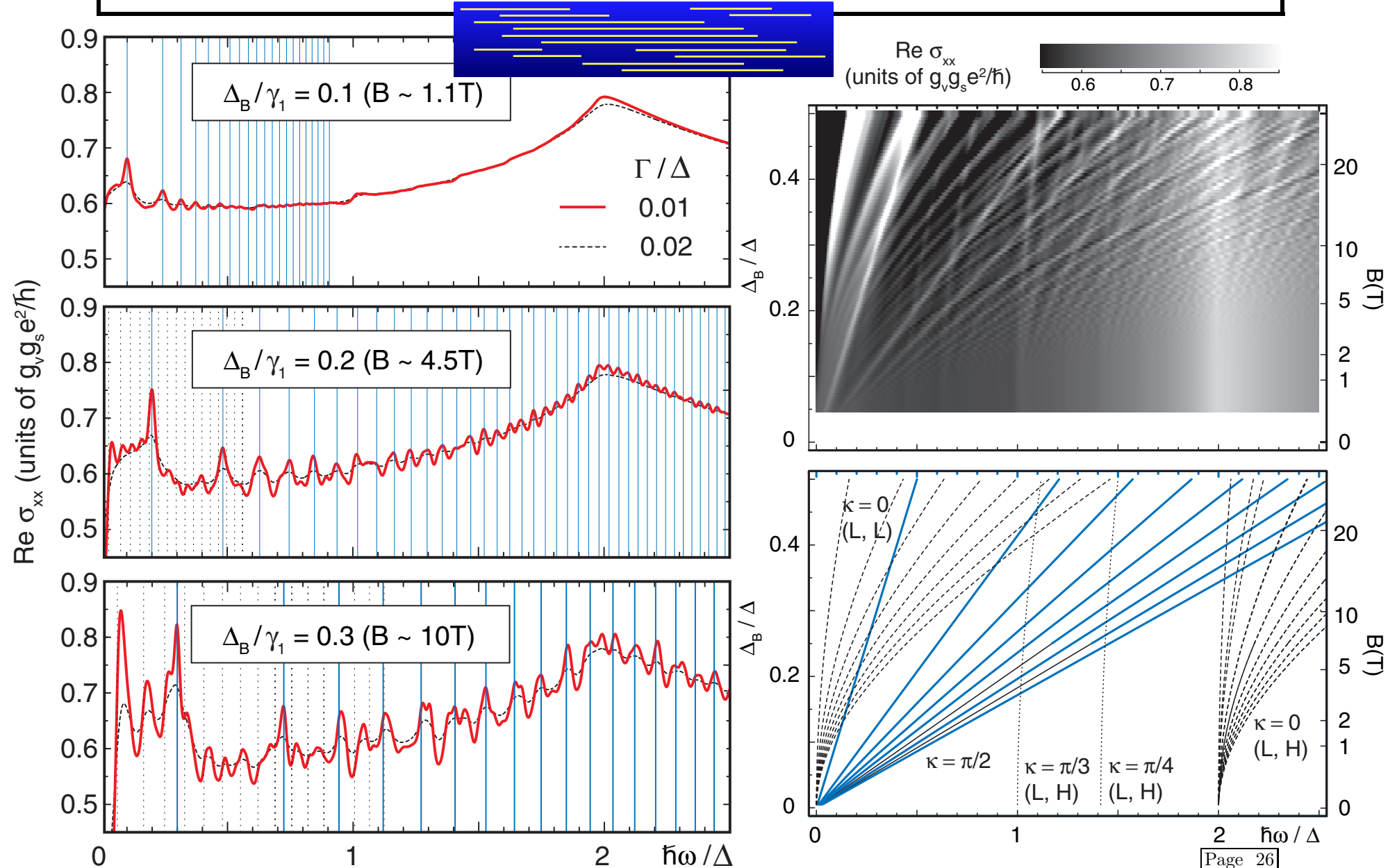
Three parameters: γ_0 , γ_1 , γ_3 (trigonal warping)



Diamagnetic susceptibility

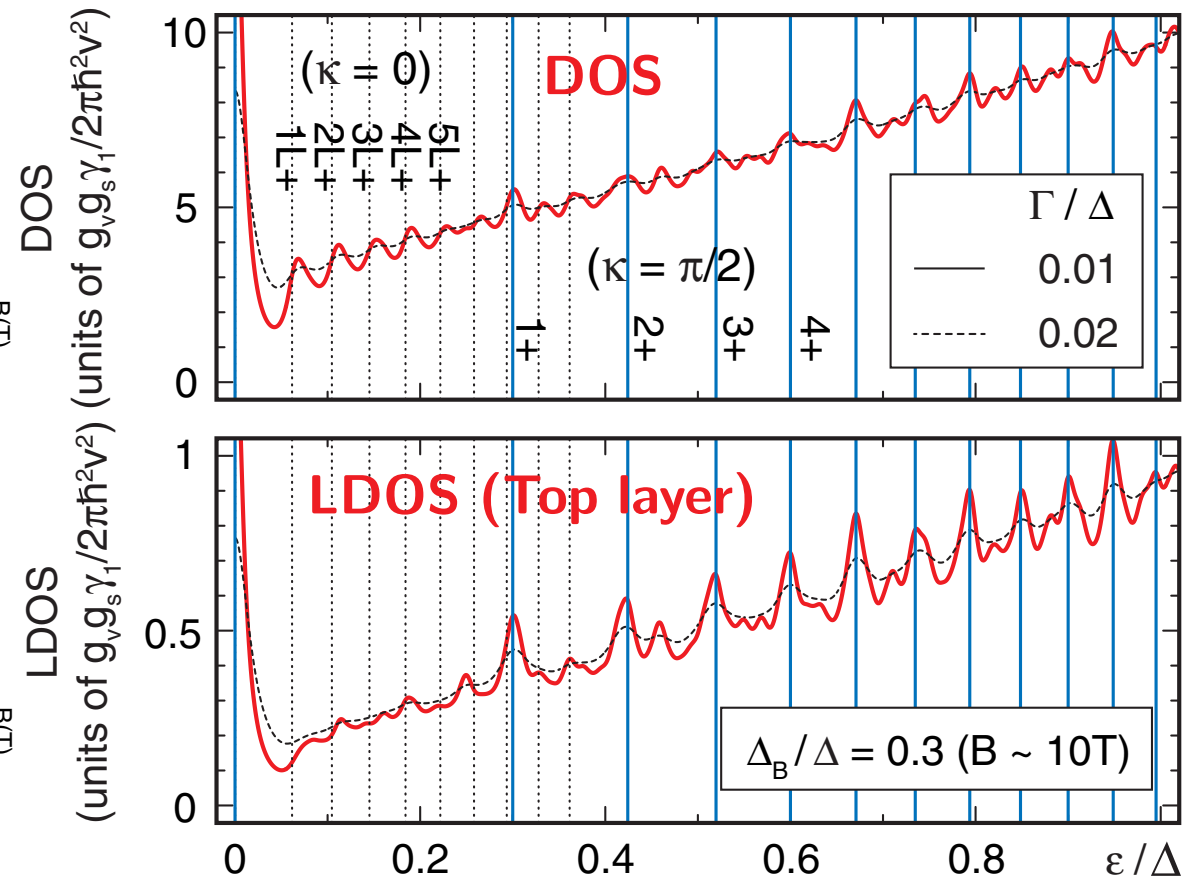
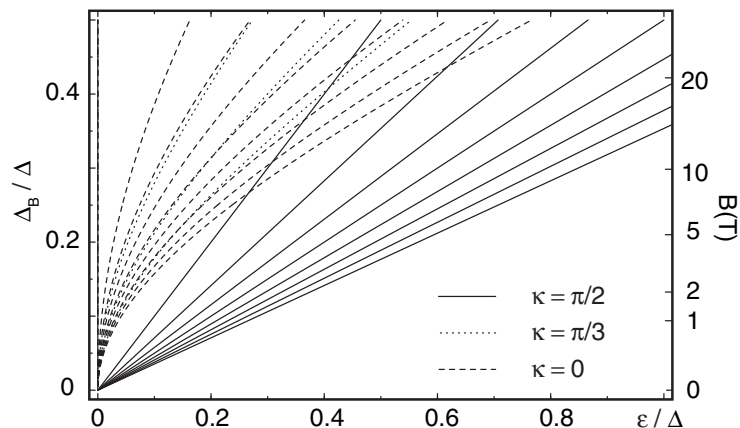
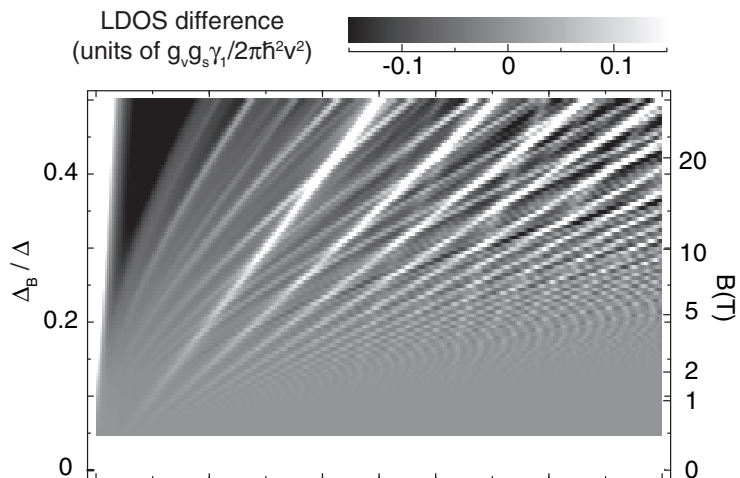
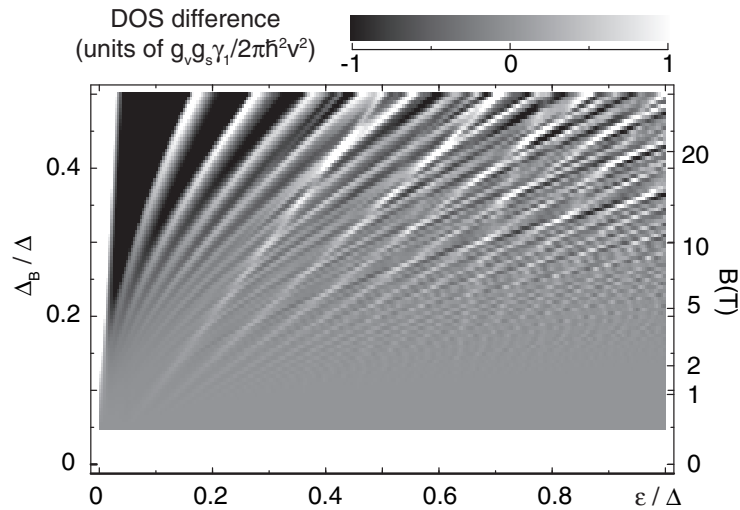
Dynamical Conductivity of Multi-Layer Graphene (Average 1–20)

M. Koshino and T. Ando, Phys. Rev. B 77, 115313 (2008)



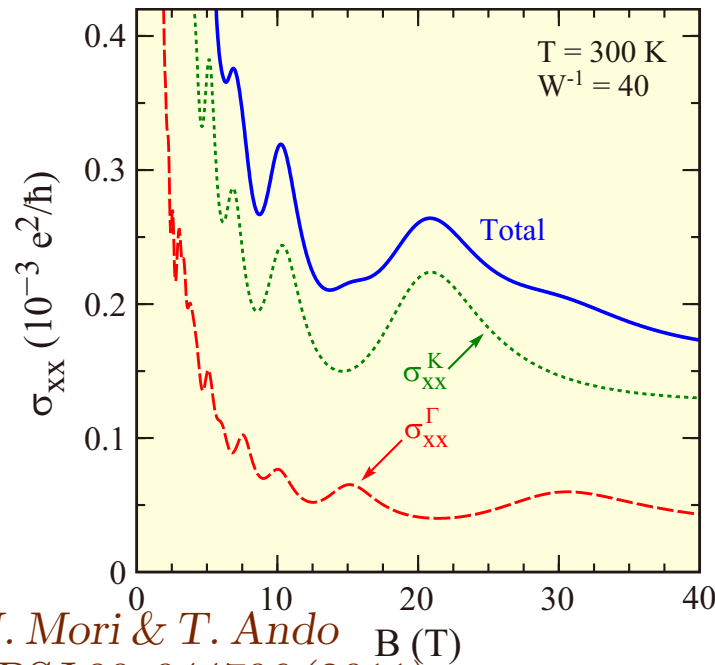
Local Density of States of Multi-Layer Graphene (Average 1–20)

M. Koshino and T. Ando,
Phys. Rev. B 77, 115313 (2008)

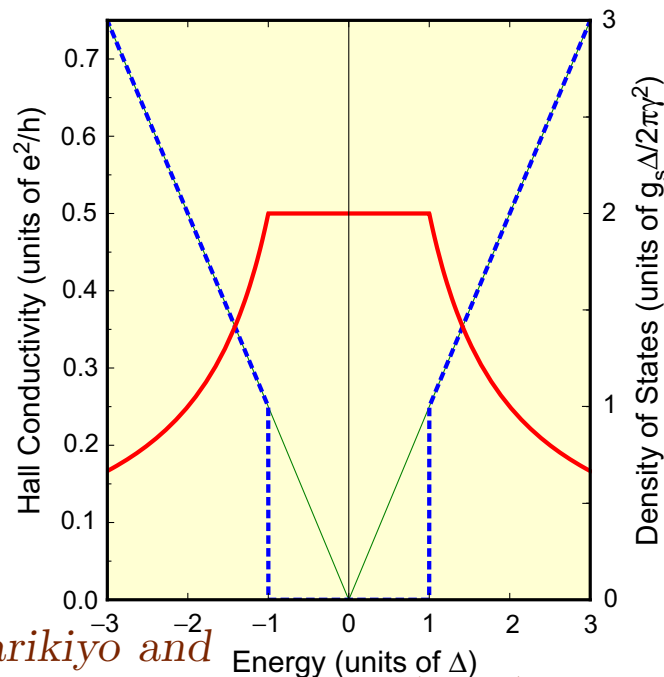


$\Delta_B / \Delta = 0.3$ ($B \sim 10T$)

Further Developments and Future Outlook



N. Mori & T. Ando
JPSJ 80, 044706 (2011)



O. Narikiyo and
K. Kuboki, JPSJ 62, 1812 (1993)

1. Time reversal symmetry and crossover
2. Electron-phonon interaction, Gauge field, instabilities, ...

- Optical phonon anomaly

- Pseudo Landau levels

⇐ Magnetophonon resonances

3. Electron-electron interaction

4. Band gap and valley polarization

- Substrate effects

- Orbital magnetism

⇐ Chiral anomaly (Hall effect without B , ...)

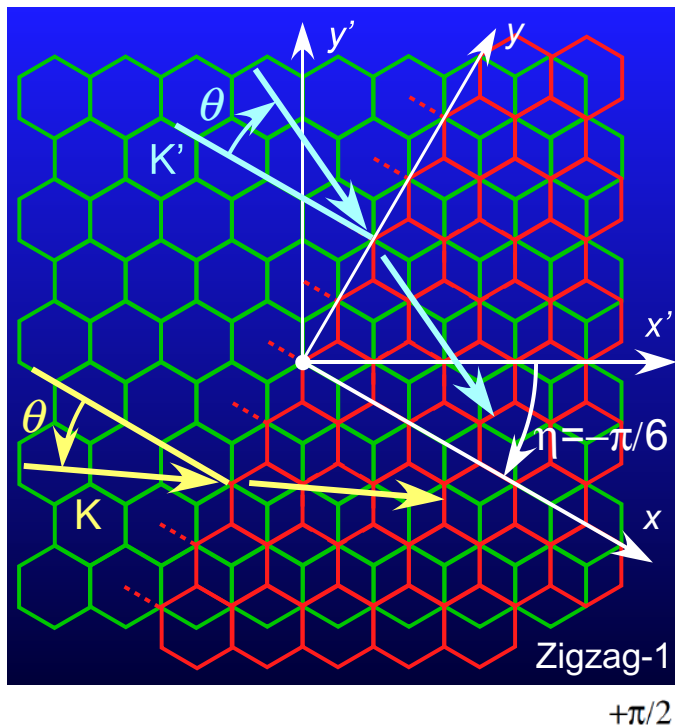
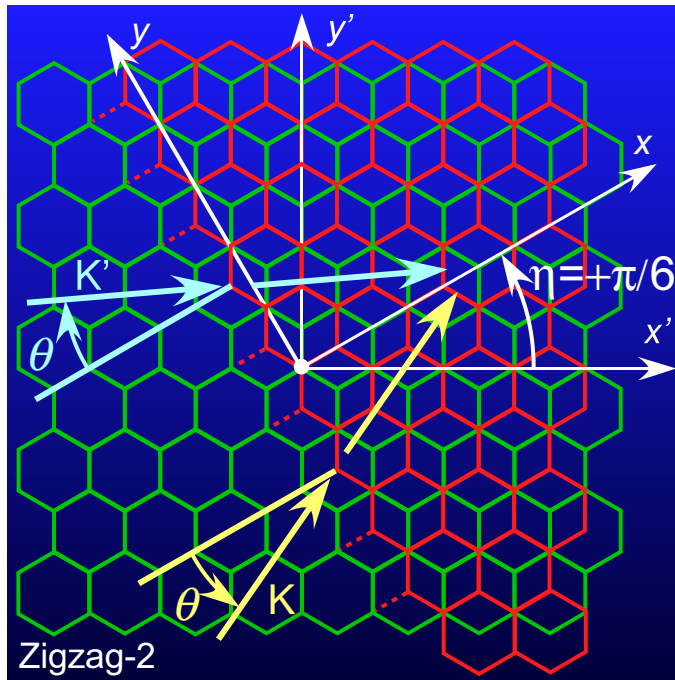
- Chiral edge states

- Exotic phenomena

5. Graphene ribbons

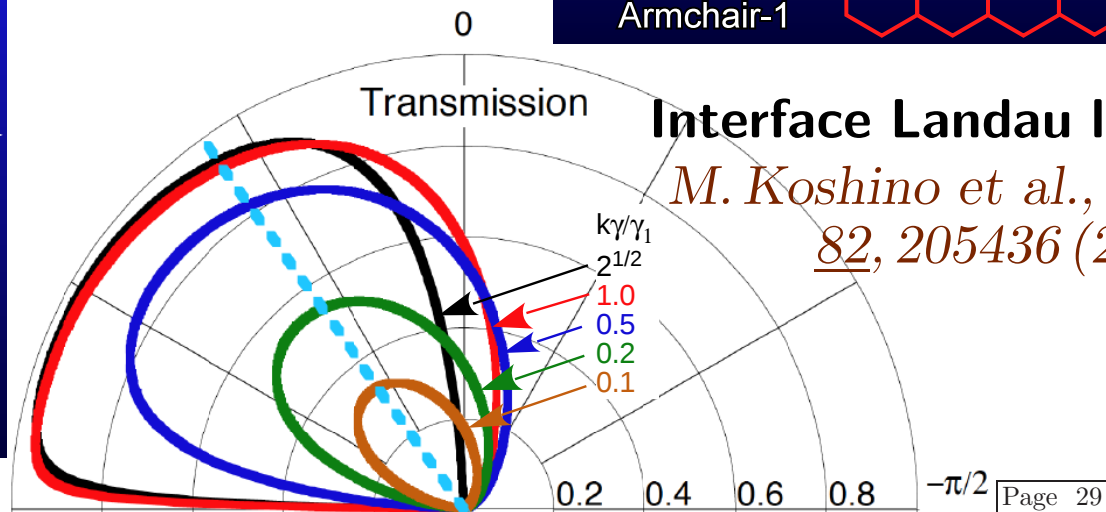
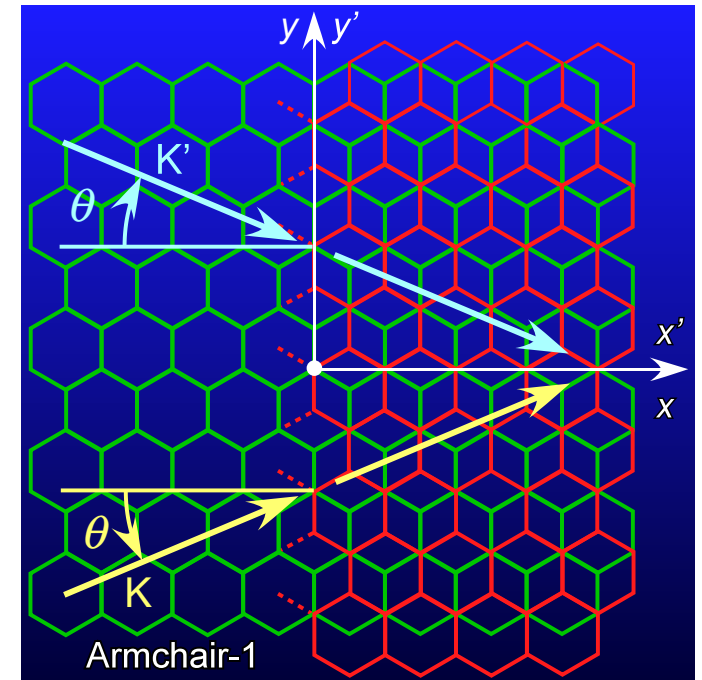
- Chiral edge states

Transmission between Mono- and Bi-Layer Graphene [T. Nakanishi, M. Koshino, and T. Ando, *PRB* 82, 125428 (2010)]



Valley polarization

$$T_K(\theta) \gg T_{K'}(\theta) \quad (\theta > 0)$$



Interface Landau levels

M. Koshino et al., *PRB* 82, 205436 (2010)

Summary: Electronic and Transport Properties of Graphene

Tsuneya ANDO

1. Introduction

- Weyl's equation for neutrino
- Berry's phase and topological anomaly

2. Singular diamagnetic susceptibility

- Band-gap effect
- Spatially varying magnetic field

3. Transport properties graphene

- Singularities at Dirac point
- Long-range scatterers

4. Multi-layer graphene

- Bilayer graphene
- Hamiltonian decomposition

5. Future outlook

- Gauge fields
- Chiral edge states

6. Summary

Nara, Jun 4 (Tue) 2013



Collaborators

Takeshi NAKANISHI (AIST)
Masaki NORO (TITech)
Mikito KOSHINO (Tohoku Univ)



16th International Workshop on Computational Electronics, Nara Prefectural New Public Hall
 Jun 4 (Tue) – 7 (Fri) 2013 [13:30–14:30 (50+10)]