

Full-Band Monte Carlo Device Simulation of a Si/SiGe-HBT with a Realistic Ge Profile

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Abstract

In this paper we present full-band Monte Carlo simulations of an advanced Silicon/Silicon-Germanium Heterojunction Bipolar Transistor. In addition to this 2D MC simulations we performed simulations with a 1D MC model and a hydrodynamic model. For main internal distributions good consistency is found. We conclude with a microscopic investigation of the full-band Monte Carlo results.

1 Introduction

Silicon/Silicon-Germanium Heterojunction Bipolar Transistors (HBT) are promising devices for low-cost high frequency applications in mobile communication. Hence, there is an increasing demand for an accurate modeling of SiGe heterodevices. The solution of Boltzmann's transport equation by means of the full-band Monte Carlo method (FB-MC) relies only on few simplifications of carrier transport physics. Thus, this method is expected to describe device behaviour of heterodevices reliable and can be used to validate classical simulators that are based on macroscopic transport equations. The transport coefficients of our hydrodynamic model (HD) are extracted from homogeneous FB-MC simulations to ensure consistency [1, 2]. The 2D HBT test structure is derived from an advanced HBT from an industrial pilot line process with a realistic Ge-profile [3, 4, 5].

2 Simulation method

The FB-MC code is based on the scattering models reported in [6] with optical phonons and acoustic phonons in the isotropic and elastic equipartition approximation. In SiGe the phonon scattering mechanisms are approximated by considering both Si and Ge phonons and weighting the respective scattering rate by the relative fraction of the component in the alloy. Furthermore alloy scattering is modeled analogous to [7]. Scattering by ionized impurities is taken into account including a doping dependent adjustment to match experimental minority and majority mobility data. The band structures are calculated with the nonlocal empirical pseudopotential method [8]. The valence and conduction bands are discretized in k-space with nonuniform unstructured grids which minimize the

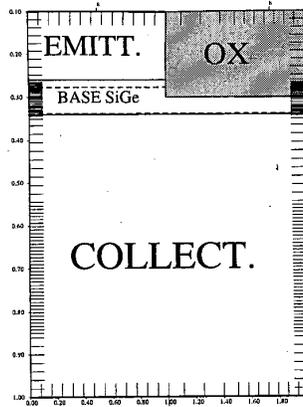


Fig. 1: The 2D SiGe-HBT test structure

discretization error [9]. Bandgap narrowing due to the doping is considered by the traditional model of [10] with the parameters given in [11]. N_T is the total doping.

$$\Delta E_{bgn} = E_0 \left(\ln \frac{N_T}{N_{Ref}} + \sqrt{\left(\ln \frac{N_T}{N_{Ref}} \right)^2 + 0,5} \right)$$

with $E_0 = 0,009 \text{ eV}$, $N_{Ref} = 1,0 \cdot 10^{17} \text{ cm}^{-3}$

Band energy shifts depending on the germanium content x_{Ge} are calculated as reported in [11]. The entire band offsets are given below.

$$\Delta E_c = (-0,228 x_{Ge} + 0,298 x_{Ge}^2) \text{ eV} - \frac{1}{2} \Delta E_{bgn}$$

$$\Delta E_v = (0,732 x_{Ge} + 0,098 x_{Ge}^2) \text{ eV} + \frac{1}{2} \Delta E_{bgn}$$

In Fig. 1 the 2D SiGe HBT test structure is shown. The maximum acceptor doping is $3 \times 10^{19} \text{ cm}^{-3}$. The lowly doped emitter and the collector have a donor doping of about $3 \times 10^{18} \text{ cm}^{-3}$ and $2 \times 10^{16} \text{ cm}^{-3}$, respectively. In the emitter and collector contact region the maximum donor concentration is limited to $1 \times 10^{19} \text{ cm}^{-3}$ in order

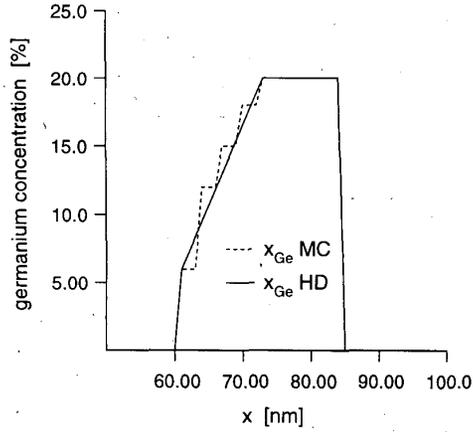


Fig. 2: Germanium profile in the base [3]

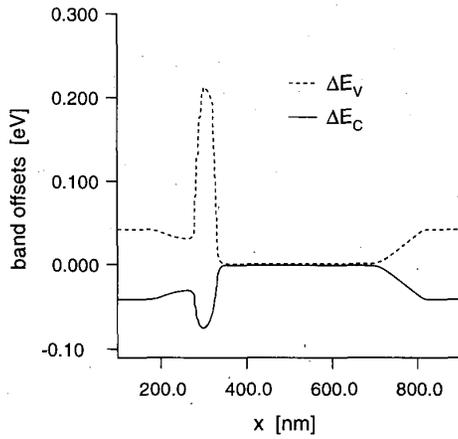


Fig. 3: Band offsets

to reduce the effect of carrier self-heating during the MC simulations. It has been checked by HD simulations that this reduction of the maximum donor doping does not obscure the simulation results in comparison with simulations based on the original doping profile. Arbitrary germanium profiles are resolved spatially in a stepwise manner. The utilized germanium profile is similar to the profile in Fig. 2. The resulting band offsets are shown in Fig. 3, which lead to an energy barrier for electrons passing the emitter/base-junction. Particle transfers via such hetero interfaces are treated consistently with the full-band structure conserving parallel momentum and energy [12].

The profile is approximated for the FB-MC simulation by 6 different germanium concentrations. Although within our FB-MC code any desired number of bandstruc-

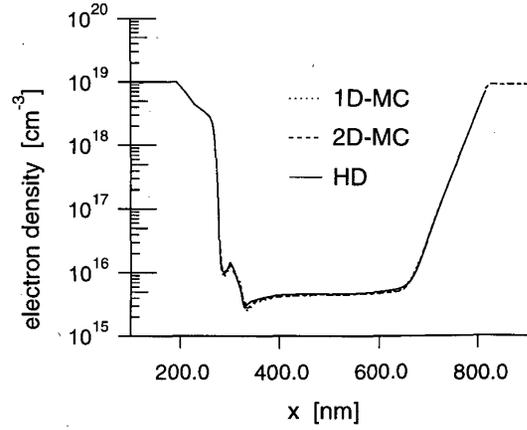


Fig. 4: Comparison of the electron density

tures is possible, the number is limited by the available CPU memory. Less than 50 MBytes are required for one band structure and the total for this simulation is about 320 MBytes including the particle data and statistics. The FB-MC model is solved self-consistently with Poisson's equation. The CPU efficiency of the FB-MC model is similar to the one of the analytical MC model described in [13].

3 Results

The results of the HD, 2D MC and 1D MC simulations are summarized in Fig. 4-6 for the electron density, drift velocity and dynamic temperature. In the case of the 1D MC simulation Poisson's equation is solved with a constant hole fermi level. Thereby the CPU times can be reduced in comparison to the 2D-MC model, because no holes have to be simulated and the dimension of the problem is reduced resulting in a smaller number of electrons.

The bias conditions are $V_{BE} = 0.75V$ and $V_{CB} = 2.0V$. The CPU time for $1ps$ with 140.000 particles is 200 minutes for the 2D-MC model on a 300MHz SUN ULTRA 2. The electron density in Fig. 4 shows a good agreement of all three cases. A significant velocity overshoot is found at the base/collector junction for the HD model and for the MC model as well (Fig. 5). The electron temperature in Fig. 6 also shows a good agreement between HD and MC. Only in the base the electron temperature of the 2D MC simulation is slightly raised compared to the other models, because the minority electrons couple strongly with the plasma oscillations of the majority holes. The overall consistency indicates the reliability of the less CPU intensive HD model. The collector current and total CPU time are compared for HD and 1D-MC in Tab. 1. The collector current agrees well within the stochastic error but

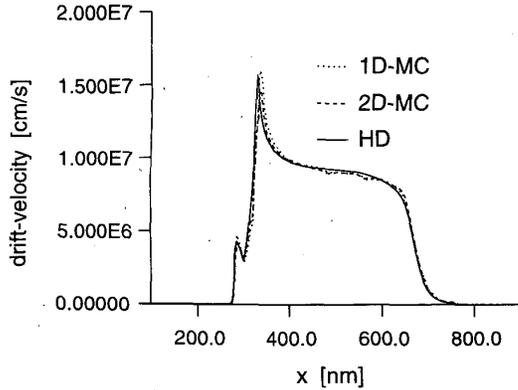


Fig. 5: Comparison of the drift velocity

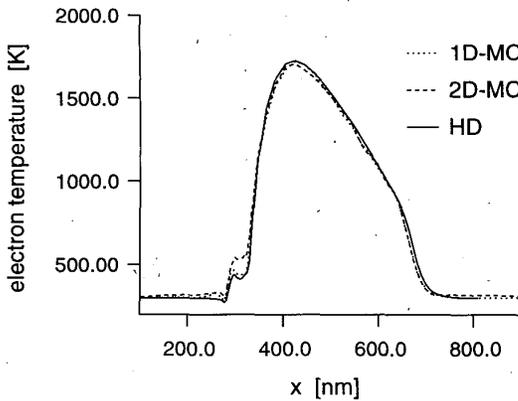


Fig. 6: Comparison of the dynamic temperature

the large difference in CPU time demonstrates the disadvantage of the FB-MC model. The stochastic error in case of the FB-MC simulations is evaluated with the method described in [14]. The calculation is converged with a probability of 95% within $\pm 19\%$. Additionally the HD simulated GUMMELplot in comparison to measurements of our SiGe HBT test structure is shown in Fig. 7. The good agreement demonstrates the advantages of the fast HD model. For more details concerning the HD simulations see [15].

It is one of the definite advantages of the FB-MC simulation that it enables detailed microscopic investigations, which are impossible by standard HD or DD device models. For example in Fig. 8 the valley occupation of the electrons for the FB-MC simulation is depicted. Due to the biaxial compressive strain in the base the two longitudinal valleys of the first conduction band are shifted upwards in energy, whereas the four transversal valleys are

	HD	1D-MC
I-C	0.70 A/cm	0.65 A/cm
CPU-time	40 sec	45 hours
error	-	$\pm 19\%$

Table 1:

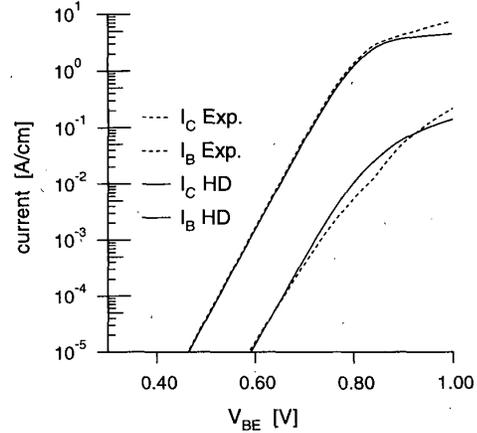


Fig. 7: GUMMELplot at $V_{CB} = 2V$

lowered. Before the base all valleys are uniformly populated, because the electron ensemble is nearly in thermal equilibrium. Since most of the electrons are passing the emitter/base-junction with low energies, the longitudinal valleys are rarely occupied in the base region.

Consequently the transversal-valleys are populated with nearly 25%. In the collector region all six valleys are again equivalent, but the strong electric field increases the population of the longitudinal valleys. In the collector contact the electron gas returns to a state near thermal equilibrium and all valleys are populated by the same number of electrons. Although from a more macroscopic point of view this effect is not important for the test HBT as demonstrated in Fig. 4-6, it might become important for more aggressively scaled devices. Therefore it is very beneficial to have the FB-MC model available for reference simulations. Another important application of the FB-MC model is the examination whether the velocity overshoot in HD simulations is spurious or real. For such investigation the faster 1D-MC model is often sufficient.

4 Conclusion

In conclusion we have reported fully consistent HD and FB-MC simulations of a SiGe HBT with a realistic Ge profile in one and two dimensions. It turns out that the CPU efficient and robust HD model and the MC model

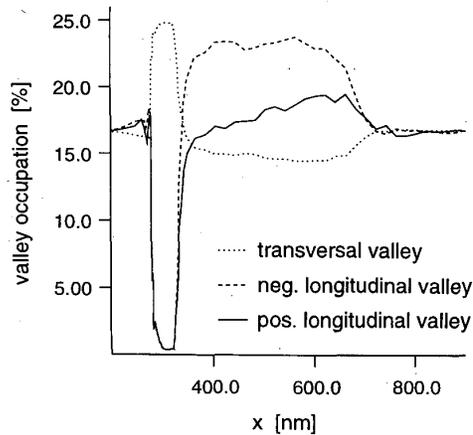


Fig. 8: Valley occupation resulting from MC simulations

agree well for all relevant internal distributions. Therefore, the HD model is currently the best model for the industrial design of SiGe HBTs. The MC model provides more physical insight and is used to validate the accuracy of the HD model.

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