

Tight-binding Study of Γ -L Bandstructure Engineering for Ballistic III-V nMOSFETs

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Abstract— A major concern for III-V nMOSFETs is the degradation of I_{ON} due to low density of states and spillover of the charge from high-mobility Γ -valley to low-mobility L-valley at high sheet charge density. In this paper, we study these Γ -L bandstructure effects for ultrathin-body $In_xGa_{1-x}Sb$ nMOSFETs with varying stoichiometry using tight-binding and ballistic transport model.

Keywords- tight-binding, ballistic transport, ultra-thin body, *InGaSb*.

I. INTRODUCTION

nMOSFETs based on III-V materials have the highest mobility/injection velocity (v_{inj}), the major concern however, is the degradation of device performance due to low density-of-states (DOS) (low effective mass of carriers) and spillover of the charge from high-mobility Γ - to low-mobility L-valley at high sheet charge [1-3]. In this paper, we study these Γ -L bandstructure effects for varying stoichiometry of $In_xGa_{1-x}Sb$, which has high mobility for both electrons [4] & holes [5] and is a promising candidate for future technology nodes.

For performance evaluation of ultrathin-body (UTB) double-gate devices, the use of bulk effective masses is not adequate [6]. E-k relations for $In_xGa_{1-x}Sb$ UTB MOSFET (Fig. 1) are calculated using $sp^3d^5s^*$ atomistic tight-binding (TB) model coupled with Poisson's equation. The effect of varying the In % on DOS, electron population among Γ -, L-, and X-valleys is studied systematically. E-k band diagrams for UTB MOSFET with GaSb and InSb channel are plotted in Fig. 2. Amongst different valleys of electrons, Γ -valley has lower effective mass, thus higher v_{inj} but lower DOS. For GaSb and low In% $In_xGa_{1-x}Sb$, due to low energy separation between Γ - and L-valleys ($\Delta_{\Gamma-L}$), L-valley can also be populated, which has higher DOS but lower v_{inj} . Increasing In% in the compound brings up L-valley, reduces the energy for Γ -valley, meanwhile reducing the effective mass of Γ -valley. However, high In% makes it difficult to achieve high electron sheet charge density (N_S) from Γ -valley (due to the low DOS). For high drive current, high v_{inj} and N_S have to be achieved simultaneously, which requires engineering of DOS, v_{inj} and electrons population in different valleys. We study the effect of all these factors (Fig. 3): The band energy (Fig. 5) & difference in DOS (Fig. 6) among Γ - and L-valleys determines the overall population among different valleys (Fig. 7). Using N_S (Fig. 8) and v_{inj} (Fig. 9), ballistic drive current (Fig. 10) is calculated and compared for varying In %.

II. METHODOLOGY

TB parameters for ternary $In_xGa_{1-x}Sb$ are calculated following virtual-crystal approximation (VCA) incorporating compositional disorder effect and fitted to bulk band gap of ternary compound [7-8]. 1D Poisson's equation perpendicular to channel direction is coupled with TB Hamiltonian by Hartree-Fock potential in the gate stack. Dangling bonds at interface are pacified by hydrogen termination of hybridized orbitals to eliminate all the states within band gap [9]. A ballistic transport model is adopted to assess transport of electrons [10]. v_{inj} is determined from full band structure with non-parabolic E-k relationship considered for all valleys.

III. SIMULATION RESULTS

A. Band energy, DOS and Valley Population

Band structures of $In_xGa_{1-x}Sb$ for $T_{BODY}=4nm$ at N_S of $\sim 3 \times 10^{12} cm^{-2}$ for different In %'s are compared in Fig. 4. For low In % $In_xGa_{1-x}Sb$, because of quantum confinement (QC) effects, $\Delta_{\Gamma-L}$ is marginal, especially under high V_G and thin T_{BODY} , resulting in Fermi level moving into L-valley. Band gap (E_g) & $\Delta_{\Gamma-L}$ are shown in Fig. 5: with higher In%, $\Delta_{\Gamma-L}$ is increased (from ~ 0 to $\sim 0.7eV$) to confine more electrons in Γ -valley counteracting the effects of quantization (Fig. 5(a)), which is more dominant for thin T_{BODY} (Fig. 5(b)) ($\Delta_{\Gamma-L} \sim 0.5eV$ for $T_{BODY}=7nm$, $\sim 0.43eV$ for $T_{BODY}=3nm$, In % 0.5). At higher N_S (V_G), though E_g is lowered due to quantum confinement stark effect [11] (from ~ 0.5 to $\sim 0.35eV$ for $T_{BODY}=4nm$, In % 0.5), $\Delta_{\Gamma-L}$ stays low (Fig. 5(c)) ($\sim 0.03eV$ for $T_{BODY}=4nm$, GaSb). This can be attributed to Γ -valley's curvature getting blunt (effective mass becomes larger) at high N_S , thus alleviated quantization effect. Fig. 6 plots the 2-D DOS, Γ -valley has 100x lower DOS compared with L-valley, which is required to achieve high N_S . In-rich compounds lead to decrease in DOS at conduction band edge ($\sim 10^{15} eV^{-1}cm^{-2}$ GaSb, $\sim 10^{13} eV^{-1}cm^{-2}$ InSb) meaning further movement of Fermi level (read higher V_G) is necessary to achieve same N_S . % occupation of electrons in Γ -valley is plotted against In %, T_{BODY} and N_S in Fig. 7. From $\Delta_{\Gamma-L}$ for reasonable In %, adequate percentage of charge can be confined in Γ -valley ($\sim 100\%$ at $1 \times 10^{11} cm^{-2}$ for In % 0.5) even at high N_S ($\sim 60\%$ for $4 \times 10^{12} cm^{-2}$, In % 0.5) and thin T_{BODY} ($\sim 50\%$ for $T_{BODY}=3nm$, In % 0.5). Fig. 8 illustrates (a) sheet charge density as a function of gate voltage and (b) sub-threshold swing comparison. Two slopes in N_S can be identified which correspond to Γ - and L-valleys respectively. For higher In

composition, loss in DOS at band edge requires much higher gate voltage to obtain sheet charge density of reasonable level for device operation (to get $2 \times 10^{12} \text{cm}^{-2}$ GaSb $\sim 0.4\text{V}$, InSb $\sim 1.3\text{V}$). The change in DOS as well as dielectric constant (from 14.4 GaSb to 16.8 InSb) is also reflected in the degradation of sub-threshold swing with increasing In % (100mV/dec for GaSb, 115mV/dec for InSb, $T_{\text{BODY}} = 4\text{nm}$). For device geometry of interest, DOS determined quantum capacitance has an evident impact on the subthreshold behavior as subthreshold swing improves for larger T_{BODY} (95mV/dec for $T_{\text{BODY}} = 4\text{nm}$, 100mV/dec for $T_{\text{BODY}} = 3\text{nm}$, GaSb).

B. v_{inj}

$I_{\text{dsat}} - V_G$ is evaluated by integrating N_S with average velocity of electrons along transport direction at each k point. Parasitic resistance is neglected in the calculation. $\langle 100 \rangle$ is set as the transport direction. The average velocity at given gate bias can be calculated by taking the ratio between the overall current density and sheet charge density [10]. Under the ballistic transport model, when Fermi level is below conduction band edge, the injection velocity stays constant; it increases as Fermi level moves into conduction band. As shown in Fig. 9, since electron population is mostly in L-valley, the overall v_{inj} of GaSb and Ga-rich $\text{In}_x\text{Ga}_{1-x}\text{Sb}$ is low ($\sim 2 \times 10^7 \text{cm/s}$) at high N_S as most of the electrons are in L-valley. In-rich compounds give high v_{inj} ($\sim 1 \times 10^8 \text{cm/s}$), because of sharper curvature of Γ -valley in calculated E-k relation, however population of L-valley leads to the decrease in v_{inj} at high N_S ($\sim 10^{12} \text{cm}^{-2}$ for In % 0.5). By engineering with $\Delta_{\Gamma-L}$ and DOS for the optimal driving capability, the overall v_{inj} can maintain high ($\sim 1 \times 10^8 \text{cm/s}$) with adequate % of charge in Γ -valley, when DOS of L-valley starts to contribute to N_S .

C. Performance Evaluation

In Fig. 8 (a), the filling of L-valley gives raise in N_S . Without excessive filling of L-valley, average injection velocity maintains high as shown in Fig. 9. Fig. 10 shows proper amount of In percentage ($\sim 25\%$) in $\text{In}_x\text{Ga}_{1-x}\text{Sb}$ can give a overall improvement in drive-current by maintaining high v_{inj} with sufficient amount of charge in Γ -valley. Further increase in In composition leads to significant loss in N_S . Drive current for $\text{In}_{0.25}\text{Ga}_{0.75}\text{Sb}$ is 50% higher than silicon, 30% higher than GaSb (highest DOS) and 120% higher than InSb (highest v_{inj}) at an over-drive voltage of 0.7V. It is shown that varying stoichiometry of III-V compound material allows careful engineering of Γ -L band structure, to achieve optimal trade-off between v_{inj} and DOS.

IV. CONCLUSION

Γ -L bandstructure effects in ultra-thin body III-V nMOSFETs are studied using tight-binding and ballistic transport model. It is shown by varying In % in $\text{In}_x\text{Ga}_{1-x}\text{Sb}$, most of electrons can be kept in the Γ -valley at relevant N_S values to avoid excessive population of electrons in the L-valley (GaSb) or significant loss of charge due to low DOS (InSb), hence achieves the best I_{DSAT} .

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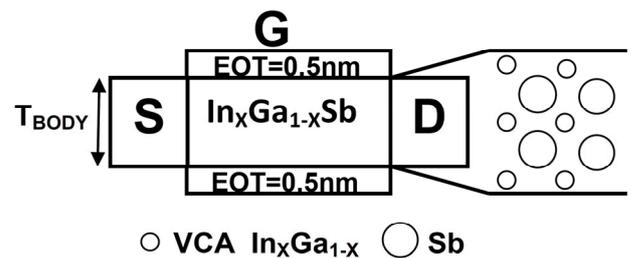


Figure 1. Structure of $\text{In}_x\text{Ga}_{1-x}\text{Sb}$ double-gate MOSFET with (100) orientation. Atom arrangement under VCA

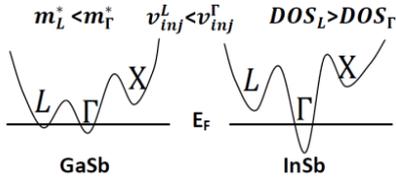


Figure 2. Band diagrams of GaSb and InSb.

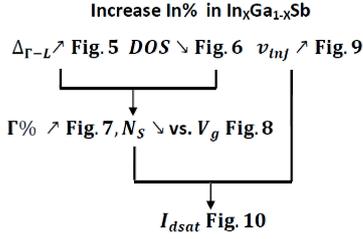


Figure 3. Bandstructure effects by varying In% in $\text{In}_x\text{Ga}_{1-x}\text{Sb}$ and their relationships. Effects studied are plotted in corresponding figures.

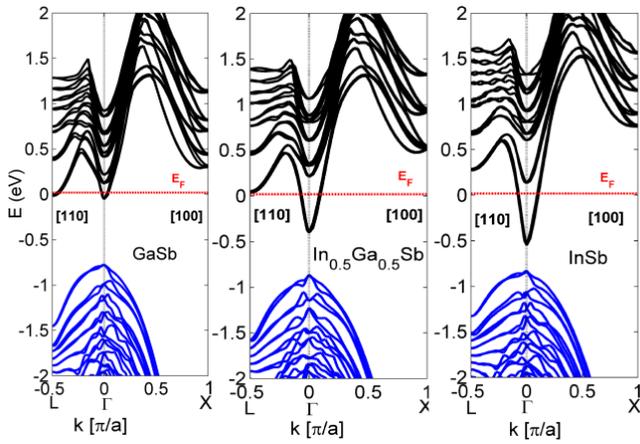


Figure 4. Calculated band structure from TB for $\text{In}_x\text{Ga}_{1-x}\text{Sb}$ at sheet charge density $\sim 3 \times 10^{12} \text{cm}^{-2}$ with a body thickness of 4nm.

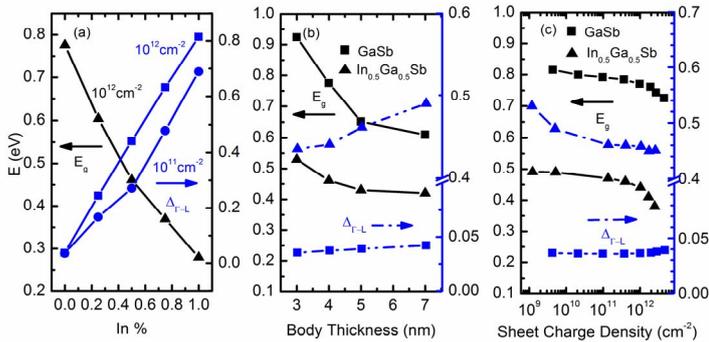


Figure 5. Calculated band gap and Γ -L energy separation vs. (a) In composition for 4nm body thickness and $\sim 10^{11}$ & 10^{12}cm^{-2} sheet charge density; (b) body thickness for sheet charge density $\sim 1 \times 10^{12} \text{cm}^{-2}$; (c) sheet charge density for 4nm body thickness.

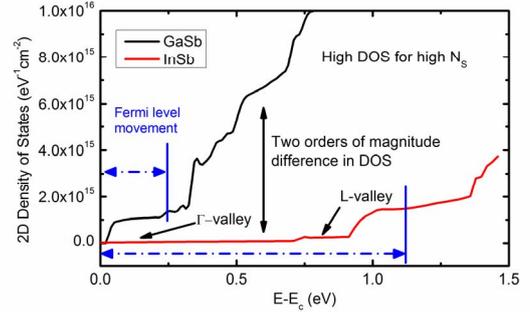


Figure 6. 2D Density of States for GaSb and InSb for body thickness of 4nm.

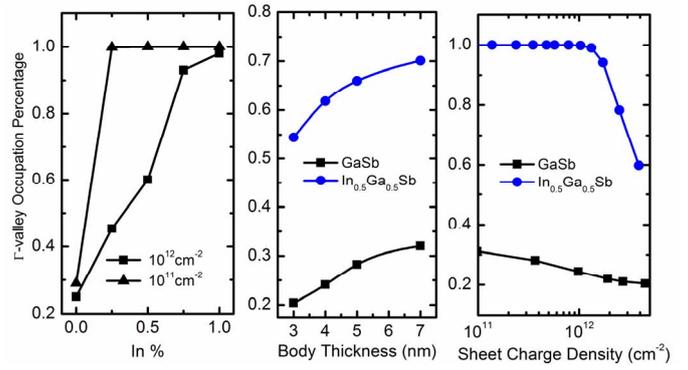


Figure 7. Percentage of electron occupation in Γ -valley vs. (a) In composition for 4nm body thickness and $\sim 10^{11} \text{cm}^{-2}$ & 10^{12}cm^{-2} sheet charge density; (b) body thickness for $\sim 1 \times 10^{12} \text{cm}^{-2}$ sheet charge density; (c) sheet charge density for 4nm body thickness.

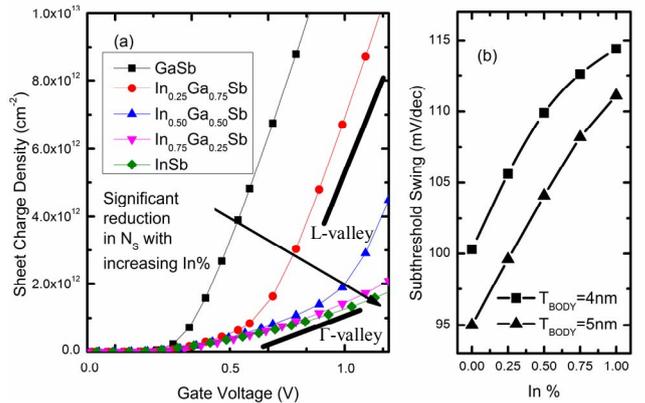


Figure 8. (a) sheet charge density as a function of gate voltage. Gate voltage is adjusted to give $5 \times 10^7 \text{cm}^{-2}$ sheet charge at 0V; (b) subthreshold swing with varying In composition for body thickness of 4nm & 5nm.

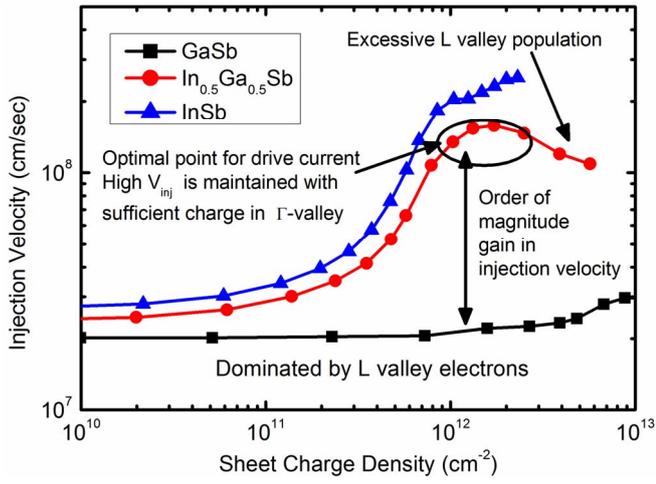


Figure 9. Injection velocity as a function of sheet charge density for different In composition, the drop in injection velocity at high sheet charge is due to L-valley population.

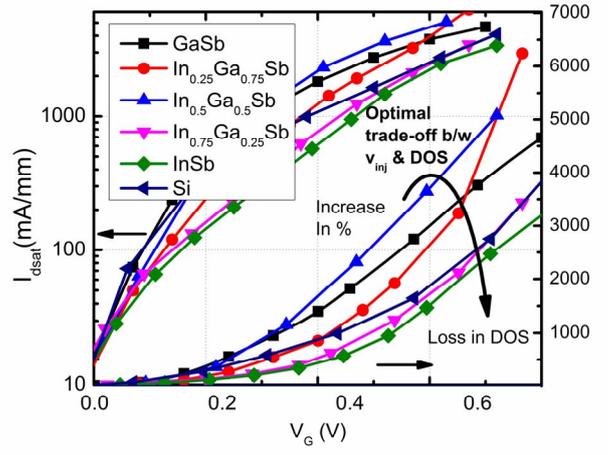


Figure 10. Saturation current as a function of gate voltage; Gate voltage is adjusted to give $15\mu\text{A}/\mu\text{m}$ current density at 0V.