# Performance investigation of uniaxially strained phosphorene n-MOSFETs

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Abstract-In this work, we report on the performance of uniaxially strained monolayer black phosphorus (phosphorene) n-MOSFETs by first-principles-based quantum transport simulations. First, without applying strain, armchair direction (AD) FETs show better performance than zigzag direction (ZD) FETs due to the orientation-dependent transport characteristics that are attributed to the anisotropic band structure of phosphorene. With increasing tensile strain in AD, however, there occurs a downward shift of the conduction band (CB) at Y where the electron effective mass is light (heavy) in ZD (AD), resulting in an increase (decrease) in ON-state current (IoN) of ZD (AD) FET. For the case of tensile strain in ZD, a sharp increase (decrease) in Ion of ZD (AD) FET was observed. This is mainly due to band switching between the first and second conduction bands and a CB at a point between  $\Gamma$  and X; CB here has an electron effective mass light enough to enhance the performance of ZD FET. Overall, compared to the compressive strain, the band structure of phosphorene is more sensitive to the tensile strain in both directions, making it possible to improve the performance of ZD FET up to the level of AD FET.

Keywords—Black phosphorus, anisotropic band strucutre, strain, n-MOSFETs.

# I. INTRODUCTION

In recent decades, mainly because of their distinctive electronic and optical properties, 2D materials have drawn extensive interest and have become promising candidates for future electronics applications [1], [2]. Among 2D materials, black phosphorus (BP) has been widely explored because it displays electronic properties superior to those of its namely metal competitors, graphene and transition dichalcogenides (TMDs). For example, monolayer BP (phosphorene) inherently possesses a highly anisotropic band structure that arises from its puckered honeycomb structure [3]. Moreover, it exhibits higher mobility than that of TMDs and has a direct band gap that can be easily tuned by the thickness of the structure, which enhance its capability for nanoelectronic applications [4].

It is quite clear that the tunability of the electronic properties of semiconductors is a crucial aspect in terms of applications for nanoscale semiconductor devices. With the ultimate goal of enhancing device performance, strain engineering has generally been used to tailor electronic properties by means of various



Fig. 1. (a) Top view of phosphorene. The dashed rectangle indicates the primitive cell lattice constants,  $a_1$ = 4.52 Å,  $a_2$ = 3.31 Å; we set the strain directions of X and Y for armchair and zigzag directions, respectively. (b) Schematic diagram of the phosphorene n-MOSFETs. (c) 2D Brillouin Zone for phosphorene.

approaches, such as lattice mismatch, functional wrapping [5], [6], material doping [7], [8] and direct mechanical application [9] . Recent studies on strain engineering have suggested that phosphorene shows dramatic variations in electronic properties, such as band gap and carrier effective mass, under various strain conditions [10]. Following the studies dealing with the effects of strain on electronic properties, continuous attempts have been made to analyze the effect of varying strain on the performance of phosphorene devices by carrying out quantum transport simulations. For example, a computational investigation was conducted on the device performance of strained BP based MOSFETs. Transport properties were assessed under a moderate region of strain (-3% to +3%) with the effective mass Hamiltonian [11].

However, the effects of strain on the transport characteristics under higher levels of strain, which the device performance is expected to change more severely, have not been thoroughly studied yet. In this work, we conducted a theoretical study of



Fig. 2. Band structures of phosphorene under uniaxial strain (a)  $\varepsilon_{AD}$  and (b)  $\varepsilon_{ZD}$ . Band structures for the strains of 0%, ±3%, and ±6% are selectively plotted to clarify the changes under strain.

uniaxially strained phosphorene n-type MOSFETs using firstprinciples-based quantum transport simulations, taking into account the full band structure of phosphorene under various strain conditions.

### II. SIMULATION DETAILS

For the geometry optimization, first-principles density functional theory (DFT) calculations were carried out on the basis of a linear combination of atomic orbitals (LCAO) and using the OpenMX code [12]. We chose a  $3s^{1}3p^{3}2d^{5}1f^{7}$  LCAO basis set and generalized gradient approximation (GGA) exchange-correlation functional by Perdew, Burke, and Ernzerhof was adopted [13]. We employed a  $10 \times 10 \times 1 k$ -grids and the atomic structures of strained phosphorene were fully relaxed until the maximum force reached 10<sup>-4</sup> Hartree/Bohr. Fig. 1(a) presents the calculated lattice constants for phosphorene which are  $a_1 = 4.52$  Å and  $a_2 = 3.31$  Å, and in good agreement with the values of other theoretical works [4], [14]. With a fully relaxed structure of phosphorene, strain was gradually applied in either the AD or ZD, within the range of  $\pm 6\%$ ; these strains are denoted as  $\varepsilon_{AD}$  and  $\varepsilon_{ZD}$ , respectively. Especially when strain was applied in ZD, a negative Poisson ratio was observed, which originated from the puckered structure; this result is in-line with those of other theoretical studies [15], [16].

After validating our calculation results, non-orthogonal DFT Hamiltonians were extracted for each strained structure. Note that, to achieve the computational efficiency, the sizes of the Hamiltonians were reduced by a mode-space transformation, as presented in [17]. Next, to analyze the transport characteristics of phosphorene n-MOSFETs, we performed quantum transport calculations; a schematic diagram of our results is provided in Fig. 1(b). The channel length ( $L_{ch}$ ) and equivalent oxide thickness (oxide thickness ( $T_{ox}$ ) is set to 1.0 nm with relative dielectric constant of 1.12 for phosphorene [18]) are 10 nm and



Fig. 3. (a) Band gap is plotted as a function of both  $\varepsilon_{AD}$  and  $\varepsilon_{ZD}$ .  $m_e^*$  under strain (b)  $\varepsilon_{AD}$  and (c)  $\varepsilon_{ZD}$ .  $\Gamma$ -AD (or ZD) denotes  $m_e^*$  of the CB at  $\Gamma$  for each direction and similar notation holds for CB at  $\Gamma$  and  $\chi$  (Y-AD and  $\chi$ -AD).

0.5 nm, respectively. Source and drain are n-doped with a density of  $1 \times 10^{21}$  cm<sup>-3</sup>. DFT Hamiltonians of reduced size were imported into our in-house tool to compute the current and charge densities by self-consistently solving the non-equilibrium Green's function and the Poisson equation [17].

#### **III. BAND STRUCTURE OF STRAINED PHOSPHORENE**

Fig. 2(a) and (b) illustrate the anisotropic band structures of phosphorene under two different strains,  $\varepsilon_{AD}$  and  $\varepsilon_{ZD}$ , respectively. The band structure was described along the band path of Y- $\Gamma$ -X, as shown in Fig. 1(c). For unstrained phosphorene, there is a direct band gap, evaluated at 0.92 eV, which is in agreement with other *ab-initio* studies [4], [19]. As is already widely known, DFT-GGA functional underestimates the band gap of semiconductors: experimentally reported value of band gap is 1.45 eV [20]. Even with this underestimation, however, it is expected that our DFT results are meaningful because we are focused on the general trends of the effects of strain on the band structure and transport properties.

The band gap of phosphorene is plotted in Fig. 3(a) as a function of both  $\varepsilon_{AD}$  and  $\varepsilon_{ZD}$ . First, the band gap shows a gradual decrease with the increase of compressive  $\varepsilon_{AD}$ ; this decrease of band gap is attributed to a downward shift of CB at  $\Gamma$ . A similar tendency of decreasing band gap can also be seen under compressive  $\varepsilon_{ZD}$ . With increasing tensile strain, the band gap



Fig. 4. Transfer characteristics of phosphorene n-MOSFETs under various strain conditions: (a)  $\varepsilon_{AD}$ -AD FET, (b)  $\varepsilon_{AD}$ -ZD FET, (c)  $\varepsilon_{ZD}$ -AD FET, and (d)  $\varepsilon_{ZD}$ -ZD FET. Note that the subscript of  $\varepsilon_{AD}$  (or  $\varepsilon_{ZD}$ ) refers to the direction of applied strain, whereas AD (ZD) FET refers to the phosphorene n-MOSFETs whose transport direction is AD (ZD).

monotonically increases under  $\varepsilon_{AD}$  as CB moves away from the Fermi level, whereas with tensile  $\varepsilon_{ZD}$ , the band gap starts to decline after its peak around  $\varepsilon_{ZD} = +4\%$  due to band switching between the first and second conduction bands.

The effective masses of electrons  $(m_e^*)$  are obtained from the curvature of the band structure. Fig. 3(b) and (c) show the alterations of  $m_e^*$  under  $\varepsilon_{AD}$  and  $\varepsilon_{ZD}$ , respectively. Regarding compressive strain, no significant change in  $m_e^*$  is observed under either  $\varepsilon_{AD}$  or  $\varepsilon_{ZD}$ , indicating the orientation dependence of phosphorene where  $m_e^*$  along ZD is larger than that along AD [21]. Under tensile  $\varepsilon_{AD}$ , however, there occurs a downward shift of CB at Y, for which  $m_e^*$  is light (heavy) in ZD (AD). For instance, at  $\varepsilon_{AD} = +6\%$ , the energy gap between the conduction band minimum (CBM) at  $\Gamma$  and the conduction band edge (CBE) at Y narrows to 0.062 eV. For tensile  $\varepsilon_{ZD}$ , band switching takes place around  $\varepsilon_{ZD} = +4\%$ , bringing sudden changes of  $m_e^*$  in both AD and ZD. With further increased tensile  $\varepsilon_{ZD}$ , CB at k = (0.31, 0.0, 0.0), denoted as  $\chi$ , with small  $m_e^*$  in both AD and ZD, changes notably. For instance, the value of CBE at  $\gamma$  becomes comparable to that of CBM at  $\varepsilon_{ZD} = +6\%$ , as plotted in Fig. 3(c).

# IV. PHOSPHORENE n-MOSFETS SIMULATION RESULTS

Fig. 4 shows the transfer curves of uniaxially strained phosphorene n-MOSFETs. Curves are shifted to have an OFF-state current ( $I_{OFF}$ ) of  $10^{-1} \,\mu A/\mu m$  at  $V_{GS} = 0$  V, which is a ITRS requirement for high performance in the  $I_{OFF}$  specifications [22]. Transfer curves for strains of 0%, ±3%, and ±6% are selectively plotted. Although  $I_{ON}$  rarely changes in the first two cases, an appreciable decrease and an appreciable increase are observed from AD FET and ZD FET, respectively under strong tensile  $\varepsilon_{ZD}$ , as can be seen in Fig. 4(c) and (d).



Fig. 5. *IoN* of the n-MOSFET as a function of uniaxial strain. *IoN* is defined at  $V_{GS}$ - $V_{th} = V_{DS} = 0.5$  V, where the OFF-state current is set at  $10^{-1} \,\mu\text{A}/\mu\text{m}$ 

To clarify the variations in transport characteristics, the values of  $I_{ON}$  of n-MOSFETs with two different transport directions, AD FET and ZD FET, as a function of strains  $\varepsilon_{AD}$  and  $\varepsilon_{ZD}$ , are shown in Fig. 5. Here,  $I_{ON}$  is defined at  $V_{GS} - V_{th} = V_{DS} = 0.5$  V. As can be expected from the results in section III, no noticeable change of  $I_{ON}$  occurs for either compressive  $\varepsilon_{AD}$  or  $\varepsilon_{ZD}$ , since there was little change in  $m_e^*$  (see Fig. 3(b) and (c)).

However, it is found that AD FET under strong tensile  $\varepsilon_{AD}$ -AD FET) shows a rapid decline in  $I_{ON}$  because a CB at Y, whose  $m_e^*$  is heavy in AD, begins to contribute to  $I_{ON}$  at  $\varepsilon_{ZD} = +4\%$ . On the contrary, ZD FET under the same strain condition ( $\varepsilon_{AD}$ -ZD FET) shows a moderate improvement in  $I_{ON}$  owing to the small value of  $m_e^*$  at Y in ZD. For the case of tensile  $\varepsilon_{ZD}$ , there occurs a sharp increase (decrease) in  $I_{ON}$  of  $\varepsilon_{ZD}$ -ZD (AD) FET around  $\varepsilon_{ZD} = +4\%$ , which can be explained by the decrease (increase) of  $m_e^*$  along ZD (AD), resulting from band switching. Particularly,  $\varepsilon_{ZD}$ -ZD FET under strong tensile  $\varepsilon_{ZD}$  exhibits the greatest improvement in  $I_{ON}$ ; this improvement is attributed to not only band switching but also to the CB at  $\chi$ , with light  $m_e^*$ , which contributes dominantly to improving  $I_{ON}$ .

#### V. CONCLUSION

A detailed study of strain effect on uniaxially strained phosphorene n-MOSFETs was performed by means of DFTbased self-consistent quantum transport simulations. We identified the variations of band structures with various strains, and found that these strain effects directly lead to alterations in the performance of phosphorene FETs. Overall, it is likely that tensile strain applied in either direction is more useful than compressive strain in modifying the electronic properties of phosphorene. Consequently, we emphasize that the strain engineering can be an effective way to modulate the effective mass of phosphorene, making it possible to narrow the performance gap between AD and ZD FETs.

# ACKNOWLEDGMENT

This research was supported by the Development of Million-Atom First-principles Simulator for Nano Transistor Devices through the National Research Foundation of Korea (NRF) funded by the Ministry of Education, Science and Technology (Grant No.NRF-2017R1A2B2005679).

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