

Multiscale Modeling of Hole Avalanche Multiplication and Excess Noise in Amorphous Selenium Semiconductors

Atreyo Mukherjee¹, Richard Akis², Dragica Vasileska³, and A.H. Goldan²

¹ *Department of Electrical Engineering, Stony Brook University, NY, USA*

² *Department of Radiology, School of Medicine, Stony Brook University, NY, USA*

³ *School of Electrical, Computer and Energy Engineering, Arizona State University, AZ, USA*
amirhossein.goldan@stonybrookmedicine.edu

Amorphous selenium (*a*-Se) is a large-area imaging detector material with avalanche gain for low-light and low-dose radiation detection applications.[1] A stark similarity in short range order exists between allotropic forms of selenium, which when augmented with a shift to non-activated extended states transport at high electric fields, allows us to describe the general details of the extended-state hole-phonon interaction in the amorphous phase by modeling band-transport lattice theory of its crystalline counterpart trigonal selenium (*t*-Se). [2] An in-house bulk Monte Carlo algorithm is employed to solve the semiclassical Boltzmann transport equation (MC-BTE). The extended state hole-phonon interaction and the lack of long-range order in *a*-Se is modeled as individual scattering processes, namely acoustic, polar, and non-polar optical phonons, hole-dipole scattering, and impact ionization gain which is modeled using a Keldysh fit. The energy and phonon band structure (**ref. Fig. 1**), along with the valence band density of states (VB-DOS), for *t*-Se, was calculated using density functional theory (DFT) and a non-parabolic approximation to the VB-DOS was used in the simulation (**ref. Fig. 2**). For the first time, we have formulated the interaction of holes in extended states with dipole type defects (the so-called valence alternation pair) present in the chalcogenide family using a non-parabolic VB-DOS approximation. To validate, we compare our drift (**ref. Fig. 3**) and time-of-flight mobility (**ref. Fig. 4**), impact ionization gain (**ref. Fig. 5**), and ensemble energy with experimental findings.[3][4][5][6] This multi-scale approach of combining DFT and MC-BTE is next used in calculating the excess noise factor (ENF) (**ref. Fig. 6**) and spatial resolution in avalanche *a*-Se layers.

[1] Andy LaBella et al., ACS Photon., **6**, 1338-1344, (2019)

[2] Atreyo Mukherjee et al., J. Appl. Phys., **124**, 235102, (2018)

[3] J Mort, Phys. Rev. Lett., **18**, 540, (1967)

[4] Robert Fairman et al., "Semiconducting Chalcogenide Glass I", Academic Press, 2004.

[5] K. Tuji et al., Optoelectron. Devices Technol., **9**, 367-378, (1994).

[6] A. Reznik et al., J. Appl. Phys., **102**, 053711, (2007)

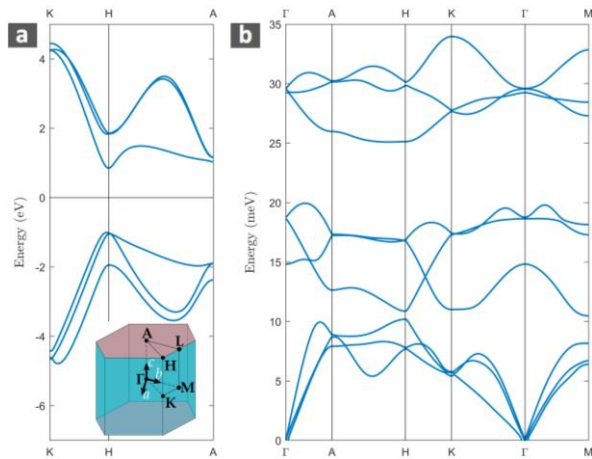


Fig.1: a) DFT calculated electronic band structure for *t*-Se showing a direct bandgap of 1.9 eV at the H point and (b) phonon dispersion. The Brillouin zone for *t*-Se is shown as an inset to Fig. 1.

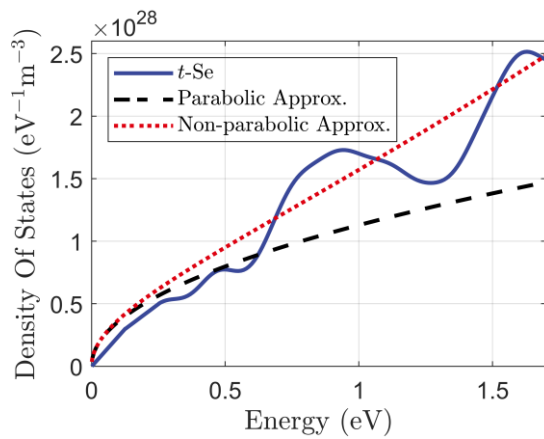


Fig.2: DFT calculations of the VB-DOS are shown by the solid blue line. The dotted red line represents the non-parabolic band approximation ($\alpha = 0.15$) to the VB-DOS. The dashed line represents the parabolic approximation.

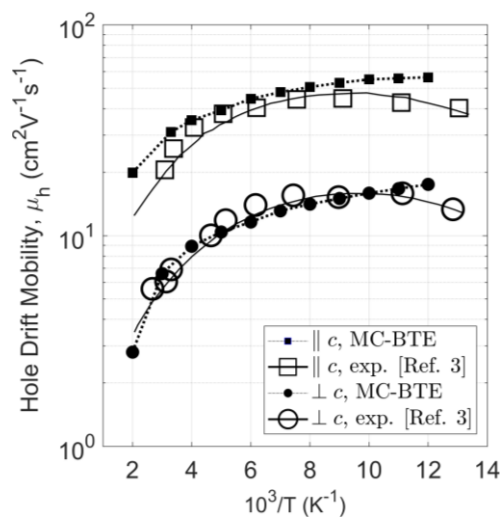


Fig.3: MC-BTE calculated low-field drift mobility values comparison with experiment in the perpendicular and parallel directions to the *c*-axis in *t*-Se, respectively. [3]

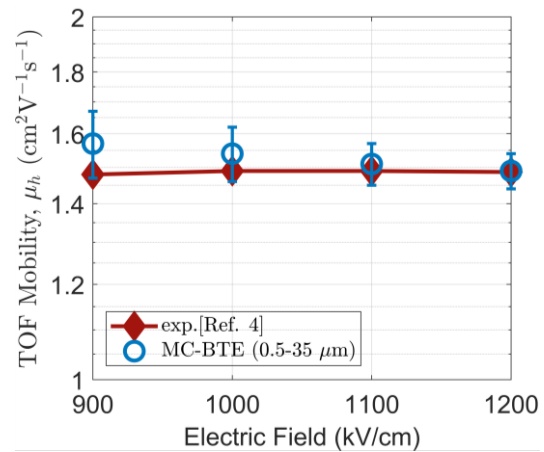


Fig.4: TOF non-parabolic MC-BTE calculated mobility (hollow markers) compared with experimental measured saturated and electric field-independent mobility (solid markers) in *a*-Se. [4]

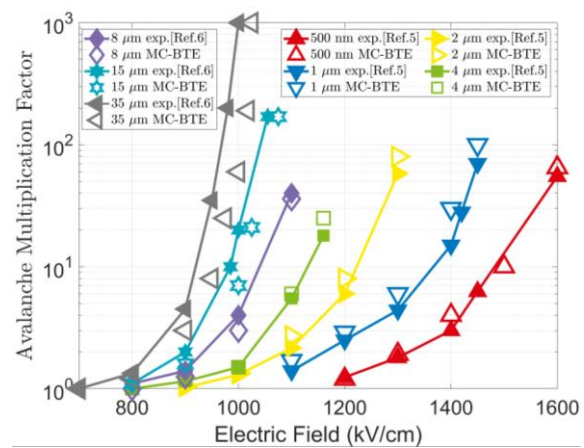


Fig.5: Impact Ionization gain calculated using non-parabolic MC-BTE and compared with experimentally measured gain for 0.5-35 μm thick *a*-Se films. [5][6]

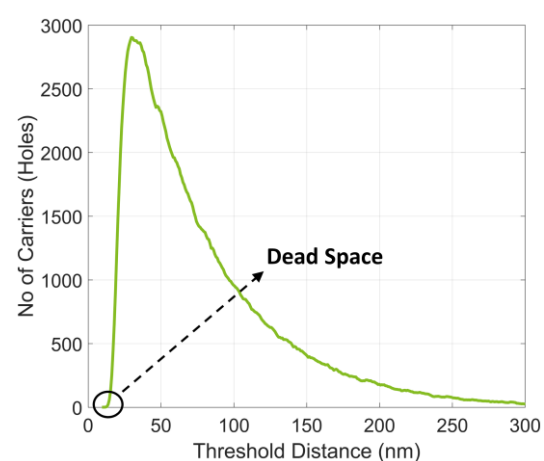


Fig.6: Threshold distance distribution showing the presence of a "dead-space" (defined as the minimum distance of travel in the direction of electric field before a carrier undergoes impact ionization avalanche) in 500 nm *a*-Se layers, resulting in a calculated ENF of ~ 1.4 .