Multigrid Becomes a Competitive Algorithm for some 3D Device Simulation Problems

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Abstract

For situations where Gummel's decoupling scheme is applicable a multigrid algorithm for the continuity equations fully consistent with the usual Scharfetter-Gummel discretization can be used to solve the van Roosbroeck equations. The main problems for applying a multigrid algorithm are that discrete spaces are not nested in the usual sense for the refined grids because of the Scharfetter-Gummel discretization and that problem coefficients vary strongly. Transforming the equations to symmetric form and applying a block MILU decomposition based on the coarse-fine splitting of the discrete spaces with a perturbed Schur complement defines the prolongation and restriction operators. The transformation back to the original variables is possible. Coarse grid matrices are M-matrices.

Let $\mathcal{M}_0, \ldots \mathcal{M}_j$ be a sequence of Euclidean vector spaces with growing dimension. In order to define a standard multigrid algorithm to solve the continuity equation of the van-Roosbroeck system $A_{i}u = f$ on the *j*-th level there need to be the following components (in the terminology of [2]):

- scalar products $((\cdot, \cdot))_k : \mathcal{M}_k \times \mathcal{M}_k \to R$
- symmetric, positive definite with respect to $((\cdot, \cdot))_k$ operators $A_k : \mathcal{M}_k \to \mathcal{M}_k$,
- interpolations $I_k : \mathcal{M}_{k-1} \to \mathcal{M}_k$
- restrictions $P_k^0 : \mathcal{M}_k \to \mathcal{M}_{k-1}$. smoothers $R_k : \mathcal{M}_k \to \mathcal{M}_k$

While the smoothers R_k are provided by one or more steps of a classical iteration method (Jacobi, Gauß-Seidel, ILU), the design of other components in cases of strongly varying coefficients or missing standard finite element background is unclear. Here, we try a 'semi-algebraic' method as described in previous stages in [4, 5, 6]. Rather similar ideas of constructing multigrid or multilevel preconditioners have been used in [1, 3, 9, 10].

In what follows, we abbreviate the level-k-indices, to mean a fine grid corresponds to space \mathcal{M}_k , a coarse grid then corresponds to \mathcal{M}_{k-1} .

On a three-dimensional grid with quadrilateral cells generated by standard refinement from a coarser one, we have the splitting of the grid vertex set $V(A) = V_C \cup V_F \cup V_E \cup V_N$ into sets of coarse grid cell midpoints, coarse grid cell face midpoints, coarse grid

cell edge midpoints and coarse grid node points, respectively. We get the matrix partitioning

$$A = \left(\begin{array}{ccc} \begin{pmatrix} A_{C} & B_{CF} & 0 \\ B_{FC} & A_{F} & B_{FE} \\ 0 & B_{EF} & A_{E} \end{pmatrix} \begin{pmatrix} 0 \\ 0 \\ B_{EN} \end{pmatrix} \\ \begin{pmatrix} 0 & 0 \\ B_{NE} \end{pmatrix} & A_{N} \end{array} \right) = \left(\begin{array}{c} A_{11} & A_{12} \\ A_{21} & A_{22} \end{array} \right),$$

where the off diagonal blocks are nonpositive and the diagonal blocks $A_C = A_{CF} + M_C$, $A_F = A_{FC} + A_{FE} + M_F$, $A_E = A_{EF} + A_{EN} + M_E$ and $A_N = A_{NE} + M_N$, are positive diagonal matrices which consist of sums of off diagonal row entries and a nonnegative "mass" term. The assumptions made on A imply at least one entry of the "mass" M_* is positive.

Let $\tilde{A}_F = A_{FE} + M_F$, $\tilde{A}_E = A_{EN} + M_E$, and choose

$$F = \begin{pmatrix} A_C & B_{CF} & 0\\ 0 & \tilde{A}_F & B_{FE}\\ 0 & 0 & \tilde{A}_E \end{pmatrix}, \qquad G = A_{12}, \qquad U = \begin{pmatrix} F & G\\ 0 & I \end{pmatrix}.$$

U can be seen as a transformation matrix to an approximate harmonic basis [7]. Then for \cdot^T being the transposition with respect to the $((\cdot, \cdot))$ -scalar product we have

$$A = U^T \begin{pmatrix} F^{-T} A_{11} F^{-1} & \Delta \\ \Delta^T & S \end{pmatrix} U,$$
 (1)

with $\Delta = F^{-T}(A_{12} - A_{11}F^{-1}G)$ and $S = \hat{S} + \Delta^T(F^{-T}A_{11}F^{-1})^{-1}\Delta$, and $\hat{S} = A_{22} - A_{21}(A_{11})^{-1}A_{12}$ is the Schur complement. To create a block diagonal preconditioner for A in the new basis, one takes the decomposition (1) and omits the off diagonal blocks Δ . Omitting the error correction in the fine grid part, too, yields a coarse grid correction by projecting the error vector onto the fine grid space in the new basis. It has the form

$$B = U^{-1} \begin{pmatrix} 0 & 0 \\ 0 & S^{-1} \end{pmatrix} U^{-T} = I_k S^{-1} P_k^0$$

with

$$I_{k} = \begin{pmatrix} -F^{-1}G \\ I \end{pmatrix} = \begin{pmatrix} \begin{pmatrix} -A_{C}^{-1}B_{CF}\tilde{A}_{F}^{-1}B_{FE}\tilde{A}_{E}^{-1}B_{EN} \\ \tilde{A}_{F}^{-1}B_{FE}\tilde{A}_{E}^{-1}B_{EN} \\ -\tilde{A}_{E}^{-1}B_{EN} \\ I \end{pmatrix} \end{pmatrix}$$

and P_k^0 being its $((\cdot, \cdot))$ -adjoint. S is the Galerkin coarse grid operator corresponding to the given choice of the intergrid transfer operators:

$$S = P_k^0 A I_k = (A_N - B_{NE} \dot{A}_E^{-1} B_{EN}) + B_{NE} \dot{A}_E^{-1} (A_{EF} - B_{EF} \dot{A}_F^{-1} B_{FE}) \dot{A}_E^{-1} B_{EN} + B_{NE} \tilde{A}_E^{-1} B_{EF} \tilde{A}_F^{-1} (A_{FC} - B_{FC} A_C^{-1} B_{CF}) \tilde{A}_F^{-1} B_{FE} \tilde{A}_E^{-1} B_{EN}.$$

Some geometrical considerations and numerical experiments suggest that in the sense of spectral equivalences, it should hold that $S \approx 4(A_N - B_N E \tilde{A}_E^{-1} B_{EN}) =: A_{k-1}$, when the coefficients are not too strongly varying. This suggests replacing S by A_{k-1} in B. At the other hand, A_{k-1} is the Schur complement of the positive definite matrix

$$4\left(\begin{array}{cc}A_{EN}+M_E & B_{EN}\\B_{NE} & A_N\end{array}\right)$$

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and inherits the $((\cdot, \cdot))$ -symmetry, the *M*-property, and the seven-diagonal structure of *A*, so the process described above can be continued recursively. To ensure the preservation of the *M*-property in a floating point representation, one has to choose a matrix data structure where, on all the levels, instead of the main diagonal entries of *A*, the difference between these entries and the sum of the remaining entries of the same column is stored.

It can be shown [4] that for $((\cdot, \cdot))$ -symmetric, positive definite operators, the convergence of a multigrid method with components defined this way depends on a number of reasonable factors:

- the spectral equivalence of A_{11} and F, and of A_{11} and its diagonal;
- the cosines of the angles between coarse grid and fine grid spaces in the A-energy scalar product;
- the spectral equivalence of S and A_{k-1} ;
- a smoothing property for R_k , which for Jacobi and Gauß-Seidel smoothers is valid for any symmetric *M*-matrix [11]

The whole multigrid operator described above is selfadjoint in the $((\cdot, \cdot))$ -scalar product provided the smoothers are selfadjoint. Without smoothing, the MG-operator has a special recursively defined MILU decomposition interpretation. One can use it as a preconditioner for conjugated gradients in this scalar product. If one considers the Scharfetter-Gummel discretization of carrier transport equations in semiconductors, the discrete operators are selfadjoint with respect to a scalar product using a weight $e^{\pm\psi}$ where ψ is the electrostatic potential.

Here, we compare the algorithm above with a classical iterative one — ILU(1) preconditioning using Chebyshev polynomials and CG with weighted inner product.

The pictures show results for a photo diode with multiple differently doped horizontal layers. The aim is to deplete the whole diode and to compute the recombination current. The kink at 3.5V in the I-U-curve is what the designers are looking for. The Gummel iteration has been stopped at $10^{-7} U_T$ to fulfill the current balance better then 10^{-4} . A second example shows the results for a MOSFET at the 1MBit DRAM design level.



MEDEA, photo diode, Multigrid versus iterative method, 3 grids: $130977 = 49 \times 33 \times 81$, $17425 = 25 \times 17 \times 41$, $2457 = 13 \times 9 \times 21$ points, CONVEX C220



MEDEA, 'half-MOSFET', 1MBit DRAM design level, $U_{\text{bulk}} = -2V$, $U_{\text{drain}} = 0.1V$, 156849 = 97 × 33 × 49, 20825 = 49 × 17 × 25, 2925 = 25 × 9 × 13, 455 = 13 × 5 × 7, 84 = 7 × 3 × 4 points, DEC ALPHA (3000/500, 64MB)

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